

YANG-BAXTER MAPS: DYNAMICAL POINT OF VIEW

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Abstract A review of some recent results on the dynamical theory of the Yang-Baxter maps (also known as set-theoretical solutions to the quantum Yang-Baxter equation) is given. The central question is the integrability of the transfer dynamics. The relations with matrix factorisations, matrix KdV solitons, Poisson Lie groups, geometric crystals and tropical combinatorics are discussed and demonstrated on several concrete examples.

INTRODUCTION

The story of Yang-Baxter maps is usually traced back to Drinfeld [1], who suggested to study the set-theoretical solutions to the quantum Yang-Baxter equation, but the first interesting examples were constructed already by Sklyanin [2]. Drinfeld's question was immediately appreciated by Weinstein and Xu [3] who developed an approach to this problem based on the theory of Poisson Lie groups and symplectic groupoids. Etingof, Schedler and Soloviev [4, 5], partly inspired by Hietarinta's work [6], took more algebraic point of view (see also [7, 8, 9, 10, 11, 12]).

Very interesting examples of such solutions have appeared in the theory of geometric crystals (Berenstein-Kazhdan [13] and Etingof [14]). This opened remarkable links with the theory of soliton cellular automata ("box-ball" systems) and tropical combinatorics extensively developed by the Japanese school [15, 16, 17, 18, 19, 20, 21, 22, 23].

The study of the dynamical aspects of this problem was initiated in the paper [24], where the term "Yang-Baxter map" was introduced. My motivation came from the soliton theory, in particular from the theory of matrix KdV equation [25]. I was puzzled by what kind of information one can learn from the interaction of two matrix solitons and what is the role of the Yang-Baxter relation here. Analogy with the theory of solvable models in statistical mechanics and quantum inverse scattering method [28, 29, 31, 30, 32], where the quantum Yang-Baxter equation has its origin, naturally led to the notion of the transfer dynamics and to the relations with the theory of integrable systems. The main idea of [24] was that the investigation of the transfer dynamics should reveal the structures encapsulated in a given Yang-Baxter map. A very good example is the Adler map (see below) from the theory of the dressing chain [33, 34]. The corresponding transfer dynamics is nothing else but the discrete version of the finite-gap KdV dynamics

(Novikov equations [35, 36]) with its rich algebro-geometric and symplectic structures.

In this extended version of the lectures given at Kyoto meeting "Combinatorial Aspects of Integrable Systems" in July 2004 I am going to review the dynamical theory of the Yang-Baxter maps and illustrate it on several concrete examples following mainly to [24, 25, 26, 27].

We start with the general discussion of the Yang-Baxter maps and related transfer dynamics. The main property of the transfer maps, which they share with the transfer matrices, is the pairwise commutativity. If the maps are polynomial or rational the commutativity relation is quite strong and sometimes implies the solvability of the corresponding dynamical system (see [37] for the discussion of some known results in this direction). Although it is not clear what is the transfer dynamics in general, our main motive is that it should be "integrable" and one can think of it also as a set-theoretical analogue of the finite-gap theory.

We illustrate this by two constructions of the Yang-Baxter equation coming from the theory of integrable systems. The first one is related to matrix factorisations and QR-type procedure, which is a discrete analogue of the Lax representation. We explain how the Lax matrices can be read off the map itself following the procedure suggested in [26].

The second class is given by the interaction of solitons with non-trivial internal parameters (e.g. matrix KdV solitons). It is a well-known phenomenon in soliton theory (see e.g. [38, 39, 40, 41, 42]) that the interaction of n solitons is completely determined by their pairwise interactions (so there are no multiparticle effects). The fact that the final result is independent of the order of collisions means that the map determining the interaction of two solitons satisfies the Yang-Baxter relation.

We briefly consider also the examples of the Yang-Baxter maps related to geometric crystals and tropical combinatorics [14, 19, 22]. Then we discuss the Hamiltonian theory of the Yang-Baxter maps following recent work by Reshetikhin and the author [27] partly based on some general considerations from [3]. The main observation is that if the Lax matrices form symplectic leaves in some Poisson Lie groups, then the corresponding transfer dynamics is symplectic. In the last section we discuss the classification problem of the Yang-Baxter maps and its relation with the classification of the quadrirational maps done recently by Adler, Bobenko and Suris [43].

YANG-BAXTER MAPS AND TRANSFER DYNAMICS

The original *quantum Yang-Baxter equation* is the following relation on a linear operator $R : V \otimes V \rightarrow V \otimes V$:

$$(1) \quad R_{12}R_{13}R_{23} = R_{23}R_{13}R_{12},$$

where R_{ij} is acting in i -th and j -th components of the tensor product $V \otimes V \otimes V$ (see e.g. [32]).

Following Drinfeld's suggestion [1] consider the following set-theoretical version of this equation.

Let X be any set, $R : X \times X \rightarrow X \times X$ be a map from its square into itself. Let $R_{ij} : X^n \rightarrow X^n$, $X^n = X \times X \times \dots \times X$ be the maps which acts as R on i -th and j -th factors and identically on the others. More precisely, if $R(x, y) = (f(x, y), g(x, y)), x, y \in X$ then $R_{ij}(x_1, x_2, \dots, x_n) = (x_1, x_2, \dots, x_{i-1}, f(x_i, x_j), x_{i+1}, \dots, x_{j-1}, g(x_i, x_j), x_{j+1}, \dots, x_n)$ for $i < j$ and $(x_1, x_2, \dots, x_{j-1}, g(x_i, x_j), x_{j+1}, \dots, x_{i-1}, f(x_i, x_j), x_{i+1}, \dots, x_n)$ otherwise. In particular for $n = 2$ $R_{21}(x, y) = (g(y, x), f(y, x))$. If $P : X^2 \rightarrow X^2$ is the permutation of x and y : $P(x, y) = (y, x)$, then obviously we have

$$R_{21} = PRP.$$

Following [24] we will call R *Yang-Baxter map* if it satisfies the Yang-Baxter relation (1) considered as the equality of the maps of $X \times X \times X$ into itself. If additionally R satisfies the relation

$$(2) \quad R_{21}R = Id,$$

we will call it *reversible Yang-Baxter map*.

The condition (2) is usually called the *unitarity* condition, but in our case the term *reversibility* is more appropriate since in dynamical systems terminology this condition means that the map R is reversible with respect to the permutation P .

Note that if we consider the linear space $V = \mathbf{C}^X$ spanned by the set X , then any Yang-Baxter map R induces a linear operator in $V \otimes V$ which satisfies the quantum Yang-Baxter equation in the usual sense. Therefore we indeed have a very special class of solutions to this equation. If X is a finite set and R is a bijection then we have permutation-type solutions discussed in [6]. However this point of view on the Yang-Baxter maps seems to be artificial because it does not reflect the nature of the maps (like birationality).

Remark. We should warn the reader that some authors (see e.g. [7, 14, 12]) consider instead of the map R a closely related map $S = PR$, where P is the permutation map. The corresponding maps satisfies instead of the Yang-Baxter relation the *braid relation*

$$(3) \quad S_{12}S_{23}S_{12} = S_{23}S_{12}S_{23}$$

and instead of reversibility the *involutivity* condition

$$(4) \quad S^2 = Id$$

One should be careful here since the groups generated by the R_{ij} and S_{ij} maps could be different. For example, the group generated by R in general is \mathbf{Z} while for S it is \mathbf{Z}_2 . This is an important difference from dynamical point of view (see below).

