## **Research** Article

# Hardy-Type Space Associated with an Infinite-Dimensional Unitary Matrix Group

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We investigate an orthogonal system of the homogenous Hilbert-Schmidt polynomials with respect to a probability measure which is invariant under the right action of an infinite-dimensional unitary matrix group. With the help of this system, a corresponding Hardy-type space of square-integrable complex functions is described. An antilinear isomorphism between the Hardy-type space and an associated symmetric Fock space is established.

## 1. Introduction

We investigate an orthogonal system of the Hilbert-Schmidt polynomials in the space  $L_{\chi}^2$  of square-integrable complex functions on the projective limit  $\mathfrak{U} = \lim_{\leftarrow} U(m)$  of unitary  $(m \times m)$ -dimensional matrix groups U(m)  $(m \in \mathbb{N})$ , called the space of virtual unitary matrices and endowed with the projective limit measure  $\chi = \lim_{\leftarrow} \chi_m$  of the probability Haar measures  $\chi_m$  on U(m). The measure  $\chi$  on the space  $\mathfrak{U}$  is invariant under the right action of the infinite-dimensional unitary group  $U(\infty) \times U(\infty)$ , where  $U(\infty) = \bigcup_m U(m)$ .

The space of virtual unitary matrices  $\mathfrak{U}$  was studied by Neretin [1] and Olshanski [2]. This notion relates to D. Pickrell's space of virtual Grassmannian [3] and to Kerov, Olshanski, and Vershik's space of virtual permutations [4]. Various spaces of integrable functions with respect to measures that are invariant under infinite-dimensional groups have been widely applied in stochastic processes [5], infinitedimensional probability [6, 7], complex analysis [8], and so forth.

The main results of the present paper are Theorems 6-7 that describe a Hardy-type subspace  $\mathscr{H}_{\chi}^2 \subset L_{\chi}^2$  spanned by the finite type homogenous Hilbert-Schmidt polynomials that are generated by an associated symmetric Fock space.

### 2. Preliminaries

We consider the following infinite-dimensional unitary matrix groups:

$$U(\infty) = \bigcup \{U(m) : m \in \mathbb{N}\},\$$
$$U^{2}(\infty) := U(\infty)U(\infty),$$
(1)

where U(m) is the group of unitary  $(m \times m)$ -matrices which is identified with the subgroup in U(m + 1) fixing the (m + 1)th basis vector. In other words,  $U(\infty)$  is the group of infinite unitary matrices  $u = [u_{ij}]_{i,j\in\mathbb{N}}$  with finitely many matrix entries  $u_{ij}$  distinct from  $\delta_{ij}$ . We equip every group U(m) with the probability Haar measure  $\chi_m$ .

Following [1, 2], every matrix  $u_m \in U(m)$  with m > 1, we write in the following block matrix form:

$$u_m = \begin{bmatrix} z_{m-1} & a \\ b & t \end{bmatrix},\tag{2}$$

corresponding to the partition m = (m-1) + 1 so that  $z_{m-1} \in U(m-1)$  and  $t \in \mathbb{C}$ . Over the group  $U(\infty)$  (resp., U(m)) the right action is well defined:

$$u \cdot g = w^{-1}uv, \tag{3}$$

where *u* belongs to  $U(\infty)$  (resp., to U(m)) and g = (v, w) belongs to  $U^2(\infty)$  (resp., to  $U^2(m) := U(m) \times U(m)$ ). In [1, Proposition 0.1], [2, Lemma 3.1], it was proven that the following Livšic-type mapping:

$$\pi_{m-1}^{m}: U(m) \ni u_{m} \longrightarrow u_{m-1} \in U(m-1), \qquad (4)$$

such that

$$\begin{bmatrix} z_{m-1} & a \\ b & t \end{bmatrix} \longmapsto \begin{cases} z_{m-1} - a(1+t)^{-1}b \colon t \neq -1, \\ z_{m-1} \colon t = -1, \end{cases}$$
(5)

(which is not a group homomorphism) is Borel and surjective onto U(m-1) and commutes with the right action of  $U^2(m-1)$ .

