## Research Article

# Constrained $C^{0}$ Finite Element Methods for Biharmonic Problem 

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This paper presents some constrained $C^{0}$ finite element approximation methods for the biharmonic problem, which include the $C^{0}$ symmetric interior penalty method, the $C^{0}$ nonsymmetric interior penalty method, and the $C^{0}$ nonsymmetric superpenalty method. In the finite element spaces, the $C^{1}$ continuity across the interelement boundaries is obtained weakly by the constrained condition. For the $C^{0}$ symmetric interior penalty method, the optimal error estimates in the broken $H^{2}$ norm and in the $L^{2}$ norm are derived. However, for the $C^{0}$ nonsymmetric interior penalty method, the error estimate in the broken $H^{2}$ norm is optimal and the error estimate in the $L^{2}$ norm is suboptimal because of the lack of adjoint consistency. To obtain the optimal $L^{2}$ error estimate, the $C^{0}$ nonsymmetric superpenalty method is introduced and the optimal $L^{2}$ error estimate is derived.

## 1. Introduction

The discontinuous Galerkin methods (DGMs) have become a popular method to deal with the partial differential equations, especially for nonlinear hyperbolic problem, which exists the discontinuous solution even when the data is well smooth, and the convection-dominated diffusion problem, and the advection-diffusion problem. For the second-order elliptic problem, according to the different numerical fluxes, there exist different discontinuous Galerkin methods, such as the interior penalty method (IP), the nonsymmetric interior penalty method (NIPG), and local discontinuous Galerkin method (LDG). A unified analysis of discontinuous Galerkin methods for the second-order elliptic problem is studied by Arnold et al. in [1].

The DGM for the fourth-order elliptic problem can be traced back to 1970s. Baker in [2] used the IP method to study the biharmonic problem and obtained the optimal error estimates. Moreover, for IP method, the $C^{0}$ and $C^{1}$ continuity can be achieved weakly by the interior penalty. Recently, using IP method and NIPG method, Süli and Mozolevski in [3-5] studied the $h p$-version DGM for the biharmonic problem, where the error estimates
are optimal with respect to the mesh size $h$ and are suboptimal with respect to the degree of the piecewise polynomial approximation $p$. However, we observe that the bilinear forms and the norms corresponding to the IP method in [3-5] are much complicated. A method to simplify the bilinear forms and the norms is using $C^{0}$ interior penalty method. $C^{0}$ interior penalty method for the biharmonic problem was introduced by Babuška and Zlámal in [6], where they used the nonconforming element and considered the inconsistent formulation and obtained the suboptimal error estimate. Motivated by the Engel and his collaborators' work [7], Brenner and Sung in [8] studied the $C^{0}$ interior penalty method for fourth-order problem on polygonal domains. They used the $C^{0}$ finite element solution to approximate $C^{1}$ solution by a postprocessing procedure, and the $C^{1}$ continuity can be achieved weakly by the penalty on the jump of the normal derivatives on the interelement boundaries.

In this paper, thanks to Rivière et al.'s idea in [9], we will study some constrained $C^{0}$ finite element approximation methods for the biharmonic problem. The $C^{1}$ continuity can be weakly achieved by a constrained condition that integrating the jump of the normal derivatives over the inter-element boundaries vanish. Under this constrained condition, we discuss three $C^{0}$ finite element methods which include the $C^{0}$ symmetric interior penalty method based on the symmetric bilinear form, the $C^{0}$ nonsymmetric interior penalty method, and $C^{0}$ nonsymmetric superpenalty method based on the nonsymmetric bilinear forms. First, we study the $C^{0}$ symmetric interior penalty method and obtain the optimal error estimates in the broken $H^{2}$ norm and in $L^{2}$ norm. However, for the $C^{0}$ nonsymmetric interior penalty method, the $L^{2}$ norm is suboptimal because of the lack of adjoint consistency. Finally, in order to improve the order of the $L^{2}$ error estimate, we give the $C^{0}$ nonsymmetric superpenalty method and show the optimal $L^{2}$ error estimates.

## 2. $C^{0}$ Finite Element Approximation

Let $\Omega \subset \mathbb{R}^{2}$ be a bounded and convex domain with boundary $\partial \Omega$. Consider the following biharmonic problem:

$$
\begin{gather*}
\Delta^{2} u=f, \quad \text { in } \Omega, \\
u=\frac{\partial u}{\partial n}=0, \quad \text { on } \partial \Omega, \tag{2.1}
\end{gather*}
$$

where $n$ denotes the unit external normal vector to $\partial \Omega$. We assume that $f$ is sufficiently smooth such that the problem (2.1) admits a unique solution $u \in H^{4}(\Omega) \cap H_{0}^{2}(\Omega)$.

Let $\tau_{h}$ be a family of nondegenerate triangular partition of $\Omega$ into triangles. The corresponding ordered triangles are denoted by $K_{1}, K_{2}, \ldots, K_{N}$. Let $h_{i}=\operatorname{diam}\left(K_{i}\right), i=$ $1, \ldots, N$, and $h=\max \left\{h_{1}, h_{2}, \ldots, h_{N}\right\}$. The nondegenerate requirement is that there exists $\rho>0$ such that $K_{i}$ contains a ball of radius $\rho h_{i}$ in its interior. Conventionally, the boundary of $K_{i}$ is denoted by $\partial K_{i}$. We denote

$$
\begin{equation*}
e_{i j}=\partial K_{i} \cap \partial K_{j}, \quad e_{i}=\partial K_{i} \cap \partial \Omega, \quad h_{e_{i j}}=\operatorname{diam}\left(e_{i j}\right), \quad h_{e_{i}}=\operatorname{diam}\left(e_{i}\right) \tag{2.2}
\end{equation*}
$$

Assume that the partition $\tau_{h}$ is quasiuniform; that is, there exists a positive constant $\mathcal{v}$ such that

$$
\begin{equation*}
h \leqslant v h_{i}, \quad i=1, \ldots, N \tag{2.3}
\end{equation*}
$$

Let $\boldsymbol{\varepsilon}_{I}$ and $\boldsymbol{\varepsilon}_{B}$ be the set of interior edges and boundary edges of $\boldsymbol{\tau}_{h}$, respectively. Let $\mathfrak{\mathcal { E }}=$ $\mathcal{E}_{I} \cup \mathfrak{\varepsilon}_{B}$. Denote by $v^{i}$ the restriction of $v$ to $K_{i}$. Let $e=e_{i j} \in \mathcal{E}_{I}$ with $i>j$. Then we denote the jump [v] and the average $\{v\}$ of $v$ on $e$ by

$$
\begin{equation*}
\left.[v]\right|_{e}=\left.v^{i}\right|_{e}-\left.\left.v^{j}\right|_{e^{\prime}} \quad\{v\}\right|_{e}=\frac{1}{2}\left(\left.v^{i}\right|_{e}+\left.v^{j}\right|_{e}\right) . \tag{2.4}
\end{equation*}
$$

If $e=e_{i} \in \mathcal{E}_{B}$, we denote [ $\left.v\right]$ and $\{v\}$ of $v$ on $e$ by

$$
\begin{equation*}
\left.[v]\right|_{e}=\left.\{v\}\right|_{e}=\left.v^{i}\right|_{e} . \tag{2.5}
\end{equation*}
$$

Define $V$ by

$$
\begin{equation*}
V=\left\{v \in H_{0}^{1}(\Omega),\left.v\right|_{K} \in H^{s}(K), \forall K \in \tau_{h}\right\}, \quad s \geqslant 3 \tag{2.6}
\end{equation*}
$$

with broken $H^{s}$ norm

$$
\begin{equation*}
\|\mid v\|_{s}=\left(\sum_{K \in \tau_{h}}\|v\|_{s, K}^{2}\right)^{1 / 2}, \quad \forall v \in V, \tag{2.7}
\end{equation*}
$$

where $\|\cdot\|_{s, K}$ is the standard Sobolev norm in $H^{s}(K)$. Define the broken $H^{2}$ norm by

$$
\begin{equation*}
\|v\|_{h}=\left(\sum_{K \in \mathcal{I}_{h}}|v|_{2, K}^{2}\right)^{1 / 2}, \quad \forall v \in V, \tag{2.8}
\end{equation*}
$$

where $|\cdot|_{2, K}$ is the seminorm in $H^{2}(K)$.
For every $K \in \tau_{h}$ and any $v \in V$, we apply the integration by parts formula to obtain

$$
\begin{align*}
\int_{K}\left(\Delta^{2} u\right) v d x & =-\int_{K} \nabla(\Delta u) \cdot \nabla v d x+\int_{\partial K} \frac{\partial \Delta u}{\partial n} v d s \\
& =\int_{K} \Delta u \Delta v d x-\int_{\partial K} \Delta u \frac{\partial v}{\partial n} d s+\int_{\partial K} \frac{\partial \Delta u}{\partial n} v d s . \tag{2.9}
\end{align*}
$$

