

The xi function

by

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1. Introduction

Despite the fact that spectral and inverse spectral properties of one-dimensional Schrödinger operators $H = -d^2/dx^2 + V$ have been extensively studied for seventy-five years, there remain large areas where our knowledge is limited. For example, while the inverse theory for operators on $L^2((-\infty, \infty))$ is well understood in case V is periodic [12], [24], [25], [35], [39]–[42], [49], it is not understood in case $\lim_{|x| \rightarrow \infty} V(x) = \infty$ and H has discrete spectrum.

Our goal here is to introduce a special function $\xi(x, \lambda)$ on $\mathbf{R} \times \mathbf{R}$ associated to H which we believe will be a valuable tool in the spectral and inverse spectral theory. In a sense we will make precise, it complements the Weyl m -functions, $m_{\pm}(x, \lambda)$.

A main application of ξ which we will make here concerns a generalization of the trace formula for Schrödinger operators to general V 's.

Recall the well-known trace formula for periodic potentials: Let $V(x) = V(x+1)$. Then, by Floquet theory (see, e.g., [10], [37], [44]),

$$\text{spec}(H) = [E_0, E_1] \cup [E_2, E_3] \cup \dots,$$

a set of bands. If V is C^1 , one can show that the sum of the gap sizes is finite, that is,

$$\sum_{n=1}^{\infty} |E_{2n} - E_{2n-1}| < \infty. \quad (1.1)$$

For a fixed y , let H_y be the operator $-d^2/dx^2 + V$ on $L^2([y, y+1])$ with $u(y) = u(y+1) = 0$ boundary conditions. Its spectrum is discrete, that is, there are eigenvalues $\{\mu_n(y)\}_{n=1}^{\infty}$ with

$$E_{2n-1} \leq \mu_n(y) \leq E_{2n}. \quad (1.2)$$