As is known [1, Theorem 1.6], the pullback of the probability Haar measure  $\chi_{m-1}$  on U(m - 1) under the mapping  $\pi_{m-1}^m$  is the probability Haar measure  $\chi_m$  on U(m), that is,

$$\chi_{m-1} \circ \pi_{m-1}^m = \chi_m. \tag{6}$$

Let  $U'(m) \subset U(m)$  be the subset of unitary matrices which do not have  $\{-1\}$ , as an eigenvalue. Then, U'(m) is open in U(m), and the complement  $U(m) \setminus U'(m)$  is a  $\chi_m$ negligible set. Moreover (see [2, Lemma 3.11]), the mapping

$$\pi_{m-1}^{m}: U'(m) \longrightarrow U'(m-1)$$
(7)

is continuous and surjective.

Consider the projective limits, taken with respect to the surjective Borel projections  $\pi_{m-1}^m$  and their continuous restrictions  $\pi_{m-1}^m|_{U'(m)}$ , respectively,

$$\mathfrak{U} = \lim_{\longleftarrow} U(m), \qquad \mathfrak{U}' = \lim_{\longleftarrow} U'(m), \qquad (8)$$

called the spaces of virtual unitary matrices. Notice that  $\mathfrak{U}$  is a Borel subset in the Cartesian product  $X_{m \in \mathbb{N}} U(m) = \{u = (u_m) : u_m \in U(m)\}$  endowed with the product topology, because all mapping  $\pi_{m-1}^m$  are Borel. Moreover, the canonical projections

$$\pi_m: \mathfrak{U} \longrightarrow U(m), \qquad \pi_m: \mathfrak{U}' \longrightarrow U'(m), \qquad (9)$$

such that  $\pi_{m-1} = \pi_{m-1}^m \circ \pi_m$ , are surjective by surjectivity of  $\pi_{m-1}^m$  and  $\pi_{m-1}^m|_{U'(m)}$ .

Following [2, Lemma 4.8], [1, Section 3.1], with the help of the Kolmogorov consistent theorem, we uniquely define a probability measure  $\chi$  on  $\mathfrak{U}'$  as the projective limit under the mapping (6),

$$\chi = \lim_{\longleftarrow} \chi_m,\tag{10}$$

which satisfies the equality  $\chi = \chi_m \circ \pi_m$  for all  $m \in \mathbb{N}$ . On  $\mathfrak{U} \setminus \mathfrak{U}'$ , the measure  $\chi$  is zero, because  $\chi_m$  is zero on  $U(m) \setminus U'(m)$  for all  $m \in \mathbb{N}$ .

Using (3), right action of the group  $U^2(\infty)$  on the space of virtual unitary matrices  $\mathfrak{U}$  can be defined (see [2, Definition 4.5]) as follows:

$$\pi_m(u \cdot g) = w^{-1}\pi_m(u) v, \quad u \in \mathfrak{U}, \tag{11}$$

where *m* is so large that  $g = (v, w) \in U^2(m)$ .

The canonical dense embedding  $\iota : U(\infty) \hookrightarrow \mathfrak{U}$  to any element  $u_m \in U(m)$  assigns the unique sequence  $u = (u_l)_{l \in \mathbb{N}}$ , such that

$$u_{l} = \begin{cases} \pi_{l}^{l+1} \circ \cdots \circ \pi_{m-1}^{m} (u_{m}) \colon l < m, \\ u_{m} \colon l = m, \\ \begin{bmatrix} u_{m} & 0 \\ 0 & 1_{l-m} \end{bmatrix} \colon l > m, \end{cases}$$
(12)

where  $\mathbb{1}_{l-m}$  is the unit in U(l - m). So, the image  $\iota \circ U(\infty)$  consists of stabilizing sequences in  $\mathfrak{U}$  (see [2, Section 4]).