Summing all $K \in \tau_{h}$, we have

$$
\begin{align*}
& \sum_{K \in \mathcal{\tau}_{n}} \int_{K} \Delta u \Delta v d x-\sum_{e \in \mathcal{E}_{I}} \int_{e}\{\Delta u\}\left[\frac{\partial v}{\partial n}\right]+[\Delta u]\left\{\frac{\partial v}{\partial n}\right\} d s+\sum_{e \in \mathcal{E}_{I}} \int_{e}\left[\frac{\partial \Delta u}{\partial n}\right]\{v\}  \tag{2.10}\\
& \quad+\left\{\frac{\partial \Delta u}{\partial n}\right\}[v] d s-\int_{\partial \Omega} \Delta u \frac{\partial v}{\partial n} d s=\int_{\Omega} f v d x .
\end{align*}
$$

Since $u \in H^{4}(\Omega) \cap H_{0}^{2}(\Omega)$ and $v \in V$, then $[\Delta u]=[\partial \Delta u / \partial n]=[v]=0$ on $e \in \varepsilon_{I}$. Thus, the previous identity can be simplified as follows:

$$
\begin{equation*}
\sum_{K \in \tau_{h}} \int_{K} \Delta u \Delta v d x-\sum_{e \in \mathcal{E}_{I}} \int_{e}\{\Delta u\}\left[\frac{\partial v}{\partial n}\right] d s-\int_{\partial \Omega} \Delta u \frac{\partial v}{\partial n} d s=\int_{\Omega} f v d x \tag{2.11}
\end{equation*}
$$

Now, we introduce the following two bilinear forms:

$$
\begin{align*}
a_{S}(u, v)= & \sum_{K \in \mathcal{Z}_{h}} \int_{K} \Delta u \Delta v d x-\sum_{e \in \mathcal{E}} \int_{e}\{\Delta u\}\left[\frac{\partial v}{\partial n}\right] d s-\sum_{e \in \mathcal{E}} \int_{e}\{\Delta v\}\left[\frac{\partial u}{\partial n}\right] d s \\
& +\sum_{e \in \mathcal{E}} \frac{\beta}{h_{e}} \int_{e}\left[\frac{\partial u}{\partial n}\right]\left[\frac{\partial v}{\partial n}\right] d s,  \tag{2.12}\\
a_{\mathrm{NS}}(u, v)= & \sum_{K \in \mathcal{Z}_{h}} \int_{K} \Delta u \Delta v d x-\sum_{e \in \mathcal{E}} \int_{e}\{\Delta u\}\left[\frac{\partial v}{\partial n}\right] d s+\sum_{e \in \mathcal{E}} \int_{e}\{\Delta v\}\left[\frac{\partial u}{\partial n}\right] d s \\
& +\sum_{e \in \mathcal{E}} \frac{1}{h_{e}} \int_{e}\left[\frac{\partial u}{\partial n}\right]\left[\frac{\partial v}{\partial n}\right] d s .
\end{align*}
$$

It is clear that $a_{S}(\cdot, \cdot)$ is a symmetric bilinear form and $a_{\mathrm{NS}}(\cdot, \cdot)$ is a nonsymmetric bilinear form. In terms of (2.11) and

$$
\begin{equation*}
\sum_{e \in \mathcal{E}_{I}} \int_{e}\{\Delta v\}\left[\frac{\partial u}{\partial n}\right] d s+\int_{\partial \Omega} \Delta v \frac{\partial u}{\partial n} d s=0 \tag{2.13}
\end{equation*}
$$

the solution $u$ to problem (2.1) satisfies the following variational problems:

$$
\begin{array}{cl}
a_{S}(u, v)=(f, v), & \forall v \in V \\
a_{N S}(u, v)=(f, v), & \forall v \in V \tag{NSP}
\end{array}
$$

Let $P_{r}(K)$ denote the space of the polynomials on $K$ of degree at most $r$. Define the following constrained $C^{0}$ finite element space:

$$
\begin{equation*}
V_{h}=\left\{v_{h} \in H_{0}^{1}(\Omega),\left.v_{h}\right|_{K} \in P_{r}(K), r \geqslant 4, \forall K \in \tau_{h}, \int_{e}\left[\frac{\partial v_{h}}{\partial n}\right] d s=0, \forall e \in \varepsilon\right\} \tag{2.14}
\end{equation*}
$$

from which we note that the $C^{1}$ continuity of $v_{h} \in V_{h}$ can be weakly achieved by the constrained condition $\int_{e}\left[\partial v_{h} / \partial n\right] d s=0$ for all $e \in \mathcal{\varepsilon}$. Next, we define the degrees of freedom for this finite element space. To this end, for any $K \in \tau_{h}$, denote by $p_{i}(i=1,2,3)$ the three
vertices of $K$. Recall that the degrees of freedom of Lagrange element on $K$ are $v(p)$, for all $p \in \mathcal{C}^{\prime}$ with (cf. [10])

$$
\begin{equation*}
\mathcal{C}^{\prime}:=\left\{p=\sum_{j=1}^{3} \lambda_{j} p_{j} ; \sum_{j=1}^{3} \lambda_{j}=1, \lambda_{j} \in\left\{0, \frac{1}{r}, \ldots, \frac{r-1}{r}, 1\right\}, 1 \leqslant j \leqslant 3\right\} \tag{2.15}
\end{equation*}
$$

Then we modify the degrees of freedom of Lagrange element to suit the constraint of normal derivatives over the edges in $V_{h}$. Specifically speaking, the degrees of freedom of $V_{h}$ are given by

$$
\begin{gather*}
v(p), \quad \forall p \in \mathcal{C} \text { with } \mathcal{C}:=\mathcal{C}^{\prime} \backslash\left\{\frac{1}{r}\left(p_{1}+p_{2}+p_{3}\right)+\frac{r-3}{r} p_{i}, i=1,2,3\right\}, \\
\int_{e} \frac{\partial v}{\partial n} d s \quad \text { for each edge } e \text { of } K . \tag{2.16}
\end{gather*}
$$

Based on the symmetric bilinear form $a_{S}(\cdot, \cdot)$, the $C^{0}$ symmetric interior penalty finite element approximation of (2.1) is

$$
\begin{align*}
& \text { find } u_{h}^{S} \in V_{h} \text { such that, } \\
& a_{S}\left(u_{h}^{S}, v_{h}\right)=\left(f, v_{h}\right), \quad \forall v_{h} \in V_{h} . \tag{2.17}
\end{align*}
$$

Based on the nonsymmetric bilinear form $a_{\mathrm{NS}}(\cdot, \cdot)$, the $C^{0}$ nonsymmetric interior penalty finite element approximation of (2.1) is

$$
\begin{align*}
& \text { find } u_{h}^{\mathrm{NS}} \in V_{h} \text { such that, } \\
& a_{\mathrm{NS}}\left(u_{h}^{\mathrm{NS}}, v_{h}\right)=\left(f, v_{h}\right), \quad \forall v_{h} \in V_{h} . \tag{2.18}
\end{align*}
$$

Moreover, the following orthogonal equations hold:

$$
\begin{gather*}
a_{S}\left(u_{h}^{S}-u, v_{h}\right)=0, \quad \forall v_{h} \in V_{h}  \tag{2.19}\\
a_{\mathrm{NS}}\left(u_{h}^{\mathrm{NS}}-u, v_{h}\right)=0, \quad \forall v_{h} \in V_{h} \tag{2.20}
\end{gather*}
$$