We can represent the relations (1, 2) in the standard diagrammatic way as shown on the Fig. 1 and 2.

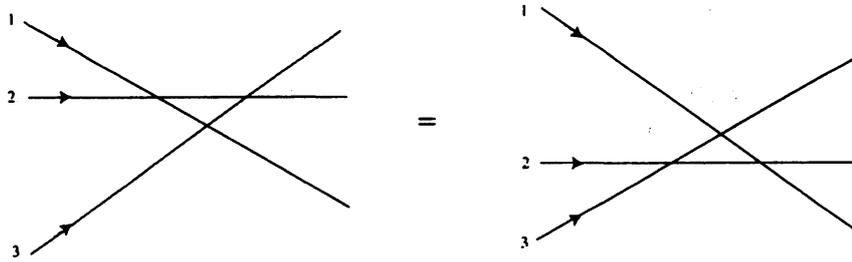


FIGURE 1. Yang-Baxter relation

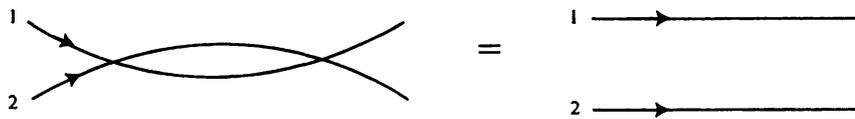


FIGURE 2. Reversibility

One should mention that there exist another (dual) way to represent the Yang-Baxter relation, which shows the relation with the so-called *3D consistency condition* in the theory of integrable discrete equations on quad-graphs [45, 46, 44]. In this representation the fields (elements of X) are assigned to the edges of elementary quadrilaterals rather than to the vertices. The corresponding "cubic" representation of the Yang-Baxter relation is illustrated in the Fig. 3.

The left-hand side of (1) corresponds to the chain of maps along the three rear faces of the cube on Fig. 3:

$$R_{12} : (x, y) \mapsto (x_2, y_1), \quad R_{13} : (x_2, z) \mapsto (x_{23}, z_1), \quad R_{23} : (y_1, z_1) \mapsto (y_{13}, z_{12}),$$

while its right-hand side corresponds to the chain of the maps along the three front faces of the cube:

$$R_{23} : (y, z) \mapsto (y_3, z_2), \quad R_{13} : (x, z_2) \mapsto (x_3, z_{12}), \quad R_{12} : (x_3, y_3) \mapsto (x_{23}, y_{13}).$$

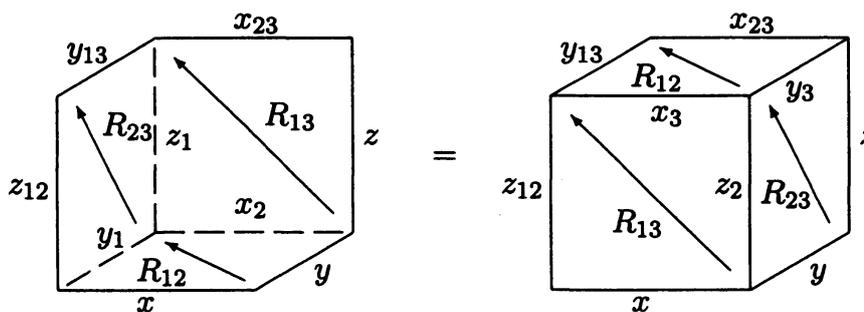


FIGURE 3. "Cubic" representation of the Yang-Baxter relation

So, (1) assures that two ways of obtaining (x_{23}, y_{13}, z_{12}) from the initial data (x, y, z) lead to the same results.

One can introduce also a more general parameter-dependent Yang-Baxter equation as the relation

$$(5) \quad R_{12}(\lambda_1, \lambda_2)R_{13}(\lambda_1, \lambda_3)R_{23}(\lambda_2, \lambda_3) = R_{23}(\lambda_2, \lambda_3)R_{13}(\lambda_1, \lambda_3)R_{12}(\lambda_1, \lambda_2)$$

and the corresponding unitarity (reversibility) condition as

$$(6) \quad R_{21}(\mu, \lambda)R(\lambda, \mu) = Id.$$

Although it can be reduced to the usual case by considering $\tilde{X} = X \times \mathbb{C}$ and $\tilde{R}(x, \lambda; y, \mu) = R(\lambda, \mu)(x, y)$ often it is useful to keep the parameter separately (see the examples below).

In the theory of the quantum Yang-Baxter equation the following *transfer matrices* $t^{(n)} : V^{\otimes n} \rightarrow V^{\otimes n}$ play a fundamental role. They are defined as the trace of the monodromy matrix

$$(7) \quad t^{(n)} = \text{tr}_{V_0} R_{0n} R_{0n-1} \dots R_{01}$$

with respect to the additional space V_0 . If the solution of the quantum Yang-Baxter equation depends on an additional spectral parameter λ then the corresponding transfer matrices commute:

$$(8) \quad t^{(n)}(\lambda)t^{(n)}(\mu) = t^{(n)}(\mu)t^{(n)}(\lambda),$$

which is the crucial fact in that theory (see [29, 30, 32]).

In the general set-theoretical case we have a problem with the trace operation (see although the remark at the end of this section), so one should look for an alternative. A possible candidate is given by the operators first considered by C.N Yang in the classical paper [28]. Their relation with transfer matrices are nicely explained in Gaudin's book [31] (see especially Chapter 10, sections 2 and 3). A related constructions were used by Frenkel and Reshetikhin in the theory of q-KZ equation [47] and by Fomin and Kirillov in the theory of Schubert polynomials [48] (see also [49, 50]).

Following [24] we introduce the *transfer maps* $T_i^{(n)}, i = 1, \dots, n$ as the maps of X^n into itself defined by the following formulas:

$$(9) \quad T_i^{(n)} = R_{ii+n-1}R_{ii+n-2} \dots R_{ii+1},$$

where the indices are considered modulo n with the agreement that we are using n rather than 0. In particular $T_1^{(n)} = R_{1n}R_{1n-1} \dots R_{12}$.

Theorem [24]. *For any reversible Yang-Baxter map R the transfer maps $T_i^{(n)}, i = 1, \dots, n$ commute with each other:*

$$(10) \quad T_i^{(n)}T_j^{(n)} = T_j^{(n)}T_i^{(n)}$$

and satisfy the property

$$(11) \quad T_1^{(n)}T_2^{(n)} \dots T_n^{(n)} = Id.$$

Conversly, suppose that the maps $T_i^{(n)}$ determined by the formula (9) commute and satisfy the relation (11) for any $n \geq 2$ then R is a reversible Yang-Baxter map.