#### 3. Invariant Probability Measure

In what follows, we will endow the space of virtual unitary matrices  $\mathfrak{U}$  with the measure  $\chi = \lim_{\leftarrow} \chi_m$ . A complex function on  $\mathfrak{U}$  is called cylindrical [2, Definition 4.5] if it has the following form:

$$f(u) = (f_m \circ \pi_m)(u), \quad u \in \mathfrak{U}, \tag{13}$$

for a certain  $m \in \mathbb{N}$  and a certain complex function  $f_m$  on U(m).

Any continuous bounded function f on  $\mathfrak{U}'$  has a unique  $\chi$ -essentially bounded extension on  $\mathfrak{U}$ , because the set  $\mathfrak{U} \setminus \mathfrak{U}'$  is  $\chi$ -negligible. Therefore, if the function  $U'(m) \ni \pi_m(u) \mapsto f_m[\pi_m(u)]$  in the definition (13) is continuous and bounded, then the corresponding cylindrical function f is  $\chi$  essentially bounded.

By  $\mathscr{L}_{\chi}^{\infty}$ , we denote closure of the algebraic hull of all cylindrical  $\chi$ -essentially bounded functions (13) with respect to the following norm:

$$\|f\|_{\mathscr{L}^{\infty}_{\chi}} = \operatorname{ess\,sup}_{u \in \mathfrak{U}} |f(u)|.$$
(14)

**Lemma 1.** The measure  $\chi = \lim_{\leftarrow} \chi_m$  on  $\mathfrak{U}$  is a Radon probability measure such that

$$\int_{\mathfrak{U}} f(u \cdot g) d\chi(u) = \int_{\mathfrak{U}} f(u) d\chi(u), \qquad (15)$$

for all  $g \in U^2(\infty)$  and  $f \in \mathscr{L}^{\infty}_{\chi}$ . For any compact set  $K \subset U(m)$  the following equality holds:

$$(\chi \circ \iota)(K) = \chi_m(K).$$
(16)

*Proof.* Recall the Prohorov criterion, which is adapted to our notation (see [9, Chapter IX.4.2, Theorem 1] or [6, Theorem 6]): there exists a Radon probability measure  $\chi'$  on  $\mathfrak{U}'$  such that

$$\chi' = \chi_m \circ \pi_m |_{\mathfrak{U}'} \quad \forall m \in \mathbb{N}, \tag{17}$$

if and only if for every  $\varepsilon > 0$  there exists a compact set  $\mathscr{K}$  in  $\mathfrak{U}'$  such that the following inequality

$$\left(\chi_m \circ \pi_m\right)(\mathscr{K}) \ge 1 - \varepsilon \quad \forall m \in \mathbb{N}$$
(18)

holds; in this case,  $\chi'$  is uniquely determined by means of the formula  $\chi'(\mathscr{K}) = \inf_{m \in \mathbb{N}} (\chi_m \circ \pi_m)(\mathscr{K})$ , where  $\mathscr{K}$  is a compact set in  $\mathfrak{U}'$ .

Let  $K_n \in U'(n)$  be a compact set with a fixed *n*. Putting  $K_{n-1} = \pi_{n-1}^n(K_n)$ , we have

$$\chi_{n-1}(K_{n-1}) = (\chi_{n-1} \circ \pi_{n-1}^n)(K_n) = \chi_n(K_n).$$
(19)

On the other hand, if we put  $K_{n+1} = \begin{bmatrix} K_n & 0 \\ 0 & 1 \end{bmatrix}$ , then via (6),

$$\chi_{n+1}(K_{n+1}) = (\chi_n \circ \pi_n^{n+1})(K_{n+1}) = (\chi_n \circ \pi_n^{n+1}) \begin{bmatrix} K_n & 0 \\ 0 & 1 \end{bmatrix} = \chi_n(K_n).$$
(20)