In order to introduce a global interpolation operator, we first define $\phi_{h}^{i} \in P_{r}\left(K_{i}\right)$ for $\phi \in H^{s}\left(K_{i}\right)$ and $K_{i} \in \tau_{h}$ according to the degrees of freedom of $V_{h}$ by

$$
\begin{gather*}
\phi_{h}^{i}(p)=\phi(p), \quad \text { if } p \in \mathcal{C} \backslash \partial \Omega, \quad \phi_{h}^{i}(p)=0, \quad \text { if } p \in \mathcal{C} \cap \partial \Omega \\
\int_{e} \frac{\partial \phi_{h}^{i}}{\partial n} d s=\int_{e} \frac{\partial \phi}{\partial n} d s, \quad \text { if edge } e \subset \partial K_{i} \backslash \partial \Omega, \quad \int_{e} \frac{\partial \phi_{h}^{i}}{\partial n} d s=0, \quad \text { if edge } e \subset \partial K_{i} \cap \partial \Omega \tag{2.21}
\end{gather*}
$$

Due to standard scaling argument and Sobolev embedding theorem (cf. [10]), we have that for every $K_{i} \in \tau_{h}, e \in \partial K_{i}, \phi \in H^{s}\left(K_{i}\right)$ with $s \geqslant 2, i=1, \ldots, N$,

$$
\begin{gather*}
\left\|\phi-\phi_{h}^{i}\right\|_{q, K_{i}} \leqslant c h^{\mu-q}\|\phi\|_{s, K_{i}^{\prime}} \quad 0 \leqslant q \leqslant \mu \\
\left\|\phi-\phi_{h}^{i}\right\|_{m, e} \leqslant c h_{e}^{\mu-1 / 2-m}\|\phi\|_{s, K_{i}{ }^{\prime}} \quad m=0,1,2  \tag{2.22}\\
\left\|\left[\frac{\partial\left(\phi-\phi_{h}^{i}\right)}{\partial n}\right]\right\|_{0, e} \leqslant c h_{e}^{\mu-3 / 2}\|\phi\|_{s, K_{i}^{\prime}}
\end{gather*}
$$

where $\mu=\min \{s, r+1\}$ and $c>0$ is independent of $h$. We also suppose that the following inverse inequalities hold:

$$
\begin{equation*}
\left\|\phi_{h}^{i}\right\|_{0, e} \leqslant c h_{e}^{-1 / 2}\left\|\phi_{h}^{i}\right\|_{0, K_{i}^{\prime}} \quad\left\|\phi_{h}^{i}\right\|_{1, e} \leqslant c h_{e}^{-1 / 2}\left\|\phi_{h}^{i}\right\|_{1, K_{i}^{\prime}} \quad\left\|\phi_{h}^{i}\right\|_{1, K_{i}} \leqslant c h^{-1}\left\|\phi_{h}^{i}\right\|_{0, K_{i}} \tag{2.23}
\end{equation*}
$$

where $c>0$ is independent of $h$. Then for every $\phi \in V$, we define the global interpolation operator $P_{h}: V \rightarrow V_{h}$ by $\left.P_{h} \phi\right|_{K_{i}}=\phi_{h}^{i}$. Moreover, from (2.22) there holds

$$
\begin{equation*}
\left\|\phi-P_{h} \phi\right\|_{h} \leqslant c h^{\mu-2}\||\phi|\|_{s^{\prime}} \tag{2.24}
\end{equation*}
$$

where $c>0$ is independent of $h$.
The following lemma is useful to establish the existence and uniqueness of the finite element approximation solution.

Lemma 2.1. There exists some constant $\mathcal{\kappa}_{0}>0$ independent of $h$ such that

$$
\begin{equation*}
\mathcal{\kappa}_{0}\left\|v_{h}\right\|_{h}^{2} \leqslant \sum_{K \in \mathcal{\tau}_{h}}\left\|\Delta v_{h}\right\|_{0, K}^{2}+\sum_{e \in \mathcal{E}} \frac{1}{h_{e}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}, \quad \forall v_{h} \in V_{h} \tag{2.25}
\end{equation*}
$$

Proof. Introduce a $H_{0}^{2}$ conforming finite element space $Z_{h}$ thanks to Guzmán and Neilan [11]. The advantage of $Z_{h}$ is that the degrees of freedom depend only on the values of functions and their first derivatives. Denote by $L_{h}$ the interpolation operator from $V_{h}$ to $Z_{h}$. Then there holds

$$
\begin{equation*}
\sum_{K \in \mathcal{\tau}_{h}}\left|v_{h}-L_{h} v_{h}\right|_{2, K}^{2} \leqslant C_{1} \sum_{e \in \mathcal{E}} \frac{1}{h_{e}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}, \quad \forall v_{h} \in V_{h} \tag{2.26}
\end{equation*}
$$

where $C_{1}>0$ is independent of $h$. Thus, we have

$$
\begin{align*}
\left\|v_{h}\right\|_{h}^{2} & \leqslant \sum_{K \in \mathcal{\tau}_{h}}\left|v_{h}-L_{h} v_{h}\right|_{2, K}^{2}+\sum_{K \in \mathcal{\tau}_{h}}\left|L_{h} v_{h}\right|_{2, K}^{2} \leqslant C_{1} \sum_{e \in \mathcal{E}} \frac{1}{h_{e}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}+C_{2} \sum_{K \in \mathcal{\tau}_{h}}\left\|\Delta\left(L_{h} v_{h}\right)\right\|_{0, K}^{2} \\
& \leqslant C_{1} \sum_{e \in \mathcal{E}} \frac{1}{h_{e}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}+C_{2} \sum_{K \in \mathcal{\tau}_{h}}\left\|\Delta\left(L_{h} v_{h}-v_{h}\right)\right\|_{0, K}^{2}+C_{2} \sum_{K \in \mathcal{\tau}_{h}}\left\|\Delta v_{h}\right\|_{0, K}^{2} \\
& \leqslant C_{1} \sum_{e \in \mathcal{E}} \frac{1}{h_{e}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}+C_{2} \sum_{K \in \mathcal{\tau}_{h}}\left|L_{h} v_{h}-v_{h}\right|_{2, K}^{2}+C_{2} \sum_{K \in \mathcal{\tau}_{h}}\left\|\Delta v_{h}\right\|_{0, K}^{2} \\
& \leqslant C_{1}\left(1+C_{2}\right) \sum_{e \in \mathcal{\varepsilon}} \frac{1}{h_{e}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}+C_{2} \sum_{K \in \mathcal{\tau}_{h}}\left\|\Delta v_{h}\right\|_{0, K}^{2} \tag{2.27}
\end{align*}
$$

which completes the proof of $(2.25)$ with $\kappa_{0}=1 / \max \left\{C_{1}\left(1+C_{2}\right), C_{2}\right\}$.

## 3. $C^{0}$ Symmetric Interior Penalty Method

In this section, we will show the optimal error estimates in the broken $H^{2}$ norm and in the $L^{2}$ norm between the solution $u$ to problem (2.1) and the solution $u_{h}^{S}$ to the problem (2.17). First, concerning the symmetric $a_{S}(\cdot, \cdot)$, we have the following coercive property in $V_{h}$.

Lemma 3.1. For sufficiently large $\beta$, there exists some constant $\kappa_{1}>0$ such that

$$
\begin{equation*}
a_{S}\left(v_{h}, v_{h}\right) \geqslant \kappa_{1}\left\|v_{h}\right\|_{h}^{2}, \quad \forall v_{h} \in V_{h} \tag{3.1}
\end{equation*}
$$

Proof. According to the definition of $a_{S}(\cdot, \cdot)$, we have

$$
\begin{equation*}
a_{S}\left(v_{h}, v_{h}\right)=\sum_{K \in \mathcal{\tau}_{h}} \int_{K}\left|\Delta v_{h}\right|^{2} d x-2 \sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta v_{h}\right\}\left[\frac{\partial v_{h}}{\partial n}\right] d s+\sum_{e \in \mathcal{E}} \frac{\beta}{h_{e}} \int_{e}\left|\left[\frac{\partial v_{h}}{\partial n}\right]\right|^{2} d s \tag{3.2}
\end{equation*}
$$

