

**SINGULAR LIMIT APPROACH  
TO STABILITY AND BIFURCATION FOR  
BISTABLE REACTION DIFFUSION SYSTEMS**

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**1. Introduction.** Patterns with sharp transition layers appear in various fields such as patchiness and segregation in ecosystems [3,10], traveling waves in excitable media [2,4,5, and 22], striking patterns in morphogenesis models [13], dendric patterns in solidification problem [1], and so on.

The most simple but substantial model system, to which most of the above ones fall, is given by the following reaction-diffusion equations in one-dimensional space:

$$(P) \quad \begin{aligned} u_s &= d_1 u_{xx} + f(u, v) \\ v_s &= d_2 v_{xx} + \delta g(u, v) \end{aligned} \quad \text{on } I,$$

where  $d_1$  and  $d_2$  are the diffusion rates of  $u$  and  $v$ , and  $\delta$  is the ratio of the reaction rates. The interval  $I$  is either  $(-l, l)$  or  $\mathbf{R}$ . The Neumann boundary conditions  $u_x = 0 = v_x$  is added to (P), when  $I = (-l, l)$ . It is usually assumed in (P) that one of the following conditions holds:

- (a) There is a significant difference in the diffusion rates of  $u$  and  $v$ .
- (b) There is a significant difference in the reaction rates of  $u$  and  $v$ .
- (c) There is a combination of (a) and (b).

Most of the symmetry breaking stationary patterns in the framework of Turing's diffusion driven instability fall into the first category. One of the well-known models is the Gierer and Meinhardt equation describing morphogenetic patterns [13]. Propagator-controller systems including a simple skeleton model for the B-Z reaction lie on the second category [23, 5]. It is essential for such systems that one of the components reacts much faster than the other. Formally speaking, the FitzHugh-Nagumo equations belong to the third category in which the second component  $v$  does not diffuse and reacts much slower than the first one [8]. However, the qualitative behavior of solutions of the FHN

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