As a consequence, the compact set  $\mathscr{K} = (K_m)$  in  $\mathfrak{U}'$ , generated by a compact set  $K_n \subset U'(n)$  with the help of mappings  $\pi_{n-1}^n$ , satisfies the following condition:

$$\chi_n\left(K_n\right) = \chi_m\left(K_m\right) \quad \forall m \in \mathbb{N}.$$
(21)

The probability Haar measure  $\chi_n$  is regular on U(n), and the complement  $U(n) \setminus U'(n)$  is a negligible set. Hence, if  $K_n$  runs over all compact sets in U'(n), then

$$\sup_{K_n \subset U'(n)} \chi_n\left(K_n\right) = 1.$$
(22)

Therefore, for every  $\varepsilon > 0$  there exists a compact set  $K_n \subset U'(n)$  such that  $\chi_n(K_n) \ge 1 - \varepsilon$ . From (21), it follows that for every  $\varepsilon > 0$  the compact set  $\mathscr{K}$  satisfies the hypothesis of Prohorov's criterion:

$$(\chi_m \circ \pi_m)(\mathscr{K}) = \chi_m(K_m) \ge 1 - \varepsilon \quad \forall m \in \mathbb{N}.$$
 (23)

So, in view of this criterion, there exists a unique Radon probability measure  $\chi'$  on  $\mathfrak{U}'$  which satisfies the condition (17). However, on the projective limits  $\mathfrak{U}' = \lim_{\leftarrow} U'(m)$ , there exists a unique  $U^2(\infty)$ -invariant Radon measure  $\chi$ , determined by the equality (15). Using the uniqueness property of projective limits, we obtain  $\chi' = \chi$ . The measure  $\chi$  on  $\mathfrak{U} \setminus \mathfrak{U}'$  is defined to be zero, because  $\chi_m$  is zero on  $U(m) \setminus U'(m)$ .

As a consequence of (21), we obtain (16), because

$$\chi\left(\mathscr{K}\right) = \inf_{m \in \mathbb{N}} \chi_m\left(K_m\right) = \chi_n\left(K_n\right).$$
(24)

As is known [1, Proposition 3.2], the measure  $\chi$  is  $U^2(\infty)$ invariant under the right actions (11) on the space  $\mathfrak{U}$ . Hence,
for every  $f \in \mathscr{L}^{\infty}_{\chi}$ , the equality (15) holds.

#### 4. Shift Groups

Consider that in the space  $\mathscr{L}^{\infty}_{\chi}$ , the group of shifts

$$Q_g f(u) = f(u \cdot g), \quad g \in U^2(\infty) \ u \in \mathfrak{U},$$
 (25)

is generated by the right action of  $U^2(\infty)$  over  $\mathfrak{U}$ . Choosing instead of  $U(\infty)$  a compact subgroup U(m) or the compact subgroups

$$U_{0} = \{g_{0}(\vartheta) = \exp(\mathfrak{i}\vartheta) : \vartheta \in (-\pi,\pi]\},\$$
$$U_{j}(m) = \{g_{mj}(\vartheta) = \mathbb{1}_{j-1} \otimes \exp(\mathfrak{i}\vartheta) \otimes \mathbb{1}_{m-j} : \vartheta \in (-\pi,\pi]\}$$
$$j = 1, \dots, m,$$
(26)

we obtain the corresponding subgroups of shifts  $Q_g$  with elements  $g \in U^2(m)$  or with elements  $g_0(\vartheta) \in U_0^2$  and  $g_{mj}(\vartheta) \in U_j^2(m)$ , respectively. Here,  $\mathbb{1}_m$  means the unit element in U(m).

**Lemma 2.** For any  $f \in \mathscr{L}^{\infty}_{\chi}$  the following equalities:

$$\int_{\mathfrak{U}} f d\chi = \int_{\mathfrak{U}} d\chi \left( u \right) \int_{U^{2}(m)} Q_{g} f\left( u \right) d\left( \chi_{m} \otimes \chi_{m} \right) \left( g \right), \quad (27)$$
$$\int_{\mathfrak{U}} f d\chi = \frac{1}{2\pi} \int_{\mathfrak{U}} d\chi \left( u \right) \int_{-\pi}^{\pi} Q_{g(\vartheta)} f\left( u \right) d\vartheta, \quad (28)$$

with  $g(\vartheta) \in U_0^2$  or  $U_j^2(m)$  hold.