Using the Hölder's inequality and the Young's inequality, we have

$$
\begin{align*}
\sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta v_{h}\right\}\left[\frac{\partial v_{h}}{\partial n}\right] d s & \leqslant\left(\sum_{e \in \mathcal{E}} h_{e}\left\|\left\{\Delta v_{h}\right\}\right\|_{0, e}^{2}\right)^{1 / 2}\left(\sum_{e \in \mathcal{E}} h_{e}^{-1}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2} \\
& \leqslant c\left(\sum_{K \in \mathcal{\tau}_{h}}\left\|\Delta v_{h}\right\|_{0, K}^{2}\right)^{1 / 2}\left(\sum_{e \in \mathcal{E}} h_{e}^{-1}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2}  \tag{3.3}\\
& \leqslant \varepsilon \sum_{K \in \mathcal{\tau}_{h}}\left\|\Delta v_{h}\right\|_{0, K}^{2}+c \varepsilon^{-1} \sum_{e \in \mathcal{E}} h_{e}^{-1}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}
\end{align*}
$$

where $c>0$ is independent of $h$ and $\varepsilon>0$ is a sufficiently small constant. Thus

$$
\begin{align*}
a_{S}\left(v_{h}, v_{h}\right) \geqslant & \sum_{K \in \tau_{h}}\left\|\Delta v_{h}\right\|_{0, K}^{2}-2 \varepsilon \sum_{K \in \tau_{h}}\left\|\Delta v_{h}\right\|_{0, K}^{2}-c \varepsilon^{-1} \sum_{e \in \mathcal{E}} h_{e}^{-1}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}  \tag{3.4}\\
& +\sum_{e \in \mathcal{E}} \frac{\beta}{h_{e}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2} d s .
\end{align*}
$$

Taking $\varepsilon=1 / 4$, then, for sufficiently large $\beta$ such that $\beta>4 c$, using (2.25) we have

$$
\begin{equation*}
a_{S}\left(v_{h}, v_{h}\right) \geqslant \frac{1}{2} \sum_{K \in \mathcal{\tau}_{h}}\left\|\Delta v_{h}\right\|_{0, K}^{2}+(\beta-4 c) \sum_{e \in \mathcal{E}} h_{e}^{-1}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2} \geqslant \kappa_{1}\left\|v_{h}\right\|_{h}^{2} \tag{3.5}
\end{equation*}
$$

with $\mathcal{\kappa}_{1}=\min \{1 / 2, \beta-4 c\} \kappa_{0}$.
A direct result of Lemma 3.1 is that the discretized problem (2.17) admits a unique solution $u_{h}^{S} \in V_{h}$ for sufficiently large $\beta$.

Lemma 3.2. For all $\phi \in V$, there holds

$$
\begin{equation*}
a_{S}\left(\phi-P_{h} \phi, v_{h}\right) \leqslant c h^{\mu-2}\left\|v_{h}\right\|_{h}\||\phi|\|_{s^{\prime}} \quad \forall v_{h} \in V_{h} \tag{3.6}
\end{equation*}
$$

where $\mu=\min \{s, r+1\}$ and $c>0$ is independent of $h$.
Proof. For all $\phi \in V$ and $v_{h} \in V_{h}$, we have

$$
\begin{align*}
a_{S}\left(\phi-P_{h} \phi, v_{h}\right)= & \sum_{K \in \mathcal{\tau}_{h}} \int_{K} \Delta\left(\phi-P_{h} \phi\right) \Delta v_{h} d x-\sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta\left(\phi-P_{h} \phi\right)\right\}\left[\frac{\partial v_{h}}{\partial n}\right] d s \\
& -\sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta v_{h}\right\}\left[\frac{\partial\left(\phi-P_{h} \phi\right)}{\partial n}\right] d s+\sum_{e \in \mathcal{E}} \frac{\beta}{h_{e}} \int_{e}\left[\frac{\partial\left(\phi-P_{h} \phi\right)}{\partial n}\right]\left[\frac{\partial v_{h}}{\partial n}\right] d s \\
\leqslant & \left\|\phi-P_{h} \phi\right\|_{h}\left\|v_{h}\right\|_{h}+\left(\sum_{e \in \mathcal{E}}\left\|\left\{\Delta\left(\phi-P_{h} \phi\right)\right\}\right\|_{0, e}^{2}\right)^{1 / 2}\left(\sum_{e \in \mathcal{E}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2} \\
& +\left(\sum_{e \in \mathcal{\varepsilon}} h_{e}\left\|\left\{\Delta v_{h}\right\}\right\|_{0, e}^{2}\right)^{1 / 2}\left(\sum_{e \in \mathcal{\varepsilon}} h_{e}^{-1}\left\|\left[\frac{\partial\left(\phi-P_{h} \phi\right)}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2} \\
& +c\left(\sum_{e \in \mathcal{E}} h_{e}^{-2}\left\|\left[\frac{\partial\left(\phi-P_{h} \phi\right)}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2}\left(\sum_{e \in \mathcal{E}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2}, \tag{3.7}
\end{align*}
$$

where $c>0$ is independent of $h$. Note that $\int_{e}\left[\left(\partial v_{h}\right) / \partial n\right] d s=0$ for $e \in \mathcal{E}$. Thus, for some constant $C^{e}$ depending on $e$ there holds

$$
\begin{equation*}
\int_{e}\left|\left[\frac{\partial v_{h}}{\partial n}\right]\right|^{2} d s=\int_{e}\left[\frac{\partial v_{h}}{\partial n}\right]\left[\frac{\partial v_{h}}{\partial n}-C^{e}\right] d s \leqslant\left(\int_{e}\left|\left[\frac{\partial v_{h}}{\partial n}\right]\right|^{2} d s\right)^{1 / 2}\left(\int_{e}\left|\left[\frac{\partial v_{h}}{\partial n}-C^{e}\right]\right|^{2} d s\right)^{1 / 2} \tag{3.8}
\end{equation*}
$$

That is

$$
\begin{equation*}
\int_{e}\left|\left[\frac{\partial v_{h}}{\partial n}\right]\right|^{2} d s \leqslant \int_{e}\left|\left[\frac{\partial v_{h}}{\partial n}-C^{e}\right]\right|^{2} d s \tag{3.9}
\end{equation*}
$$

From [9], we have

$$
\begin{equation*}
\left\|\left[\frac{\partial v_{h}}{\partial n}-C^{e}\right]\right\|_{0, e} \leqslant c h^{1 / 2}\left|v_{h}\right|_{2, K} \tag{3.10}
\end{equation*}
$$

where $c>0$ is independent of $h$. Thus

$$
\begin{equation*}
\left(\sum_{e \in \mathcal{E}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2} \leqslant c h^{1 / 2}\left\|v_{h}\right\|_{h} . \tag{3.11}
\end{equation*}
$$

Substituting (3.11) into (3.7) and using (2.22)-(2.23) give

$$
\begin{align*}
a_{S}\left(\phi-P_{h} \phi, v_{h}\right) \leqslant & \left\|\phi-P_{h} \phi\right\|_{h}\left\|v_{h}\right\|_{h}+c h^{1 / 2}\left(\sum_{e \in \mathcal{E}}\left\|\left\{\Delta\left(\phi-P_{h} \phi\right)\right\}\right\|_{0, e}^{2}\right)^{1 / 2}\left\|v_{h}\right\|_{h} \\
& +c\left(\sum_{e \in \mathcal{E}} h_{e}^{-1}\left\|\left[\frac{\partial\left(\phi-P_{h} \phi\right)}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2}\left\|v_{h}\right\|_{h} \\
& +c h^{1 / 2}\left(\sum_{e \in \mathcal{E}} h_{e}^{-2}\left\|\left[\frac{\partial\left(\phi-P_{h} \phi\right)}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2}\left\|v_{h}\right\|_{h} \leqslant c h^{\mu-2}\||\phi|\|_{s}\left\|v_{h}\right\|_{h} . \tag{3.12}
\end{align*}
$$

Theorem 3.3. Suppose that $u \in V$ and $u_{h}^{S} \in V_{h}$ are the solutions to problems (SP) and (2.17), respectively; then the following optimal broken $H^{2}$ error estimate holds:

$$
\begin{equation*}
\left\|u-u_{h}^{S}\right\|_{h} \leqslant c h^{\mu-2}\||u|\|_{s^{\prime}} \tag{3.13}
\end{equation*}
$$

where $\mu=\min \{s, r+1\}$ and $c>0$ is independent of $h$.