Proof of the first part is very similar to the proof of the commutativity of the transfer matrices and follows from the consideration of the corresponding diagrams representing the products $T_i^{(n)}T_j^{(n)}$ and $T_j^{(n)}T_i^{(n)}$ respectively (cf [32]):

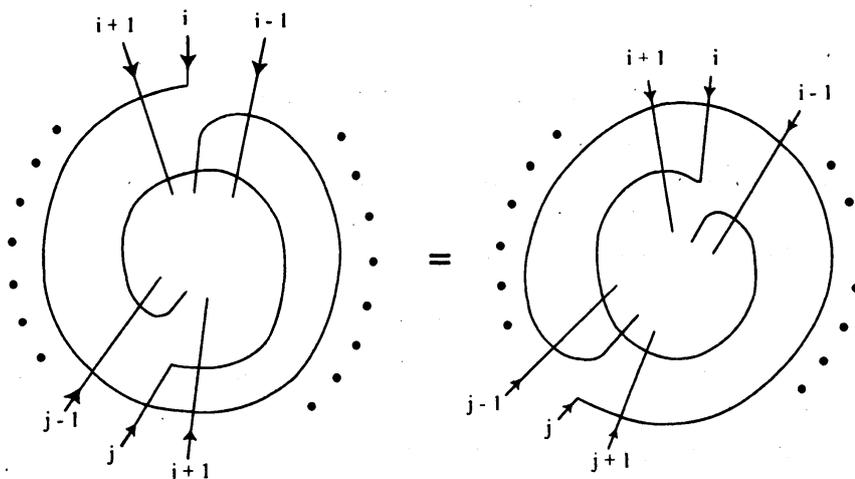


FIGURE 4. Commutativity of the transfer maps

One can easily check that the second diagram is a result of the several operations presented above on figures 1 and 2 applied to the first diagram. The second identity also follows from similar consideration (see fig. below where the case $n = 3$ is presented).

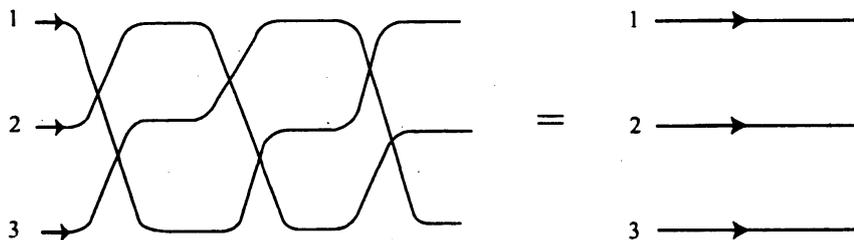


FIGURE 5. Product of transfer maps

To prove the second part it is actually enough to consider only the cases $n = 2$ and $n = 3$. When $n = 2$ we have two maps $T_1^{(2)} = R$ and $T_2^{(2)} = R_{21}$, so the relation (11) becomes simply the reversibility condition (2). The commutativity condition is trivial in this case. For $n = 3$ then we have three transfer maps: $T_1^{(3)} = R_{13}R_{12}$, $T_2^{(3)} = R_{21}R_{23}$, $T_3^{(3)} = R_{32}R_{31}$. The product $T_1^{(3)}T_2^{(3)}T_3^{(3)} = R_{13}R_{12}R_{21}R_{23}R_{32}R_{31}$ is obviously an identity because

$R_{12}R_{21} = R_{23}R_{32} = R_{13}R_{31} = Id$. Commutativity $T_1^{(3)}T_2^{(3)} = T_2^{(3)}T_1^{(3)}$ means that $R_{13}R_{12}R_{21}R_{23} = R_{21}R_{23}R_{13}R_{12}$ which is equivalent to Yang-Baxter equation (1) modulo the reversibility relation which was already shown.

The properties of the transfer maps can also be explained in terms of the corresponding extended affine Weyl group $\tilde{A}_{n-1}^{(1)}$. As we have already mentioned above the maps

$$S_i = P_{ii+1}R_{ii+1},$$

where P_{ij} is the permutation of i -th and j -th factors in X^n , satisfy the relations

$$(12) \quad S_i S_{i+1} S_i = S_{i+1} S_i S_{i+1}, \quad i = 1, \dots, n-1$$

and

$$(13) \quad S_i^2 = Id$$

which are the defining relations of the affine Weyl group $A_{n-1}^{(1)}$. However (see the warning after the definition of the Yang-Baxter maps) the transfer maps **do not** belong to this group but do belong to its extension $\tilde{A}_{n-1}^{(1)}$ generated by S_1, \dots, S_n and the cyclic permutation $\omega = P_{1n}P_{1n-1} \dots P_{12}$, $\omega^n = Id$ (known as *extended affine Weyl group*, see [51]). Moreover one can easily check that they corresponds to the translations in this groups and thus must commute (cf.[50]).

Note that if we consider only the set of S_i with $i = 1, \dots, n-1$ (without S_n) then they generate the (finite) permutation group S_n . In that case as it was shown by Etingof et al [4] the corresponding dynamics is equivalent to the usual action of S_n by permutations (and thus is trivial). It is very important that in our case the transfer maps generate the action of an infinite abelian group \mathbf{Z}^{n-1} so the question of whether two such actions are equivalent or not is a non-trivial problem of dynamical systems theory. The fact that the corresponding group is abelian indicates that the dynamics is integrable but to make this claim precise is not so easy.

Remark. We should mention that one can introduce another dynamical analogue of the transfer matrices, which may be even more natural but leads to the multi-valued maps (correspondences). Indeed consider the map

$$F = R_{0n}R_{0n-1} \dots R_{01} : (x_0, x_1, \dots, x_n) \rightarrow (y_0, y_1, \dots, y_n)$$

and assume that $x_0 = y_0$ (to mimic the trace operation). Generically one can express x_0 as a function of x_1, x_2, \dots, x_n from this relation (in not necessarily unique way though). Substituting this x_0 into the rest of the relations $y_1 = y_1(x_0, x_1, \dots, x_n), \dots, y_n = y_n(x_0, x_1, \dots, x_n)$ we come to the *transfer correspondence* $t^{(n)} : X^n \rightarrow X^n$, which can be considered as analogue of transfer matrices (7). The commutativity property (8) follows as in the usual case. More detail can be found in [52] where the corresponding multi-valued dynamics related to Adler maps and discrete KdV equation is investigated.

Now we are going to present the main constructions and the examples of the Yang-Baxter maps coming from the theory of integrable systems.

MATRIX FACTORISATIONS AND LAX REPRESENTATIONS

The relation between matrix factorisations and Yang-Baxter property goes back at least to A.B. Zamolodchikov's works in the late 70-th on factorised S-matrices and related algebras. I would like to mention here also the papers [2, 54, 55, 56, 57, 58]. The first concrete example of the Yang-Baxter map I know which came from this construction is the Adler map [34](see below).

Let $A(x, \zeta)$ be a family of matrices depending on the point $x \in X$ and a "spectral" parameter $\zeta \in \mathbf{C}$. One should think of X being an algebraic variety and A depending polynomially/rationally on ζ although the procedure works always when the corresponding factorisation problems are uniquely solvable. For matrix polynomials usually this is the case once the factorisation of the corresponding determinant is fixed (see [53]).

The following procedure is a version of standard QR-algorithm in linear algebra and known to be very useful in the theory of the discrete integrable systems, see e.g. [54].

