A First Approximate Solution of the Morinaga's Equation: $\frac{\sqrt{\Delta}}{2}\epsilon_{stpq}K_{lm}^{\cdot\cdot pq} = K_{lmst}$.

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1. From the parallel displacement, which makes $ds\psi = 0$ invariant, the differential equation for ψ in the most general form:

$$\frac{\partial\psi}{\partial x^m} = (-\Lambda_m + 2T^5_m\gamma_5 + R_m I)\psi \tag{1.1}$$

has been obtained by K. Morinaga. Specially when the coefficients of connection in the x-space, are Riemannian *i. e.* $\Gamma^i_{jk} = \{^{\,i}_{jk}\}$, (1.1) can be written in

$$\frac{\partial\psi}{\partial x^m} = (\Gamma_m + T^5_m\gamma_5 - L_m I)\psi\,, \tag{1.2}$$

where Γ_m is determined uniquely from the equation:

$$\frac{\partial\gamma_i}{\partial x^m} = \{^{\,j}_{im}\}\gamma_j + \Gamma_m\gamma_i - \gamma_i\Gamma_m \tag{1.3}$$

and T^5_m, L_m may be any functions of x's.

In this paper in order to apply this theory to the problem of gravitation, I shall consider only the case in which $T^5_m = 0$, $L_m = 0$. In our case, therefore, the fundamental equation for ψ becomes

$$\frac{\partial\psi}{\partial x^m} = \Gamma_m\psi\,. \tag{1.4}$$

Further if we assume that $ds^2 = g_{ij}dx^i dx^j$ is a positive definite where $g_{ij} = \gamma_{(i}\gamma_{j)}$, then the condition of integrability of (1.4) is either