Proof. According to Lemma 3.1, we have

$$
\begin{align*}
\kappa_{1}\left\|P_{h} u-u_{h}^{S}\right\|_{h}^{2} & \leqslant a_{S}\left(P_{h} u-u_{h}^{S}, P_{h} u-u_{h}^{S}\right)=a_{S}\left(P_{h} u-u, P_{h} u-u_{h}^{S}\right)+a_{S}\left(u-u_{h^{\prime}}^{S}, P_{h} u-u_{h}^{S}\right) \\
& =a_{S}\left(P_{h} u-u, P_{h} u-u_{h}^{S}\right) \\
& \leqslant c h^{\mu-2}\|\mid u\|_{s}\left\|P_{h} u-u_{h}^{S}\right\|_{h^{\prime}} \tag{3.14}
\end{align*}
$$

where we use the orthogonal equation (2.19) and Lemma 3.2. The previous estimate implies

$$
\begin{equation*}
\left\|P_{h} u-u_{h}^{S}\right\|_{h} \leqslant c h^{\mu-2}\| \| u \|_{S} \tag{3.15}
\end{equation*}
$$

Finally, the triangular inequality and (2.24) yield

$$
\begin{equation*}
\left\|u-u_{h}^{S}\right\|_{h} \leqslant\left\|u-P_{h} u\right\|_{h}+\left\|P_{h} u-u_{h}^{S}\right\|_{h} \leqslant c h^{\mu-2}\| \| u \|_{S} \tag{3.16}
\end{equation*}
$$

Next, we will show the optimal $L^{2}$ error estimate in terms of the duality technique. Suppose $g \in L^{2}(\Omega)$ and consider the following biharmonic problem:

$$
\begin{gather*}
\Delta^{2} w=g, \quad \text { in } \Omega \\
w=\frac{\partial w}{\partial n}=0, \quad \text { on } \partial \Omega \tag{3.17}
\end{gather*}
$$

Suppose that problem (3.17) admits a unique solution $w \in H_{0}^{2}(\Omega) \cap H^{4}(\Omega)$ such that

$$
\begin{equation*}
\|w\|_{4} \leqslant c\|g\| \tag{3.18}
\end{equation*}
$$

where $\|\cdot\|_{4}$ denotes the $H^{4}$ norm in $\Omega$ and $\|\cdot\|$ denotes the $L^{2}$ norm in $\Omega$ and $c>0$ is independent of $h$.

Denote by $\Pi_{h}$ the $C^{1}$ continuous interpolate operator from $V$ to $H_{0}^{2}(\Omega) \cap V_{h}$, and $\Pi_{h}$ satisfies the approximation property (2.22). Then for the solution $w \in H^{4}(\Omega) \cap H_{0}^{2}(\Omega)$ to problem (3.17), there hold

$$
\begin{gather*}
\left\|w-\Pi_{h} w\right\|_{2} \leqslant c h^{2}\|w\|_{4} \leqslant c h^{2}\|g\| \\
\left\|w-\Pi_{h} w\right\|_{2, e} \leqslant c h_{e}^{3 / 2}\|w\|_{4} \leqslant c h_{e}^{3 / 2}\|g\|, \quad \forall e \in \mathcal{\varepsilon} \tag{3.19}
\end{gather*}
$$

where $c>0$ is independent of $h$.

Theorem 3.4. Suppose that $u \in V$ and $u_{h}^{S} \in V_{h}$ are the solutions to problems (SP) and (2.17), respectively; then the following optimal $L^{2}$ error estimate holds:

$$
\begin{equation*}
\left\|u-u_{h}^{S}\right\| \leqslant c h^{\mu}\||u|\|_{s^{\prime}} \tag{3.20}
\end{equation*}
$$

where $\mu=\min \{s, r+1\}$ and $c>0$ is independent of $h$.
Proof. Setting $g=u-u_{h}^{S}$ in (3.17), multiplying (3.17) by $u-u_{h}^{S}$, and integrating over $\Omega$, we have

$$
\begin{align*}
\left\|u-u_{h}^{S}\right\|^{2} & =\int_{\Omega}\left(\Delta^{2} w\right)\left(u-u_{h}^{S}\right) d x=a_{S}\left(w, u-u_{h}^{S}\right)=a_{S}\left(w-\Pi_{h} w, u-u_{h}^{S}\right)  \tag{3.21}\\
& =a_{S}\left(w-\Pi_{h} w, P_{h} u-u_{h}^{S}\right)+a_{S}\left(w-\Pi_{h} w, u-P_{h} u\right)
\end{align*}
$$

where we use the orthogonal equation (2.19). We estimate two terms in the right-hand side of (3.21) as follows:

$$
\begin{align*}
a_{S}(w & \left.-\Pi_{h} w, P_{h} u-u_{h}^{S}\right) \\
= & \sum_{K \in \mathcal{Z}_{h}} \int_{K} \Delta\left(w-\Pi_{h} w\right) \Delta\left(P_{h} u-u_{h}^{S}\right) d x-\sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta\left(w-\Pi_{h} w\right)\right\}\left[\frac{\partial\left(P_{h} u-u_{h}^{S}\right)}{\partial n}\right] d s \\
& -\sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta\left(P_{h} u-u_{h}^{S}\right)\right\}\left[\frac{\partial\left(w-\Pi_{h} w\right)}{\partial n}\right] d s+\sum_{e \in \mathcal{E}} \frac{\beta}{h_{e}} \int_{e}\left[\frac{\partial\left(w-\Pi_{h} w\right)}{\partial n}\right]\left[\frac{\partial\left(P_{h} u-u_{h}^{S}\right)}{\partial n}\right] d s \\
= & \sum_{K \in \mathcal{Z}_{h}} \int_{K} \Delta\left(w-\Pi_{h} w\right) \Delta\left(P_{h} u-u_{h}^{S}\right) d x-\sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta\left(w-\Pi_{h} w\right)\right\}\left[\frac{\partial\left(P_{h} u-u_{h}^{S}\right)}{\partial n}\right] d s \\
\leqslant & \left\|w-\Pi_{h} w\right\|_{2}\left\|P_{h} u-u_{h}^{S}\right\|_{h}+\left(\sum_{e \in \mathcal{E}}\left\|w-\Pi_{h} w\right\|_{2, e}^{2}\right)^{1 / 2}\left(\sum_{e \in \mathcal{E}}\left\|P_{h} u-u_{h}^{S}\right\|_{1, e}^{2}\right)^{1 / 2} \\
\leqslant & c h^{2}\left\|P_{h} u-u_{h}^{S}\right\|_{h}\left\|u-u_{h}^{S}\right\| \leqslant c h^{\mu}\|u l\|_{s}\left\|u-u_{h}^{S}\right\| \tag{3.22}
\end{align*}
$$

where we use the estimate (3.15). In terms of the inequalities (2.22)-(2.23), we have

$$
\begin{aligned}
a_{S}(w- & \left.\Pi_{h} w, u-P_{h} u\right) \\
= & \sum_{K \in \tau_{h}} \int_{K} \Delta\left(w-\Pi_{h} w\right) \Delta\left(u-P_{h} u\right) d x-\sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta\left(w-\Pi_{h} w\right)\right\}\left[\frac{\partial\left(u-P_{h} u\right)}{\partial n}\right] d s \\
& -\sum_{e \in \mathcal{\varepsilon}} \int_{e}\left\{\Delta\left(u-P_{h} u\right)\right\}\left[\frac{\partial\left(w-\Pi_{h} w\right)}{\partial n}\right] d s+\sum_{e \in \mathcal{\varepsilon}} \frac{\beta}{h_{e}} \int_{e}\left[\frac{\partial\left(w-\Pi_{h} w\right)}{\partial n}\right]\left[\frac{\partial\left(u-P_{h} u\right)}{\partial n}\right] d s
\end{aligned}
$$

$$
\begin{align*}
& =\sum_{K \in \mathcal{I}_{h}} \int_{K} \Delta\left(w-\Pi_{h} w\right) \Delta\left(u-P_{h} u\right) d x-\sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta\left(w-\Pi_{h} w\right)\right\}\left[\frac{\partial\left(u-P_{h} u\right)}{\partial n}\right] d s \\
& \leqslant\left\|w-P_{h} w\right\|_{2}\left\|u-P_{h} u\right\|_{h}+\left(\sum_{e \in \mathcal{E}}\left\|w-\Pi_{h} w\right\|_{2, e}^{2}\right)^{1 / 2}\left(\sum_{e \in \mathcal{E}}\left\|\left[\frac{\partial\left(u-P_{h} u\right)}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2} \\
& \leqslant c h^{\mu}\|u\|_{s}\left\|u-u_{h}^{S}\right\| . \tag{3.23}
\end{align*}
$$

Substituting the estimates (3.22)-(3.23) into (3.21) yields

$$
\begin{equation*}
\left\|u-u_{h}^{S}\right\| \leqslant c h^{\mu}\|\mid u\|_{s} . \tag{3.24}
\end{equation*}
$$

## 4. $C^{0}$ Nonsymmetric Interior Penalty Method

In this section, we will show the error estimates in the broken $H^{2}$ norm and in $L^{2}$ norm between the solution $u$ to problem (2.1) and the solution $u_{h}^{\mathrm{NS}}$ to the problem (2.18). The optimal broken $H^{2}$ error estimate is derived. However, the $L^{2}$ error estimate is suboptimal because of the lack of adjoint consistency. According to Lemma 2.1, we have

$$
\begin{equation*}
a_{\mathrm{NS}}\left(v_{h}, v_{h}\right) \geqslant \kappa_{0}\left\|v_{h}\right\|_{h}^{2}, \quad \forall v_{h} \in V_{h} . \tag{4.1}
\end{equation*}
$$

Moreover, for the nonsymmetric bilinear form $a_{\mathrm{NS}}(\cdot, \cdot)$, proceeding as in the proof of Lemma 3.2, we have the following lemma.