Consider the product $L = A(y, \zeta)A(x, \zeta)$, then change the order of the factors $L \rightarrow \tilde{L} = A(x, \zeta)A(y, \zeta)$ and refactorise it again: $\tilde{L} = A(\tilde{y}, \zeta)A(\tilde{x}, \zeta)$. Now define the map R by the formula

$$(14) \quad R(x, y) = (\tilde{x}, \tilde{y}).$$

In other words the relation

$$(15) \quad A(x, \lambda; \zeta)A(y, \mu; \zeta) = A(\tilde{y}, \mu; \zeta)A(\tilde{x}, \lambda; \zeta),$$

is satisfied iff $(\tilde{x}, \tilde{y}) = R(\lambda, \mu)(x, y)$.

In that case the matrix $A(x, \zeta)$ is called the *Lax matrix* of the map R , which is automatically reversible Yang-Baxter map. Indeed the Yang-Baxter equation easily follows from associativity and the reversibility is obvious.

Once the Lax matrix is known one can introduce the *monodromy matrix*

$$M = A(x_n, \zeta)A(x_{n-1}, \zeta) \dots A(x_1, \zeta).$$

Theorem [24]. *The transfer maps $T_i^{(n)}$, $i = 1, \dots, n$ preserve the spectrum of the corresponding monodromy matrix $M(x_1, \dots, x_n; \zeta)$ for all ζ .*

The proof is simple. For the map $T_1^{(n)} = R_{1n}R_{1n-1} \dots R_{12}$ one can easily see that

$$A(x_n, \zeta)A(x_{n-1}, \zeta) \dots A(x_1, \zeta) = A(\tilde{x}_1, \zeta)A(\tilde{x}_n, \zeta)A(\tilde{x}_{n-1}, \zeta) \dots A(\tilde{x}_2, \zeta),$$

where $(\tilde{x}_1, \dots, \tilde{x}_n) = T_1^{(n)}(x_1, \dots, x_n)$. Since the spectrum of the product of the matrices is invariant under the cyclic permutation of the factors, we have

$$\text{Spec}M(x_1, \dots, x_n; \zeta) = \text{Spec}M(\tilde{x}_1, \dots, \tilde{x}_n; \zeta).$$

To prove the same for $T_i^{(n)}$ one should consider the matrix $A(x_{n+i-1}, \zeta) \dots A(x_i, \zeta)$ which obviously has the same spectrum as M and use the same arguments.

As a corollary we have that the coefficients of the characteristic polynomial $\chi = \det(M - \lambda I)$ are the integrals of the transfer maps $T_i^{(n)}$. If the dependence of A (and therefore M) on λ is polynomial then the dynamics can be linearised on the Jacobi varieties of the corresponding spectral curves $\chi(\zeta, \lambda) = 0$ (see [54]). If these maps are also symplectic then one can use also the discrete version of the Liouville theorem [37] to claim their integrability.

We should mention here that although the procedure is similar to the one proposed in [54] for the discrete Lagrangian systems there is one important difference: in the construction of Yang-Baxter maps the factors A should depend only on one of the variables, while in general scheme they usually depend both on x and y .

Example. Symmetries of the periodic dressing chain and recutting of polygons [33, 34].

In [34] Adler discovered a remarkable geometric dynamical system, which he called the *recutting of polygons*. It turned out that (after a slight modification) it can be described by the following refactorisation procedure related to the *periodic dressing chain* investigated by Shabat and the author [33].

Here $X = \mathbf{C} \times \mathbf{C}$ and the matrix $A(x)$, $x = (f, \beta)$ has the form:

$$(16) \quad A = \begin{pmatrix} f & f^2 + \beta - \zeta \\ 1 & f \end{pmatrix}$$

One can check that the factorisation procedure described above leads in this case to the following birational map $R : (f_1, \beta_1; f_2, \beta_2) \rightarrow (\tilde{f}_1, \beta_1; \tilde{f}_2, \beta_2) :$

$$(17) \quad \tilde{f}_1 = f_2 - \frac{\beta_1 - \beta_2}{f_1 + f_2} \quad \tilde{f}_2 = f_1 - \frac{\beta_2 - \beta_1}{f_1 + f_2},$$

which is the *Adler map*. Strictly speaking Adler considered the S -version of this map: $S = PR$. The observation that it satisfies the braid relations and thus leads to the action of the affine Weyl group was important for him, because this was related to the polynomial growth property [37]. In a different context this action was discussed in much more details by Noumi and Yamada [50].

The transfer dynamics in this example is symplectic (see also below) and integrable in the sense of discrete Liouville theorem [37] (see [34]). This follows from the theory of the dressing chain [33] which is closely related to the finite-gap theory of the KdV equation [35, 36]. In fact the transfer dynamics $T_i^{(n)}$ with odd $n = 2g + 1$ can be considered as a discrete version of the Novikov equations [35] describing the g -gap Schrödinger operators (see for details [34, 33]).

It turned out that quite often (in particular, in the previous example) one can derive the Lax matrix directly from the Yang-Baxter map. The

following observation is due to Suris and the author [26] and was inspired by the results [45, 46] on 3D consistent discrete equations.

Suppose that on the set X we have an effective action of the linear group $G = GL_N$, and that the Yang-Baxter map $R(\lambda, \mu)$ has the following special form:

$$(18) \quad \tilde{x} = B(y, \mu, \lambda)[x], \quad \tilde{y} = A(x, \lambda, \mu)[y],$$

where $A, B : X \times \mathbf{C} \times \mathbf{C} \rightarrow GL_N$ are some matrix valued functions on X depending on parameters λ and μ and $A[x]$ denotes the action of $A \in G$ on $x \in X$. Then the claim is that both $A(x, \lambda, \zeta)$ and $B^T(x, \lambda, \zeta)$ are Lax matrices for R . The claim about B is equivalent to the fact that $B(x, \lambda, \zeta)$ is a Lax matrix for R_{21} .

Indeed, look at the values of z_{12} produced by the both parts of the Yang-Baxter relation (5): the left-hand side gives $z_{12} = A(y_1, \mu, \nu)A(x_2, \lambda, \nu)[z]$, while the right-hand side gives $z_{12} = A(x, \lambda, \nu)A(y, \mu, \nu)[z]$. Now since we assume that the action of G is effective, we immediately arrive at the relation $A(x, \lambda, \nu)A(y, \mu, \nu) = A(y_1, \mu, \nu)A(x_2, \lambda, \nu)$, which holds whenever $(x_2, y_1) = R(\lambda, \mu)(x, y)$. This coincides with (14), an arbitrary parameter ν playing the role of the spectral parameter ζ .

Similarly, one could look at the values of x_{23} produced by the both parts of (5): the left-hand side gives $x_{23} = B(z, \nu, \lambda)B(y, \mu, \lambda)[x]$, while the right-hand side gives $x_{23} = B(y_3, \mu, \lambda)B(z_2, \nu, \lambda)[x]$. Effectiveness of the action of G again implies: $B(z, \nu, \lambda)B(y, \mu, \lambda) = B(y_3, \mu, \lambda)B(z_2, \nu, \lambda)$, whenever $(y_3, z_2) = R(\mu, \nu)(y, z)$. This turns into (14) for the transposed matrices B^T (or for the inverse matrices B^{-1}); the role of spectral parameter is here played by an arbitrary parameter λ .

