

Values of Brownian intersection exponents, I: Half-plane exponents

by

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1. Introduction

Theoretical physics predicts that conformal invariance plays a crucial role in the macroscopic behavior of a wide class of two-dimensional models in statistical physics (see, e.g., [4], [6]). For instance, by making the assumption that critical planar percolation behaves in a conformally invariant way in the scaling limit, and using ideas involving conformal field theory, Cardy [7] produced an exact formula for the limit, as $N \rightarrow \infty$, of the probability that, in two-dimensional critical percolation, there exists a cluster crossing the rectangle $[0, aN] \times [0, bN]$. Also, Duplantier and Saleur [13] predicted the “fractal dimension” of the hull of a very large percolation cluster. These are just two examples among many such predictions.

In 1988, Duplantier and Kwon [12] suggested that the ideas of conformal field theory can also be applied to predict the intersection exponents between random walks in \mathbf{Z}^2 (and Brownian motions in \mathbf{R}^2). They predicted, for instance, that if B and B' are independent planar Brownian motions (or simple random walks in \mathbf{Z}^2) started from distinct points in the upper half-plane $\mathbf{H} = \{(x, y) : y > 0\} = \{z \in \mathbf{C} : \text{Im}(z) > 0\}$, then when $n \rightarrow \infty$,

$$\mathbf{P}[B[0, n] \cap B'[0, n] = \emptyset] = n^{-\zeta + o(1)} \quad (1.1)$$

and

$$\mathbf{P}[B[0, n] \cap B'[0, n] = \emptyset \text{ and } B[0, n] \cup B'[0, n] \subset \mathbf{H}] = n^{-\tilde{\zeta} + o(1)}, \quad (1.2)$$

where

$$\zeta = \frac{5}{8}, \quad \tilde{\zeta} = \frac{5}{3}.$$