Lemma 4.1. For all $\phi \in V$, there holds

$$
\begin{equation*}
a_{\mathrm{NS}}\left(\phi-P_{h} \phi, v_{h}\right) \leqslant c h^{\mu-2}\left\|v_{h}\right\|_{h}\||\phi|\|_{s^{\prime}} \quad \forall v_{h} \in V_{h}, \tag{4.2}
\end{equation*}
$$

where $\mu=\min \{s, r+1\}$ and $c>0$ is independent of $h$.
Theorem 4.2. Suppose that $u \in V$ and $u_{h}^{\text {NS }} \in V_{h}$ are the solutions to problems (NSP) and (2.18), respectively; then there holds

$$
\begin{equation*}
\left\|u-u_{h}^{\mathrm{NS}}\right\|_{h} \leqslant c h^{\mu-2}\|\mid u\|_{s^{\prime}} \tag{4.3}
\end{equation*}
$$

where $\mu=\min \{s, r+1\}$ and $c>0$ is independent of $h$.

Proof. According to (2.20), (4.1), and Lemma 4.1, we have

$$
\begin{align*}
\kappa_{0}\left\|P_{h} u-u_{h}^{\mathrm{NS}}\right\|_{h}^{2} & \leqslant a_{\mathrm{NS}}\left(\mathrm{P}_{h} u-u_{h}^{\mathrm{NS}}, P_{h} u-u_{h}^{\mathrm{NS}}\right) \\
& =a_{\mathrm{NS}}\left(P_{h} u-u, P_{h} u-u_{h}^{\mathrm{NS}}\right)+a_{\mathrm{NS}}\left(u-u_{h}^{\mathrm{NS}}, P_{h} u-u_{h}^{\mathrm{NS}}\right)  \tag{4.4}\\
& =a_{\mathrm{NS}}\left(P_{h} u-u, P_{h} u-u_{h}^{\mathrm{NS}}\right) \leqslant c h^{\mu-2}\|\mid u\|\left\|_{S}\right\| P_{h} u-u_{h}^{\mathrm{NS}} \|_{h^{\prime}}
\end{align*}
$$

where $c>0$ is independent of $h$. That is

$$
\begin{equation*}
\left\|P_{h} u-u_{h}^{\mathrm{NS}}\right\|_{h} \leqslant c h^{\mu-2}\||u|\|_{s} . \tag{4.5}
\end{equation*}
$$

Using the triangular inequality yields

$$
\begin{equation*}
\left\|u-u_{h}^{\mathrm{NS}}\right\|_{h} \leqslant\left\|u-P_{h} u\right\|_{h}+\left\|P_{h} u-u_{h}^{\mathrm{NS}}\right\|_{h} \leqslant c h^{\mu-2}\| \| u \mid \|_{S} \tag{4.6}
\end{equation*}
$$

Theorem 4.3. Suppose that $u \in V$ and $u_{h}^{N S} \in V_{h}$ are the solutions to problems (NSP) and (2.18), respectively; then there holds

$$
\begin{equation*}
\left\|u-u_{h}^{\mathrm{NS}}\right\| \leqslant c h^{\mu-1}\||u|\|_{s^{\prime}} \tag{4.7}
\end{equation*}
$$

where $\mu=\min \{s, r+1\}$ and $c>0$ is independent of $h$.
Proof. Setting $g=u-u_{h}^{\mathrm{NS}}$ in (3.17), multiplying (3.17) by $u-u_{h}^{\mathrm{NS}}$, and integrating over $\Omega$, we have

$$
\begin{aligned}
\left\|u-u_{h}^{\mathrm{NS}}\right\|^{2}= & \int_{\Omega}\left(\Delta^{2} w\right)\left(u-u_{h}^{\mathrm{NS}}\right) d x=a_{\mathrm{NS}}\left(w, u-u_{h}^{\mathrm{NS}}\right) \\
= & a_{\mathrm{NS}}\left(w-\Pi_{h} w, u-u_{h}^{\mathrm{NS}}\right)+2 \sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta \Pi_{h} w\right\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{NS}}\right)}{\partial n}\right] d s \\
& -2 \sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta\left(u-u_{h}^{\mathrm{NS}}\right)\right\}\left[\frac{\partial \Pi_{h} w}{\partial n}\right] d s \\
= & a_{\mathrm{NS}}\left(w-\Pi_{h} w, u-u_{h}^{\mathrm{NS}}\right)+2 \sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta \Pi_{h} w\right\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{NS}}\right)}{\partial n}\right] d s
\end{aligned}
$$

$$
\begin{align*}
= & \sum_{K \in \mathcal{\tau}_{h}} \int_{K} \Delta\left(w-\Pi_{h} w\right) \Delta\left(u-u_{h}^{\mathrm{NS}}\right) d x \\
& +3 \sum_{e \in \mathfrak{E}} \int_{e}\left\{\Delta\left(\Pi_{h} w-w\right)\right\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{NS}}\right)}{\partial n}\right] d s+2 \sum_{e \in \mathcal{E}} \int_{e}\{\Delta w\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{NS}}\right)}{\partial n}\right] d s \\
= & I_{1}+I_{2}+I_{3} \tag{4.8}
\end{align*}
$$

where we use $\Pi_{h} w \in H_{0}^{2}(\Omega)$. We estimate $I_{1}$ as follows:

$$
\begin{align*}
I_{1} & =\sum_{K \in \tau_{h}} \int_{K} \Delta\left(w-\Pi_{h} w\right) \Delta\left(u-u_{h}^{\mathrm{NS}}\right) d x \leqslant\left\|u-u_{h}^{\mathrm{NS}}\right\|_{h}\left\|w-\Pi_{h} w\right\|_{2}  \tag{4.9}\\
& \leqslant c h^{\mu}\||u|\|_{s}\left\|u-u_{h}^{\mathrm{NS}}\right\|
\end{align*}
$$

We estimate $I_{2}$ as follows:

$$
\begin{align*}
I_{2}= & 3 \sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta\left(\Pi_{h} w-w\right)\right\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{NS}}\right)}{\partial n}\right] d s \\
= & 3 \sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta\left(\Pi_{h} w-w\right)\right\}\left[\frac{\partial\left(u-P_{h} u\right)}{\partial n}\right] d s+3 \sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta\left(\Pi_{h} w-w\right)\right\}\left[\frac{\partial\left(P_{h} u-u_{h}^{\mathrm{NS}}\right)}{\partial n}\right] d s \\
\leqslant & c\left(\sum_{e \in \mathcal{E}}\left\|w-\Pi_{h} w\right\|_{2, e}^{2}\right)^{1 / 2}\left(\sum_{e \in \mathcal{E}}\left\|\left[\frac{\partial\left(u-P_{h} u\right)}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2} \\
& +\left(\sum_{e \in \mathcal{E}}\left\|w-\Pi_{h} w\right\|_{2, e}^{2}\right)^{1 / 2}\left(\sum_{e \in \mathcal{E}} \|\left[\frac{\partial\left(P_{h} u-u_{h}^{\mathrm{NS}}\right)}{\partial n} \|^{2}\right)_{0, e}^{1 / 2}\right. \\
\leqslant & c h^{3 / 2}\left(h^{\mu-3 / 2}\| \| u\left\|_{s}+c h^{1 / 2}\right\| P_{h} u-u_{h}^{\mathrm{NS}} \|_{h}\right)\left\|u-u_{h}^{\mathrm{NS}}\right\| \leqslant c h^{\mu}\| \| u\left\|_{s}\right\| u-u_{h}^{\mathrm{NS}} \| \tag{4.10}
\end{align*}
$$