In fact the effectiveness of the action can be replaced by the following weaker property. Let us call the action of $G = GL_N$ on X *projective* if the action of A on X is trivial only if A is a scalar. Similarly we say that $A(x, \lambda, \zeta)$ gives a *projective Lax representation* for the Yang-Baxter map R if the relation (15) holds up to multiplication by a scalar matrix cI , where c may depend on all the variables in the relation. One can easily modify the previous arguments to produce the integrals for the transfer maps using the projective Lax matrix: all the ratios of the eigenvalues of the monodromy matrix are obviously preserved by these maps.

For the projective action the matrices $A(x, \lambda, \zeta)$ and $B^T(x, \lambda, \zeta)$ give projective Lax representations for the corresponding Yang-Baxter maps (14). In practice however for a natural choice of matrices A, B we have actually proper Lax representations, as we can see from the example of Adler map (17).

Indeed one can rewrite (17) as follows

$$\tilde{y} = x - \frac{\mu - \lambda}{x + y} = \frac{x^2 + xy - (\mu - \lambda)}{x + y} = A(x, \lambda, \mu)[y],$$

where as before

$$A(x, \lambda, \zeta) = \begin{pmatrix} x & x^2 + \lambda - \zeta \\ 1 & x \end{pmatrix}.$$

and the group $G = GL_2$ is acting on \mathbf{CP}^1 by Möbius transformations.

Another interesting example when the Lax representation can be read directly from the form of the Yang-Baxter map is related to the theory of solitons, in particular matrix solitons.

INTERACTION OF SOLITONS AS A YANG-BAXTER MAP.

Consider any integrable by the inverse scattering method PDE in 1+1 dimensions which has multi-soliton solutions. Suppose that each soliton has a non-trivial internal degrees of freedom described by the set X (which is usually a manifold). The soliton interaction has the following remarkable property: the result of interaction of three solitons is independent of the order of collisions and thus is completely defined by the pairwise interaction. Although the proof of this claim in full generality seems to be absent in the literature it is confirmed in all known cases (see recent papers [40, 41, 42] for the discussion of this). This property means that the corresponding soliton interaction map R satisfies the Yang-Baxter relation.

A good example is the matrix KdV equation [59] :

$$U_t + 3UU_x + 3U_xU + U_{xxx} = 0,$$

U is $d \times d$ matrix.

It is easy to check that it has the soliton solution of the form

$$U = 2\lambda^2 P \operatorname{sech}^2(\lambda x - 4\lambda^3 t),$$

where P must be a projector: $P^2 = P$. If P has rank 1 then it should have the form $P = \frac{\xi \otimes \eta}{(\xi, \eta)}$, where ξ is a vector in a vector space V of dimension d , η is a vector from the dual space V^* (covector) and bracket (ξ, η) means the canonical pairing between V and V^* .

The change of the matrix amplitudes P ("polarisations") of two solitons with the velocities λ_1 and λ_2 after their interaction is described by the following map [40, 25]:

$$R(\lambda_1, \lambda_2) : (\xi_1, \eta_1; \xi_2, \eta_2) \rightarrow (\tilde{\xi}_1, \tilde{\eta}_1; \tilde{\xi}_2, \tilde{\eta}_2)$$

$$(19) \quad \tilde{\xi}_1 = \xi_1 + \frac{2\lambda_2(\xi_1, \eta_2)}{(\lambda_1 - \lambda_2)(\xi_2, \eta_2)} \xi_2, \quad \tilde{\eta}_1 = \eta_1 + \frac{2\lambda_2(\xi_2, \eta_1)}{(\lambda_1 - \lambda_2)(\xi_2, \eta_2)} \eta_2,$$

$$(20) \quad \tilde{\xi}_2 = \xi_2 + \frac{2\lambda_1(\xi_2, \eta_1)}{(\lambda_2 - \lambda_1)(\xi_1, \eta_1)} \xi_1, \quad \tilde{\eta}_2 = \eta_2 + \frac{2\lambda_1(\xi_1, \eta_2)}{(\lambda_2 - \lambda_1)(\xi_1, \eta_1)} \eta_1.$$

In this example X is the set of projectors P of rank 1 which is the variety $\mathbf{CP}^{N-1} \times \mathbf{CP}^{N-1}$, and the group $G = GL_N$ is acting on the projectors by conjugation (which corresponds to the natural action of $GL(V)$ on $V \otimes V^*$).

It is easy to see that the formulas (19), (20) are of the form (18) with the matrices

$$(21) \quad A(\xi, \eta, \lambda; \zeta) = I + \frac{2\lambda}{\zeta - \lambda} \frac{\xi \otimes \eta}{(\xi, \eta)} = I + \frac{2\lambda}{\zeta - \lambda} P$$

(note that again $R_{21} = R$). Our results show that the matrix $A(P, \lambda, \zeta)$ gives a projective Lax representation for the interaction map.

One can show that this is actually a genuine Lax representation. We will demonstrate this for the class of more general Yang-Baxter maps on Grassmannians following to [25].

Let V be an n -dimensional real (or complex) vector space, $P : V \rightarrow V$ be a projector of rank k : $P^2 = P$. Any such projector is uniquely determined by its kernel $K = \text{Ker}P$ and image $L = \text{Im}P$, which are two subspaces of V of dimensions k and $n - k$ complementary to each other: $K \oplus L = V$. The space of all projectors X of rank k is an open set in the product of two Grassmannians $G(k, n) \times G(n - k, n)$.

Consider the following matrix

$$(22) \quad A(P, \lambda; \zeta) = I + \frac{2\lambda}{\zeta - \lambda} P$$

and the related re-factorization relation

$$(23) \quad \left(I + \frac{2\lambda_1}{\zeta - \lambda_1} P_1\right) \left(I + \frac{2\lambda_2}{\zeta - \lambda_2} P_2\right) = \left(I + \frac{2\lambda_2}{\zeta - \lambda_2} \tilde{P}_2\right) \left(I + \frac{2\lambda_1}{\zeta - \lambda_1} \tilde{P}_1\right)$$

which we can rewrite in the polynomial form as

$$(24) \quad ((\zeta - \lambda_1)I + 2\lambda_1 P_1)((\zeta - \lambda_2)I + 2\lambda_2 P_2) = ((\zeta - \lambda_2)I + 2\lambda_2 \tilde{P}_2)((\zeta - \lambda_1)I + 2\lambda_1 \tilde{P}_1).$$

The claim is that if $\lambda_1 \neq \pm\lambda_2$ it has a unique solution. This follows from the general theory of matrix polynomials (see e.g. [53]) but in this case we can see this directly.