*Proof.* For any  $f \in \mathscr{L}^{\infty}_{\chi}$ , the function  $(u, g) \mapsto Q_g f(u) = f(u \cdot g)$  is integrable on the Cartesian product  $\mathfrak{U} \times U^2(m)$ . By the Fubini theorem, we obtain

$$\int_{\mathfrak{U}} d\chi (u) \int_{U^{2}(m)} Q_{g} f(u) d(\chi_{m} \otimes \chi_{m}) (g)$$

$$= \int_{U^{2}(m)} d(\chi_{m} \otimes \chi_{m}) (g) \int_{\mathfrak{U}} Q_{g} f(u) d\chi (u).$$
(29)

This equality yields the required formula (27), because the internal integral on the right-hand side is independent of g and  $\int_{U^2(m)} d(\chi_m \otimes \chi_m) = 1$ . In turn, putting instead of U(m) the subgroups  $U_0$  and  $U_j(m)$ , we obtain equalities (28).

## 5. The Homogeneous Hilbert-Schmidt Polynomials

Consider the countable orthogonal Hilbertian sum

$$\mathsf{E} := \bigoplus_{m \in \mathbb{N}} \mathbb{C}^m = \left\{ x = (x_m) : x_m \in \mathbb{C}^m, \|x\|_{\mathsf{E}} < \infty \right\}, \quad (30)$$

with the scalar product  $\langle x \mid y \rangle_{\mathsf{E}} = \sum_{m} \langle x_{m} \mid y_{m} \rangle_{\mathbb{C}^{m}}$ , where every coordinate  $x_{m} \in \mathbb{C}^{m}$  is identified with its image  $(0, \ldots, 0, x_{m}, 0, \ldots) \in \mathsf{E}$  under the embedding  $\mathbb{C}^{m} \hookrightarrow \mathsf{E}$ .

Let  $\bigotimes_{\mathfrak{h}}^{n} \mathsf{E}$  stand for the complete *n*th tensor power of the Hilbert subspace  $\mathsf{E}$ , endowed with the Hilbertian scalar product and norm, respectively,

$$\langle x_1 \otimes \cdots \otimes x_n \mid \psi_n \rangle_{\otimes_{\mathfrak{h}}^n \mathsf{E}} = \sum_j \langle x_1 \mid y_{1j} \rangle_{\mathsf{E}} \dots \langle x_n \mid y_{nj} \rangle_{\mathsf{E}},$$

$$\| \psi_n \|_{\otimes_{\mathfrak{h}}^n \mathsf{E}} = \langle \psi_n \mid \psi_n \rangle_{\otimes_{\mathfrak{h}}^n \mathsf{E}}^{1/2},$$
(31)

where  $x_1 \otimes \cdots \otimes x_n$ ,  $y_{1j} \otimes \cdots \otimes y_{nj} \in \bigotimes_{\mathfrak{h}}^n \mathsf{E}$  with  $x_{tj}$ ,  $y_{tj} \in \mathsf{E}$  for all t = 1, ..., n and  $\psi_n = \sum_j y_{1j} \otimes \cdots \otimes y_{nj}$  denotes a finite sum. Put  $\otimes_{6}^{0} \mathsf{E} = \mathbb{C}$ . We use the following short denotation:

$$x^{\otimes n} = x \otimes \dots \otimes x, \quad x \in \mathsf{E}.$$
(32)

Replacing the space E by the subspace  $\mathbb{C}^m$ , we similarly define the tensor product  $\otimes_{\mathfrak{h}}^{n} \mathbb{C}^{m}$ . There is the unitary embedding  $\otimes_{\mathfrak{h}}^{n} \mathbb{C}^{m} \hookrightarrow \otimes_{\mathfrak{h}}^{n} \mathsf{E}$ . If m = 1, then  $\otimes_{\mathfrak{h}}^{n} \mathbb{C} = \mathbb{C}$ .

