

149. A New Theory of Relativity under the Non-Locally Extended Lorentz Transformation Group

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The present author has established [16] an ameliorated theory of relativity under the group of *extended* Lorentz transformations:

$$(1) \quad \varepsilon_i \xi^i = a_m^i(\xi^m) \varepsilon_m \xi^m + \varepsilon_i a_0^i, \quad (a_0^i = \text{const.}, \varepsilon_i = (-1)^{\frac{1}{2}(1+\delta_i^i)}),$$

$$(2) \quad \varepsilon_i \xi^i = \omega_\mu^i(x^\sigma) \varepsilon_\mu x^\mu + \varepsilon_i \omega_0^i, \quad (\omega_0^i = \text{const.}, \varepsilon_\mu = (-1)^{\frac{1}{2}(1+\delta_\mu^\mu)}),$$

$l, m, \dots; \lambda, \mu, \dots = 1, 2, 3, 4; x^1 = x, x^2 = y, x^3 = z, x^4 = ir = ict; ((x, y, z):$ rectangular Cartesian coordinates, $t = \text{time}$); $(a_m^i(\xi^n))$ and $(\omega_\mu^i(x^\sigma))$: orthogonal matrices with determinant $\neq 0$; $(x^\sigma), (\xi^i)$, and (ξ^i) : II-geodesic rectangular coordinates # [1...16]; δ 's: Kronecker deltas which are 3-dimensional extended equiform Laguerre transformations *, the Einstein space $(R_{\mu\nu} = 0)$ [$dS^2 = g_{\mu\nu}(x^\sigma) dx^\mu dx^\nu = (-1)^{1+\delta_\mu^\mu} \omega^\mu \omega^\mu > 0, g_{\mu\nu} = \omega_\mu^i \omega_\nu^i, \omega^i = \omega_\mu^i(x^\sigma) dx^\mu$] being the map of the Minkowski space (x^σ) by the inverse transformation of the extended Lorentz transformation (2), so that *connection is not necessary* [28]. Thereby the physical interpretations of the geometrical objects were as follows:

$$(3) \quad \left\{ \begin{array}{l} dS = \text{action, } \omega_\mu^i(x^\sigma) = \text{momentum-potential vector; principle of} \\ \text{equivalence} = \text{invariancy of physical laws under the group *} \\ \text{(physical change); "relativity" = referring to \#; physical lines} \\ = \text{II-geodesic curves (straight lines inclusive);} \end{array} \right.$$

(4) *Hamilton's principle:* $\delta S = 0 \rightarrow$ equations of motion:

$$\frac{d^2 \xi^i}{dS^2} = \frac{d}{dS} \frac{\omega^i}{dS} = \omega_\lambda^i \left\{ \frac{d^2 x^\lambda}{dS^2} + A_{\mu\nu}^\lambda \frac{dx^\mu}{dS} \frac{dx^\nu}{dS} \right\} = \omega_\lambda^i \left\{ \frac{d^2 x^\lambda}{dS^2} + \left\{ \begin{array}{c} \lambda \\ \mu\nu \end{array} \right\} \frac{dx^\mu}{dS} \frac{dx^\nu}{dS} \right\} = 0,$$

where

$$(5) \quad A_{\mu\nu}^\lambda = \Omega_\mu^i \partial_\nu \omega_\mu^i \equiv -\omega_\mu^i \partial_\nu \Omega_\mu^i, \quad [\Omega_\mu^i \omega_\mu^i = \delta_\mu^i \iff \Omega_\mu^i \omega_\mu^i = \delta_\mu^i],$$

the (4) representing II-geodesics (in the present author's sense) in 4-dimension, which are in 3-dimension "Kanalf"l"achen" enveloped by oriented II-geodesic spheres (in the present author's sense) with the particle (x^1, x^2, x^3) as center and a II-geodesic radius $r = \int \frac{\omega^4}{dS} dS$. The theory was resumed ([16], p. 623) in the comparison of the present author's theory with the Einstein's, proving the immortal character (comparable with that of the Newton's law) of the former.

In this note, *the said theory will be extended further* by extending the extended Lorentz transformations to "non-locally" ([17-20]) extended Lorentz transformations. *The general procedure consists in considering*