Indeed let us compare the kernels of both sides of the relation (24) when the spectral parameter $\zeta = \lambda_1$. In the right hand side we obviously have \tilde{K}_1 while the left hand side gives

$$((\lambda_1 - \lambda_2)I + 2\lambda_2 P_2)^{-1} K_1 = \left(I + \frac{2\lambda_2}{\lambda_1 - \lambda_2} P_2\right)^{-1} K_1.$$

Now we use the following property of the matrix (22):

$$(25) \quad A(P, -\lambda; \zeta) = A(P, \lambda; \zeta)^{-1}$$

to have

$$(26) \quad \tilde{K}_1 = \left(I - \frac{2\lambda_2}{\lambda_1 + \lambda_2} P_2\right) K_1.$$

Similarly taking the image of both sides of (24) at $\zeta = \lambda_2$ we will have

$$(27) \quad \tilde{L}_2 = \left(I + \frac{2\lambda_1}{\lambda_2 - \lambda_1} P_1\right) L_2.$$

To find \tilde{K}_2 and \tilde{L}_1 one should take first the inverse of both sides of (23), use the property (25) and then repeat the procedure. This will lead us to the formulas:

$$(28) \quad \tilde{K}_2 = \left(I - \frac{2\lambda_1}{\lambda_1 + \lambda_2} P_1\right) K_2$$

and

$$(29) \quad \tilde{L}_1 = \left(I + \frac{2\lambda_2}{\lambda_1 - \lambda_2} P_2\right) L_1.$$

The formulas (26,27,28,29) determine a parameter-dependent Yang-Baxter map on the set of projectors. One can easily check that for $k = 1$ one has the formulas (19, 20) for two matrix soliton interaction.

If we supply now our vector space V with the Euclidean (Hermitian) structure and consider the self-adjoint projectors P of rank k then the corresponding space X will coincide with the Grassmannian $G(k, n)$: such a projector is completely determined by its image L (which is a k -dimensional subspace in V and thus a point in $G(k, n)$) since the kernel K in this case is the orthogonal complement to L .

The corresponding Yang-Baxter map R on the Grassmannian is determined by the formulas

$$(30) \quad \tilde{L}_1 = \left(I + \frac{2\lambda_2}{\lambda_1 - \lambda_2} P_2\right) L_1,$$

$$(31) \quad \tilde{L}_2 = \left(I + \frac{2\lambda_1}{\lambda_2 - \lambda_1} P_1\right) L_2.$$

Yang-Baxter maps related to geometric crystals. In this section we briefly discuss the examples of Yang-Baxter maps related to geometric crystals and tropical combinatorics [13, 14, 19, 22]. We show how one can derive the Lax representation for them following to [26].

Let $X = \mathbf{C}^n$, and define the map $R : X \times X \rightarrow X \times X$ by the formulas

$$(32) \quad \tilde{x}_j = x_j \frac{P_j}{P_{j-1}}, \quad \tilde{y}_j = y_j \frac{P_{j-1}}{P_j}, \quad j = 1, \dots, n,$$

where

$$(33) \quad P_j = \sum_{a=1}^n \left(\prod_{k=1}^{a-1} x_{j+k} \prod_{k=a+1}^n y_{j+k} \right)$$

(in this formula subscripts $j+k$ are taken (mod n)). The map (32) obviously preserves the following subsets invariant: $X_\lambda \times X_\mu \subset X \times X$, where $X_\lambda = \{(x_1, \dots, x_n) \in X : \prod_{k=1}^n x_k = \lambda\}$. Let us show that the restriction of R to $X_\lambda \times X_\mu$ can be written in the form (18). For this, let us embed this set into $\mathbf{CP}^{n-1} \times \mathbf{CP}^{n-1}$:

$$J(x, y) = (z(x), w(y)), \quad z(x) = (1 : z_1 : \dots : z_{n-1}), \quad w(y) = (w_1 : \dots : w_{n-1} : 1),$$

$$z_j = \prod_{k=1}^j x_k, \quad w_j = \prod_{k=j+1}^n y_k.$$

Then it is easy to see that in coordinates (z, w) the map R is written as

$$\tilde{z} = B(y, \mu, \lambda)[z], \quad \tilde{w} = A(x, \lambda, \mu)[w],$$

with certain matrices B, A from $G = GL_n$, where the standard projective action of GL_n on \mathbf{CP}^{n-1} is used. A simple calculation shows that the inverse matrices are the following matrices well-known in this context [14, 19, 22]:

$$(34) \quad B^{-1}(y, \mu, \lambda) = \begin{pmatrix} y_1 & -1 & 0 & \dots & 0 & 0 \\ 0 & y_2 & -1 & \dots & 0 & 0 \\ 0 & 0 & y_3 & \dots & 0 & 0 \\ & \dots & & & \dots & \\ 0 & 0 & 0 & \dots & y_{n-1} & -1 \\ -\lambda & 0 & 0 & \dots & 0 & y_n \end{pmatrix},$$

$$(35) \quad A^{-1}(x, \lambda, \mu) = \begin{pmatrix} x_1 & 0 & 0 & \dots & 0 & -\mu \\ -1 & x_2 & 0 & \dots & 0 & 0 \\ 0 & -1 & x_3 & \dots & 0 & 0 \\ & \dots & & & \dots & \\ 0 & 0 & 0 & \dots & x_{n-1} & 0 \\ 0 & 0 & 0 & \dots & -1 & x_n \end{pmatrix}.$$

Thus we derived again the Lax matrices directly from the map itself. It is worthy to mention that this is closely related to the notion of the *structure group* G_R of the Yang-Baxter map R [14]. It was shown by Etingof in [14] that for the map (32) the so-called reduced structure groups G_R^+ and G_R^- can be realised as the subgroups of the loop group $PGL_n(C(\lambda))$ generated by the matrix functions $A^{-1}(x, \cdot, \lambda)$ with $x \in X$, resp. by $B^{-1}(x, \cdot, \lambda)$ with $x \in X$.

POISSON LIE GROUPS AND HAMILTONIAN STRUCTURE FOR THE TRANSFER DYNAMICS

In this section we describe the Hamiltonian structure for the transfer maps following the recent work by Reshetikhin and the author [27] partly based on some general constructions from Weinstein and Xu [3]. The main observation is that if the Lax matrices form symplectic leaves in certain Poisson Lie group then the transfer dynamics is Poisson.

For a good account of the theory of the Poisson Lie group we recommend Chapter 5 in the Reiman and Semenov-Tian-Shansky book [60] or original Drinfeld paper [61].

Recall that *Poisson Lie group* is a Lie group G with the Poisson structure on it such that the multiplication map $\mu : G \times G \rightarrow G$ is Poisson. Recall also that the map between two Poisson manifolds is called *Poisson* if it preserves the Poisson bracket.

We will need also the notion of *Poisson correspondence* (or Poisson relation) (see e.g. Weinstein [62] and references therein). The *correspondence* (relation) Φ from a set X into itself is a multivalued, partially defined map determined by its graph Γ_Φ , which is a subset of $X \times X$. Now let $X = M$ be a smooth manifold with Poisson structure J and the graph Γ_Φ is a submanifold of $M \times M$. Let us supply the manifold $M \times M$ with the Poisson structure $J \oplus (-J)$. Recall now that a submanifold N of a Poisson manifold P is called *coisotropic* if the set of functions on P which vanish on N is closed under Poisson bracket. The correspondence Φ is called *Poisson* if its graph Γ_Φ is coisotropic submanifold of $M \times M$. For the usual maps this equivalent to the standard notion of Poisson map. If some smoothness condition is fulfilled then the composition of two Poisson correspondences is again Poisson.

For any Poisson Lie group G one can construct the following correspondence $\Phi_G : G \times G \rightarrow G \times G$: its graph Γ_Φ is the set of $(g_1, g_2; h_1, h_2) \in G \times G \times G \times G$ such that

$$(36) \quad g_1 g_2 = h_2 h_1.$$

From the definition of the Poisson Lie group it follows that this correspondence is Poisson.