We estimate $I_{3}$ as follows:

$$
\begin{aligned}
I_{3} & =2 \sum_{e \in \mathcal{E}} \int_{e}\{\Delta w\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{NS}}\right)}{\partial n}\right] d s=2 \sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta w-C^{e}\right\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{NS}}\right)}{\partial n}\right] d s \\
& =2 \sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta w-C^{e}\right\}\left[\frac{\partial\left(\mathrm{u}-P_{h} u\right)}{\partial n}\right] d s+2 \sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta w-C^{e}\right\}\left[\frac{\partial\left(P_{h} u-u_{h}^{\mathrm{NS}}\right)}{\partial n}\right] d s
\end{aligned}
$$

$$
\begin{align*}
\leqslant & c\left(\sum_{e \in \mathcal{E}}\left\|\Delta w-C^{e}\right\|_{0, e}^{2}\right)^{1 / 2}\left(\sum_{e \in \mathcal{E}}\left\|u-P_{h} u\right\|_{1, e}^{2}\right)^{1 / 2} \\
& +\left(\sum_{e \in \mathcal{E}}\left\|\Delta w-C^{e}\right\|_{0, e}^{2}\right)^{1 / 2}\left(\sum_{e \in \mathcal{E}}\left\|\left[\frac{\partial\left(P_{h} u-u_{h}^{\mathrm{NS}}\right)}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2} \\
\leqslant & c\left(\sum_{e \in \mathcal{E}}\left\|\Delta w-C^{e}\right\|_{0, e}^{2}\right)^{1 / 2}\left(h^{\mu-3 / 2}\|\mid u\|_{s}+c h^{1 / 2}\left\|P_{h} u-u_{h}^{\mathrm{NS}}\right\|_{h}\right) \\
\leqslant & c h^{\mu-3 / 2}\left(\sum_{e \in \mathcal{E}}\left\|\Delta w-C^{e}\right\|_{0, e}^{2}\right)^{1 / 2}\| \| u \|_{s} \tag{4.11}
\end{align*}
$$

where $C^{e}$ is some positive constant. Substituting the following estimate

$$
\begin{equation*}
\sum_{e \in \mathcal{E}}\left\|\Delta w-C^{e}\right\|_{0, e}^{2} \leqslant c \sum_{e \in \mathcal{E}} h_{e}\|w\|_{3, e}^{2} \leqslant c h\|w\|_{4}^{2} \leqslant c h\left\|u-u_{h}^{\mathrm{NS}}\right\|^{2} \tag{4.12}
\end{equation*}
$$

into (4.11) gives

$$
\begin{equation*}
I_{3} \leqslant c h^{\mu-1}\||u|\|_{s}\left\|u-u_{h}^{\mathrm{NS}}\right\| \tag{4.13}
\end{equation*}
$$

Finally, substituting (4.9)-(4.13) into (4.8), we obtain

$$
\begin{equation*}
\left\|u-u_{h}^{\mathrm{NS}}\right\| \leqslant c h^{\mu-1}\| \| u \mid \|_{S} . \tag{4.14}
\end{equation*}
$$

## 5. $C^{0}$ Superpenalty Nonsymmetric Method

In order to obtain the optimal $L^{2}$ error estimate for the nonsymmetric method, in this section we will consider the $C^{0}$ superpenalty nonsymmetric method. First, we introduce a new nonsymmetric bilinear form:

$$
\begin{align*}
a_{\mathrm{SNS}}(\mathrm{u}, v)= & \sum_{K \in \mathcal{\tau}_{h}} \int_{K} \Delta u \Delta v d x-\sum_{e \in \mathcal{E}} \int_{e}\{\Delta u\}\left[\frac{\partial v}{\partial n}\right] d s+\sum_{e \in \mathcal{E}} \int_{e}\{\Delta v\}\left[\frac{\partial u}{\partial n}\right] d s \\
& +\sum_{e \in \mathcal{E}} \frac{1}{h_{e}^{3}} \int_{e}\left[\frac{\partial u}{\partial n}\right]\left[\frac{\partial v}{\partial n}\right] . \tag{5.1}
\end{align*}
$$

The broken $H^{2}$ norm is modified to

$$
\begin{equation*}
\left\|\left|v_{h}\right|\right\|_{h}=\left(\sum_{K \in \mathcal{Z}_{h}}\left|v_{h}\right|_{2, K_{i}}^{2}+\sum_{e \in \mathcal{E}} \frac{1}{h_{e}^{3}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2}, \quad \forall v_{h} \in V_{h} \tag{5.2}
\end{equation*}
$$

From Lemma 2.1, it is easy to show that there exists some constant $\kappa_{2}>0$ such that

$$
\begin{equation*}
a_{\mathrm{SNS}}\left(v_{h}, v_{h}\right) \geqslant \kappa_{2}\| \| v_{h} \mid \|_{h}^{2}, \quad \forall v_{h} \in V_{h} \tag{5.3}
\end{equation*}
$$

In fact, we have

$$
\begin{align*}
a_{\mathrm{SNS}}\left(v_{h}, v_{h}\right) & =\sum_{K \in \mathcal{\tau}_{h}}\left\|\Delta v_{h}\right\|_{0, K}^{2}+\sum_{e \in \mathcal{E}} \frac{1}{h_{e}^{3}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2} \\
& \geqslant \frac{1}{2} \sum_{e \in \mathcal{E}} \frac{1}{h_{e}^{3}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}+\frac{1}{2}\left(\sum_{K \in \tau_{h}}\left\|\Delta v_{h}\right\|_{0, K}^{2}+\sum_{e \in \mathcal{E}} \frac{1}{h_{e}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}\right)  \tag{5.4}\\
& \geqslant \frac{1}{2} \sum_{e \in \varepsilon} \frac{1}{h_{e}^{3}}\left\|\left[\frac{\partial v_{h}}{\partial n}\right]\right\|_{0, e}^{2}+\frac{\kappa_{0}}{2} \sum_{K \in \tau_{h}}\left|v_{h}\right|_{2, K}^{2} \geqslant \kappa_{2}\left\|\left|v_{h}\right|\right\|_{h}^{2}
\end{align*}
$$

for $\kappa_{2}=\min \left\{1 / 2, \kappa_{0} / 2\right\}$. Since the solution $u$ to problem (2.1) belongs to $H^{4}(\Omega) \cap H_{0}^{2}(\Omega)$, then it satisfies

$$
\begin{equation*}
a_{\mathrm{SNS}}(u, v)=(f, v), \quad \forall v \in V \tag{SNSP}
\end{equation*}
$$

In this case, the the $C^{0}$ superpenalty nonsymmetric finite element approximation of (2.1) is

$$
\begin{align*}
& \text { find } u_{h}^{\mathrm{SNS}} \in V_{h} \text { such that, } \\
& a_{\mathrm{SNS}}\left(u_{h}^{\mathrm{SNS}}, v_{h}\right)=\left(f, v_{h}\right), \quad \forall v_{h} \in V_{h} \tag{5.5}
\end{align*}
$$

Then, we have the following orthogonal equation:

$$
\begin{equation*}
a_{\mathrm{SNS}}\left(u_{h}^{\mathrm{SNS}}-u, v_{h}\right)=0, \quad \forall v_{h} \in V_{h} \tag{5.6}
\end{equation*}
$$

Let $\Pi_{h}$ be the $C^{1}$ continuous interpolated operator defined in Section 3. Observe that $\left[\left(\partial u-\Pi_{h} u\right) / \partial n\right]=0$ on every $e \in \varepsilon$, and proceeding as in the proof of Lemmas 3.2 and 4.1, we have the following.

Lemma 5.1. For all $\phi \in V$, there holds

$$
\begin{equation*}
a_{N S}\left(\phi-\Pi_{h} \phi, v_{h}\right) \leqslant c h^{\mu-2}\left\|\left|v_{h}\right|\right\|_{h}\||\phi|\|_{s^{\prime}} \quad \forall v_{h} \in V_{h} \tag{5.7}
\end{equation*}
$$

where $\mu=\min \{s, r+1\}$ and $c>0$ is independent of $h$.