For any finite sum  $\psi_n = \sum_j y_{1j} \otimes \cdots \otimes y_{nj}$  from the space  $\otimes_{\mathfrak{h}}^{n} \mathbb{C}^{m}$  (or  $\otimes_{\mathfrak{h}}^{n} \mathsf{E}$ ), we can to define the finite type *n*-homogeneous Hilbert-Schmidt polynomials:

$$\mathbb{C}^{m} \ni x \longmapsto \left\langle x^{\otimes n} \mid \psi_{n} \right\rangle_{\otimes_{\mathfrak{h}}^{n} \mathbb{C}^{m}} = \sum_{j} \prod_{t=1}^{n} \left\langle x \mid y_{tj} \right\rangle_{\mathbb{C}^{m}}.$$
 (33)

Consider the canonical orthonormal bases:

$$\mathscr{E}(\mathbb{C}^{m}) = \{\mathbf{e}_{m1}, \dots, \mathbf{e}_{mm}\} \quad \text{in } \mathbb{C}^{m},$$
  
$$\mathscr{E}(\mathsf{E}) = \bigcup \{\mathscr{E}(\mathbb{C}^{m}) : m \in \mathbb{N}\} \quad \text{in } \mathsf{E},$$
  
(34)

where  $\mathbf{e}_{ml} = (\underbrace{0, \dots, 0, 1}^{l}, 0, \dots, 0)_{m}$ . If  $\mathfrak{G}$  :  $\{1, \dots, n\} \mapsto \{\mathfrak{F}(1), \dots, \mathfrak{F}(n)\}$  runs over all *n*elements permutations  $\mathfrak{S}(n)$ , then the symmetric *n*th tensor power  $\odot_{h}^{n} \mathbb{C}^{m}$  is defined to be a codomain of the symmetrization mapping:

$$\bigotimes_{\mathfrak{h}}^{n} \mathbb{C}^{m} \ni x_{1} \otimes \cdots \otimes x_{n} \longmapsto x_{1} \odot \cdots \odot x_{n},$$

$$x_{1} \odot \cdots \odot x_{n} := \frac{1}{n!} \sum_{\mathfrak{s} \in \mathfrak{S}(n)} x_{\mathfrak{s}(1)} \otimes \cdots \otimes x_{\mathfrak{s}(n)},$$
(35)

which is an orthogonal projector. Similarly, the symmetric *n*th tensor power  $\odot_{\mathfrak{h}}^{n} \mathsf{E}$  can be defined. Clearly,  $\odot_{\mathfrak{h}}^{n} \mathbb{C}^{m}$  is a closed subspace in  $\odot_{h}^{n} E$ .

Given a pair of numbers  $(m, n) \in \mathbb{N} \times \mathbb{Z}_+$ , we consider the *n*-fold tensor power of the canonical mapping  $\pi_m : \mathfrak{U} \ni$  $u \mapsto \pi_m(u) \in U(m),$ 

$$\mathfrak{U} \ni u \longmapsto \pi_m^{\otimes n}(u) \in \mathscr{L}(\mathfrak{O}_{\mathfrak{h}}^n \mathbb{C}^m), \qquad (36)$$

where  $\pi_m^{\otimes n}(u) := \underline{\pi_m(u) \otimes \cdots \otimes \pi_m(u)}_n$ . If n = 0, we put  $\pi_m^{\otimes 0}(u) = 1$  for all  $u \in \mathfrak{U}$  and  $m \in \mathbb{N}$ . The mapping (36) is Borel and has a continuous restriction to  $\mathfrak{U}'$ , because  $\pi_m$  has the same property (see Section 2).