The first simple observation is that this correspondence satisfies the Yang-Baxter relation. Note that for the correspondence $R : X \times X \rightarrow X \times X$ the map R_{12} is defined by its graph which is a subset of $X^{(6)}$ of the form $\Gamma_R \times \Delta$, where $\Delta = (x, x)$, $x \in X$ is the diagonal in $X \times X$. Similarly one can define R_{ij} and thus *Yang-Baxter correspondence* as such R which satisfies the relation (1). Reversibility is a bit more tricky. The problem is with the notion of the inverse correspondence Φ^{-1} : the standard formula $\Phi\Phi^{-1} = Id$ seems to be too restrictive, so we should avoid it. There are two involutions on $X \times X \times X \times X$:

$$\pi : (x_1, x_2; y_1, y_2) = (x_2, x_1; y_2, y_1)$$

and

$$\tau : (x_1, x_2; y_1, y_2) = (y_1, y_2; x_1, x_2).$$

We will call the correspondence R *reversible* if $\pi\Gamma_R = \tau\Gamma_R$, or equivalently its graph Γ_R is invariant under the involution

$$\sigma = \pi\tau : (x_1, x_2; y_1, y_2) = (y_2, y_1; x_2, x_1).$$

For the maps this is equivalent to the definition given above.

We claim that for any Poisson Lie group G the Poisson correspondence Φ_G is a reversible Yang-Baxter correspondence. Indeed the Yang-Baxter property follows from the associativity of multiplication in G : both left and right sides are the correspondences with the same graph given by the relation

$$g_1 g_2 g_3 = h_3 h_2 h_1.$$

Reversibility is obvious.

In such generality this claim is probably not interesting enough, since we are primarily interested in genuine maps but not correspondences. A natural idea therefore is to restrict the previous construction to the symplectic leaves in G in order to make the re-factorisation relation (36) uniquely solvable.

Suppose that there exist a one-parameter family of embeddings of the set X as a symplectic leaf in the Poisson Lie group G : $\phi_\lambda : X \rightarrow G$ and define the correspondence $R(\lambda, \mu) : X \times X \rightarrow X \times X$ by the relation

$$(37) \quad \phi_\lambda(x)\phi_\mu(y) = \phi_\mu(\tilde{y})\phi_\lambda(\tilde{x}).$$

Note that the embedding ϕ_λ induce on X the symplectic structure

$$(38) \quad \omega_\lambda = \phi_\lambda^*(\omega_G),$$

where ω_G is the symplectic structure on the corresponding leaf in G . Let us introduce on $X \times X$ a symplectic structure as the direct sum $\omega_\lambda \oplus \omega_\mu$. In a similar way we define the symplectic structure $\Omega^{(N)}$ on $X^{(N)}$ as the direct sum of the forms $\omega_i = \omega_{\lambda_i}$:

$$(39) \quad \Omega^{(N)} = \omega_1 \oplus \omega_2 \oplus \dots \oplus \omega_N.$$

Theorem [3, 27]. *The correspondence $R(\lambda, \mu)$ defined by (37) is a reversible Yang-Baxter Poisson correspondence. The related transfer maps (9) are Poisson correspondences with respect to the symplectic structure $\Omega^{(N)}$.*

This construction is already good enough to produce many interesting examples of symplectic Yang-Baxter maps, including the interaction of matrix solitons [24, 25]. To show this we consider the standard Poisson Lie loop group $G = LGL(n, \mathbf{R})$ and the symplectic leaves of a special form.

The standard Poisson structure on the loop group $G = LGL(n, \mathbf{R})$ is given by the Sklyanin formula

$$(40) \quad \{L_1(\zeta), \otimes L_2(\eta)\} = \left[\frac{P_{12}}{\zeta - \eta}, L_1(\zeta) \otimes L_2(\eta) \right],$$

where P_{12} is the permutation matrix: $P(x \otimes y) = y \otimes x$.

We will consider the symplectic leaves of the form $L(\zeta) = A + B\zeta$. The relation (40) is equivalent to the following set of relations

$$(41) \quad \{A_1, \otimes A_2\} = P_{12}(A_1 B_2 - B_1 A_2), \quad \{A_1, \otimes B_2\} = 0, \quad \{B_1, \otimes B_2\} = 0.$$

From these relations it follows that the matrix elements of B are in the centre of the corresponding Poisson algebra, so B can be fixed.

If B is non-degenerate one can assume that $B = I$ is identity matrix. The rest of the relations (41) determine the standard Lie-Poisson structure on the matrix group $GL(n, \mathbf{R})$ considered as an open set in the $n \times n$ -matrix algebra M_n which can be identified as a space with $gl^*(n, \mathbf{R})$.

Let \mathcal{O}_k be the special symplectic leaf in $gl^*(n, \mathbf{R})$ consisting of involutions $S = 2P - I$, where P is a projector of rank k : $P^2 = P, S^2 = I$. Such a projector P is determined by its kernel $K = Ker P$ and its image $L = Im P$, which are two subspaces of dimension k and $n - k$ respectively such that $K \oplus L = \mathbf{R}^n$. Thus X is an open subset of the product of two Grassmannians

$G(n, k)$ and $G(n, n-k)$. Note that the Grassmannian $G(n, n-k)$ is naturally isomorphic to the Grassmannian $G(n, k)$ in the dual space.

Let $X = \mathcal{O}_k$ and define the embedding ϕ_λ of X into G by the formula

$$(42) \quad \phi_\lambda(S) = \zeta I + \lambda S = (\zeta - \lambda)I + 2\lambda P.$$

It is easy to see that the formulas (37), (42) lead to the Yang-Baxter map (26,27,28,29) considered in the previous section in relation with matrix KdV equation. In the rank 1 case we have the formulas (19), (20).

The corresponding transfer maps $T_i^{(N)}$ (9) are symplectic and preserve the spectrum of the monodromy matrix

$$(43) \quad M(\zeta) = \mathcal{L}(P_1, \lambda_1, \zeta)\mathcal{L}(P_2, \lambda_2, \zeta) \dots \mathcal{L}(P_N, \lambda_N, \zeta),$$

where $\mathcal{L}(P, \lambda, \zeta) = (\zeta - \lambda)I + 2\lambda P$. The coefficients of the characteristic polynomial $\chi(\zeta, \eta) = \det(\mathcal{L} - \eta I)$ give the set of integrals for transfer dynamics, which are in involution with respect to the symplectic structure $\Omega^{(N)}$ on $X^{(N)}$.

We believe that this implies integrability of the transfer dynamics in Liouville sense, but the proof of completeness is still to be found. As usual the dynamics is linear on the Jacobi variety of the spectral curve $\chi(\zeta, \eta) = 0$.

Let us consider now the symplectic leaves in $LGL(2, \mathbf{R})$ of the form $L = A + B\zeta$ with degenerate

$$B = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}.$$

Writing A as

$$A = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

we have the following Poisson brackets for the matrix elements from the relation (40):

$$\{a, b\} = a, \quad \{a, d\} = c, \quad \{a, c\} = 0, \quad \{b, c\} = 0, \quad \{b, d\} = d, \quad \{c, d\} = 0.$$

It is easy to check that the centre of this Poisson algebra is generated by two functions: $C_1 = c$ and $C_2 = ad - bc$. If we choose $C_1 = c = 1$ and $C_2 = ad - bc = \lambda$ we come to the symplectic leaf $\mathcal{O}(\lambda)$ consisting of the matrices of the form

$$(44) \quad L(\zeta, \lambda) = \begin{pmatrix} a & ad + \lambda - \zeta \\ 1 & d \end{pmatrix},$$

where $\{a, d\} = 1$.