Using Lemma 5.1, the following optimal broken $H^{2}$ error estimate holds.
Theorem 5.2. Suppose that $u \in V$ and $u_{h}^{\mathrm{SNS}} \in V_{h}$ are the solutions to problems (SNSP) and (5.5), respectively; then there holds

$$
\begin{equation*}
\left\|\left|u-u_{h}^{\mathrm{SNS}}\right|\right\|_{h} \leqslant c h^{\mu-2}\||u|\|_{s^{\prime}} \tag{5.8}
\end{equation*}
$$

where $\mu=\min \{s, r+1\}$ and $c>0$ is independent of $h$.
The main result in this section is the following optimal $L^{2}$ error estimate.
Theorem 5.3. Suppose that $u \in V$ and $u_{h}^{\text {SNS }} \in V_{h}$ are the solutions to problems (SNSP) and (5.5), respectively; then there holds

$$
\begin{equation*}
\left\|u-u_{h}^{\mathrm{SNS}}\right\| \leqslant c h^{\mu}\||u|\|_{s^{\prime}} \tag{5.9}
\end{equation*}
$$

where $\mu=\min \{s, r+1\}$ and $c>0$ is independent of $h$.
Proof. Let $g=u-u_{h}^{\mathrm{SNS}}$. Multiplying (3.17) by $u-u_{h}^{\mathrm{SNS}}$ and integrating over $\Omega$, we have

$$
\begin{align*}
\left\|u-u_{h}^{\mathrm{SNS}}\right\|^{2}= & \int_{\Omega}\left(\Delta^{2} w\right)\left(u-u_{h}^{\mathrm{SNS}}\right) d x=a_{\mathrm{SNS}}\left(w, u-u_{h}^{\mathrm{SNS}}\right) \\
= & a_{\mathrm{SNS}}\left(w-\Pi_{h} w, u-u_{h}^{\mathrm{SNS}}\right)+2 \sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta \Pi_{h} w\right\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{SNS}}\right)}{\partial n}\right] d s \\
& \left.-2 \sum_{e \in \mathcal{\varepsilon}} \int_{e}\left\{\Delta\left(u-u_{h}^{\mathrm{SNS}}\right)\right\}\left[\frac{\partial \Pi_{h} w}{\partial n}\right]^{2}\right] d s \\
= & a_{\mathrm{SNS}}\left(w-\Pi_{h} w, u-u_{h}^{\mathrm{SNS}}\right)+2 \sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta \Pi_{h} w\right\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{SNS}}\right)}{\partial n}\right] d s \\
= & \sum_{K \in \tau_{h}} \int_{K} \Delta\left(w-\Pi_{h} w\right) \Delta\left(u-u_{h}^{\mathrm{SNS}}\right) d x \\
& +3 \sum_{e \in \mathcal{\varepsilon}} \int_{e}\left\{\Delta\left(\Pi_{h} w-w\right)\right\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{SNS}}\right)}{\partial n}\right] d s+2 \sum_{e \in \mathcal{E}} \int_{e}\{\Delta w\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{SNS}}\right)}{\partial n}\right] d s \\
= & I_{4}+I_{5}+I_{6}, \tag{5.10}
\end{align*}
$$

where we use $\Pi_{h} w \in C^{1}(\Omega) \cap H_{0}^{2}(\Omega)$. Proceeding as in the proof of Theorem 4.2, we can estimate $I_{4}$ and $I_{5}$ as follows:

$$
\begin{gather*}
I_{4}=\sum_{K \in \mathcal{Z}_{h}} \int_{K} \Delta\left(w-\Pi_{h} w\right) \Delta\left(u-u_{h}^{\mathrm{SNS}}\right) d x \leqslant c h^{\mu}\||u|\|_{s}\left\|u-u_{h}^{\mathrm{SNS}}\right\|  \tag{5.11}\\
I_{5}=3 \sum_{e \in \mathcal{E}} \int_{e}\left\{\Delta\left(\Pi_{h} w-w\right)\right\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{SNS}}\right)}{\partial n}\right] d s \leqslant c h^{\mu}\||u|\|_{s}\left\|u-u_{h}^{\mathrm{SNS}}\right\| . \tag{5.12}
\end{gather*}
$$

The different estimate compared to Theorem 4.3 is the estimate of $I_{6}$. Under the new norm ||| • |||h, we have

$$
\begin{align*}
I_{6} & =2 \sum_{e \in \mathcal{E}} \int_{e}\{\Delta w\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{SNS}}\right)}{\partial n}\right] d s=2 \sum_{e \in \mathcal{E}} \int_{e}\{\Delta w-c\}\left[\frac{\partial\left(u-u_{h}^{\mathrm{SNS}}\right)}{\partial n}\right] d s \\
& \leqslant c\left(\sum_{e \in \mathcal{E}} h_{e}^{3}\|\Delta w-C\|_{0, e}^{2}\right)^{1 / 2}\left(\sum_{e \in \mathcal{E}} \frac{1}{h_{e}^{3}}\left\|\left[\frac{\partial\left(u-u_{h}^{\mathrm{SNS}}\right)}{\partial n}\right]\right\|_{0, e}^{2}\right)^{1 / 2}  \tag{5.13}\\
& \leqslant c h^{2}\left\|u-u_{h}^{\mathrm{SNS}}\right\| \cdot\left\|\left|u-u_{h}^{\mathrm{SNS}}\| \|_{h} \leqslant c h^{\mu}\||u|\|_{s}\left\|u-u_{h}^{\mathrm{SNS}}\right\| .\right.\right.
\end{align*}
$$

Substituting (5.11)-(5.13) into (5.10) yields the following optimal $L^{2}$ error estimate:

$$
\begin{equation*}
\left\|u-u_{h}^{\mathrm{SNS}}\right\| \leqslant c h^{\mu}\||u|\|_{s} . \tag{5.14}
\end{equation*}
$$

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## References

[1] D. N. Arnold, F. Brezzi, B. Cockburn, and L. D. Marini, "Unified analysis of discontinuous Galerkin methods for elliptic problems," SIAM Journal on Numerical Analysis, vol. 39, no. 5, pp. 1749-1779, 2002.
[2] G. A. Baker, "Finite element methods for elliptic equations using nonconforming elements," Mathematics of Computation, vol. 31, no. 137, pp. 45-59, 1977.
[3] I. Mozolevski and E. Süli, "A priori error analysis for the hp-version of the discontinuous Galerkin finite element method for the biharmonic equation," Computational Methods in Applied Mathematics, vol. 3, no. 4, pp. 596-607, 2003.
[4] I. Mozolevski, E. Süli, and P. R. Bösing, "hp-version a priori error analysis of interior penalty discontinuous Galerkin finite element approximations to the biharmonic equation," Journal of Scientific Computing, vol. 30, no. 3, pp. 465-491, 2007.
[5] E. Süli and I. Mozolevski, "hp-version interior penalty DGFEMs for the biharmonic equation," Computer Methods in Applied Mechanics and Engineering, vol. 196, no. 13-16, pp. 1851-1863, 2007.
[6] I. Babuška and M. Zlámal, "Nonconforming elements in the finite element method with penalty," SIAM Journal on Numerical Analysis, vol. 10, pp. 863-875, 1973.
[7] G. Engel, K. Garikipati, T. J. R. Hughes, M. G. Larson, L. Mazzei, and R. L. Taylor, "Continuous/discontinuous finite element approximations of fourth-order elliptic problems in structural and continuum mechanics with applications to thin beams and plates, and strain gradient elasticity," Computer Methods in Applied Mechanics and Engineering, vol. 191, no. 34, pp. 3669-3750, 2002.
[8] S. C. Brenner and L.-Y. Sung, " $C^{0}$ interior penalty methods for fourth order elliptic boundary value problems on polygonal domains," Journal of Scientific Computing, vol. 22-23, pp. 83-118, 2005.
[9] B. Rivière, M. F. Wheeler, and V. Girault, "Improved energy estimates for interior penalty, constrained and discontinuous Galerkin methods for elliptic problems. I," Computational Geosciences, vol. 3, no. 34, pp. 337-360, 1999.
[10] P. G. Ciarlet, The Finite Element Method for Elliptic Problems, North-Holland Publishing, Amsterdam, The Netherlands, 1978, Studies in Mathematics and Its Applications.
[11] J. Guzmán and M. Neilan, "Conforming and divergence free stokes elements on general triangularmeshes," Mathematics of Computation. In press.

