

## 27. On the Radiation Pressure in a Planetary Nebula. I.

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(Comm. by Y. HAGIHARA, M.J.A., June 12, 1950.)

### Abstract.

The radiation pressure of the Lyman  $\alpha$  line-radiation in a planetary nebula is discussed. Zanstra's<sup>1)</sup> idea of redistribution in frequency in the line-contour is taken into account in detail. The equation of transfer of the Lyman  $\alpha$  radiation with redistribution mechanism is solved in contrast with Zanstra's rough treatment in which a definite form of emission and complete redistribution are assumed. The result obtained is found to be nearly the same as in Zanstra's theory. The radiation pressure due to the Lyman  $\alpha$  radiation is so much reduced that we should be able to get rid of the blowing-up difficulty of planetary nebula in Ambarzumian's<sup>2)</sup> theory. Thus it is confirmed that the complete redistribution is a good approximation to the solution of this problem.

### 1. The fundamental equation.

The basic equation for the transfer of the Lyman  $\alpha$  radiation in the shell of a planetary nebula is taken to be

$$\begin{aligned} \cos \theta \frac{dI(\nu, z, \theta)}{\rho dz} = & -I(\nu, z, \theta) \int \sigma \left( \nu \left[ 1 - \frac{1}{c} (vr) \right] \right) \psi(v) dv \\ & + \frac{1}{4\pi} \iint dv dr' \sigma \left( \nu \left[ 1 - \frac{1}{c} (vr) \right] \right) \psi(v) I \left( \nu \left[ 1 - \frac{1}{c} (vr) \right. \right. \\ & \left. \left. + \frac{1}{c} (vr') \right], z, \theta' \right) + Se^{-\tau} \int \sigma \left( \nu \left[ 1 - \frac{1}{c} (vr) \right] \right) \psi(v) dv, \quad (1) \end{aligned}$$

where  $z$ -axis is taken in the direction of direct radiation from the central star,  $I(\nu, z, \theta)$  the intensity of the  $L_\alpha$  radiation at an angle  $\theta$  with the  $z$ -axis at the distance  $z$  from the inner boundary of the nebula,  $\psi(v) dv$  is the well-known Maxwellian velocity distribution of the hydrogen atom in its ground state,  $\sigma$  the natural damping contour, and  $Se^{-\tau}$  is assumed to be the amount of the  $L_\alpha$  emission followed by the absorption of the Lyman continuum at the optical thickness  $\tau$ :

$$S = \frac{\nu_c}{\nu} \frac{\kappa_c}{2\sigma_0} \frac{1}{4} S_c (1 - p),$$

where  $S_c$  is the intensity of the Lyman continuum at the inner boundary  $\nu_c$  the frequency of Lyman limit,  $\kappa_c$  and  $\sigma_0$  are the absorption coefficients at the Lyman limit and the line center of Doppler contour of the  $L_\alpha$ . The first term of the right hand side in the equation (1) is due to the absorption and the second and the third

1) H. Zanstra: B. A. N., **11**, No. 401, 1, 1949.

2) V. A. Ambarzumian: M. N., **93**, 50, 1932. Y. Hagihara: Jap. J. Astr. Geophys., **15**, 1, 1938; **20**, 113, 1943. Hagihara & Hatanaka: *ibid.*, **19**, 135, 1942.