This is not quite Lax matrix (16) for the Adler map yet. To have it one should apply the Hamiltonian reduction to the product of the corresponding orbits $\mathcal{O}(\lambda)$ (see [27] for the details). This leads to the Poisson structure for the transfer dynamics for Adler maps, which is known from the theory of the periodic dressing chain [33, 34] (with an odd period N).

CLASSIFICATION OF YANG-BAXTER MAPS: QUADRIRATIONAL CASE

The classification of the Yang-Baxter maps is probably the most interesting but hopelessly difficult problem even in the simplest case of $X = \mathbf{CP}^1$. A natural class where the complete classification seems to be realistic is given by the quadrirational maps introduced and investigated by Adler, Bobenko and Suris [43].

The map $R : (x, y) \rightarrow (u, v)$ is called *quadrirational* if both R and so called *companion map* $\bar{R} : (x, v) \rightarrow (y, u)$ are birational maps of $\mathbf{CP}^1 \times \mathbf{CP}^1$ into itself. All such maps have the form

$$(45) \quad R : u = \frac{a(y)x + b(y)}{c(y)x + d(y)}, \quad v = \frac{A(x)y + B(x)}{C(x)y + D(x)},$$

where $a(y), \dots, d(y)$ and $A(x), \dots, D(x)$ are polynomials of degree at most 2. There are three subclasses of such maps, denoted as [1:1], [1:2] and [2:2], corresponding to the highest degrees of the coefficients of both fractions in (45). In [43] the classification of such maps was done modulo group $(\mathcal{M}\ddot{o}b)^4$, where $\mathcal{M}\ddot{o}b$ is the group of projective transformations of \mathbf{CP}^1 :

$$z \rightarrow w = \frac{az + b}{cz + d}.$$

In the richest class [2:2] it looks as follows.

Theorem [43]. *Any quadrirational map on $\mathbf{CP}^1 \times \mathbf{CP}^1$ of subclass [2:2] is equivalent, under some change of variables from $(\mathcal{M}\ddot{o}b)^4$, to exactly one of the following five maps:*

$$(F_I) \quad u = \alpha y P, \quad v = \beta x P, \quad P = \frac{(1 - \beta)x + \beta - \alpha + (\alpha - 1)y}{\beta(1 - \alpha)x + (\alpha - \beta)yx + \alpha(\beta - 1)y},$$

$$(F_{II}) \quad u = \frac{y}{\alpha} P, \quad v = \frac{x}{\beta} P, \quad P = \frac{\alpha x - \beta y + \beta - \alpha}{x - y},$$

$$(F_{III}) \quad u = \frac{y}{\alpha} P, \quad v = \frac{x}{\beta} P, \quad P = \frac{\alpha x - \beta y}{x - y},$$

$$(F_{IV}) \quad u = y P, \quad v = x P, \quad P = 1 + \frac{\beta - \alpha}{x - y},$$

$$(F_V) \quad u = y + P, \quad v = x + P, \quad P = \frac{\alpha - \beta}{x - y},$$

with some suitable constants α, β .

It is remarkable that all these maps can be described by the following natural geometric construction [43].

Consider a pair of nondegenerate conics Q_1, Q_2 on the plane \mathbf{CP}^2 , which are rational curves isomorphic to \mathbf{CP}^1 . Take $X \in Q_1, Y \in Q_2$, and let

$\ell = \overline{XY}$ be the line through X, Y . Generically the line ℓ intersects Q_1 and Q_2 at two points U and V (see Fig. 6). This defines a map $\mathcal{R} : (X, Y) \mapsto (U, V)$.

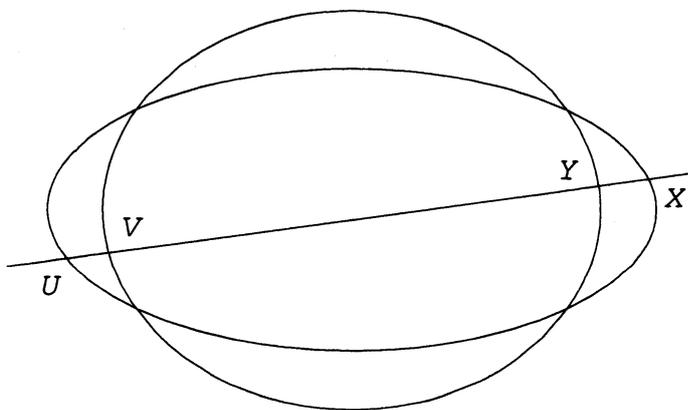


FIGURE 6. A quadrirational map on a pair of conics

There are five possible types I – V of intersection of two conics (see [63]):

- I: four simple intersection points;
- II: two simple intersection points and one point of tangency;
- III: two points of tangency;
- IV: one simple intersection point and one point of the second order tangency;
- V: one point of the third order tangency.

In suitable rational parametrisations of the conics we have exactly the five maps listed in the theorem above.

It is remarkable that all five representative maps from the last theorem satisfy the Yang-Baxter relation (see [43], where a nice geometric interpretation of this relation is also given). This surprising fact, which must be related to the classification approach of [43] based on the analysis of the singularities, does not imply however that all quadrirational maps are Yang-Baxter (as one might conclude from [43]) since the action of $(\mathcal{M}\ddot{o}b)^4$ in general destroy the Yang-Baxter property.

For example, if we change in the fifth map F_V the variables $x \rightarrow -x, y \rightarrow -y$ (preserving u, v) we come to the quadrirational map $u = -y - P, v = -x - P, P = \frac{\alpha - \beta}{x - y}$, which does not satisfy the Yang-Baxter relation. However, if we change $u \rightarrow -u, y \rightarrow -y$ we come to the Adler map (17), which is a Yang-Baxter map. Note that neither of these changes is a proper conjugation of maps of $\mathbf{CP}^1 \times \mathbf{CP}^1$ into itself (in particular, F_V is not conjugated to Adler map).

A question remains therefore to classify all quadrirational Yang-Baxter maps modulo diagonal action of $\mathcal{M}\ddot{o}b$, which seems to be a natural equivalence within this class. The previous example shows that in the same

($\mathcal{M}\ddot{o}b$)⁴-orbit there are several Yang-Baxter representatives, non-equivalent under $\mathcal{M}\ddot{o}b$ -conjugation. The answer to this question can be probably extracted from the results of the paper [43].

Let us mention in this relation that a complete classification of integrable (3D-consistent) discrete equations on quad-graphs is known due to the same authors [44]. Since these two problems are closely related (see [44, 65]) a very interesting question is what could be a Yang-Baxter analogue of the "master equation" Q_4 in their classification (see [64] for the discussion of this very important case). Probably one should look for a wider class of Yang-Baxter maps (like correspondences or Cremona maps). A class of matrix factorisations suggested by Odesskii [58] could give one of the possible answers to this question.

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