Let  $\mathbf{a}_m \in \mathbb{C}^m$  be an arbitrary fixed element such that  $\|\mathbf{a}_m\|_{\mathbb{C}^m} = 1$ . Then,  $\mathbf{a}_m^{\otimes n} \in \odot_{\mathfrak{h}}^n \mathbb{C}^m$ . Using the mapping (36), we can write

$$\left[\pi_m^{\otimes n}\left(u\right)\right]\left(\mathfrak{a}_m^{\otimes n}\right) = \underbrace{\left[\pi_m\left(u\right)\right]\left(\mathfrak{a}_m\right) \otimes \cdots \otimes \left[\pi_m\left(u\right)\right]\left(\mathfrak{a}_m\right)}_{n}.$$
(37)

To any *n*-homogeneous Hilbert-Schmidt polynomial (33), there corresponds the function

$$\psi_{n}^{*}(u) := \left\langle \left[\pi_{m}^{\otimes n}(u)\right]\left(\mathfrak{a}_{m}^{\otimes n}\right) \mid \psi_{n}\right\rangle_{\otimes_{\mathfrak{h}}^{\mathfrak{n}}\mathbb{C}^{m}}$$

$$= \sum_{j} \prod_{t=1}^{n} \left\langle \left[\pi_{m}(u)\right]\left(\mathfrak{a}_{m}\right) \mid y_{tj}\right\rangle_{\mathbb{C}^{m}}$$
(38)

of the variable  $u \in \mathfrak{U}$ . Any cylindrical function of the form  $\mathfrak{U} \ni u \mapsto \langle [\pi_m(u)](\mathfrak{a}_m) \mid y_{tj} \rangle_{\mathbb{C}^m}$  has a continuous bounded restriction to  $\mathfrak{U}'$ . Therefore, it is  $\chi$ -essentially bounded on  $\mathfrak{U}$ , because  $\mathfrak{U} \setminus \mathfrak{U}'$  is a  $\chi$ -negligible set. Consequently,  $\psi_n^* \in L_{\chi}^{\infty}$ and  $\psi_n^*|_{\mathfrak{U}'}$  is continuous and bounded.

*Definition 3.* We define  $\mathscr{P}^n_{\mathfrak{h}}(\mathbb{C}^m)$  to be the space of all functions  $\psi_n^*$  of the variable  $u \in \mathfrak{U}$ , determined by the finite type *n*-homogeneous Hilbert-Schmidt polynomials (33).

**Lemma 4.** For any element  $\mathfrak{a}_m \in \mathbb{C}^m$  such that  $\|\mathfrak{a}_m\|_{\mathbb{C}^m} = 1$ the set

$$\mathbf{S}^{m} = \left\{ x = \left[ \pi_{m}\left( u \right) \right] \left( \mathfrak{a}_{m} \right) : u \in \mathfrak{U} \right\}$$
(39)

coincides with the unit sphere in  $\mathbb{C}^m$ . As a consequence, the one-to-one antilinear corresponding

$$\odot^{n}_{\mathfrak{h}}\mathbb{C}^{m} \ni \psi_{n} \rightleftharpoons \psi_{n}^{*} \in \mathscr{P}^{n}_{\mathfrak{h}}(\mathbb{C}^{m}).$$

$$(40)$$

Holds, and any function  $\psi_n^*$  is independent of the choice of an element  $\mathfrak{a}_m \in S^m$ .

*Proof.* Suppose, on the contrary, that there is an element  $\psi_n \in$  $\odot_{\mathfrak{h}}^{n}\mathbb{C}^{m}$  such that  $\langle x^{\otimes n} | \psi_{n} \rangle_{\otimes_{\mathfrak{h}}^{n}\mathbb{C}^{m}} = 0$  for all  $x = [\pi_{m}(u)](\mathfrak{a}_{m}) \in \mathbb{C}^{m}$  $S^m$  with  $u \in \mathfrak{U}$ . The mapping

$$\pi_m: \mathfrak{U} \ni u \longmapsto \pi_m(u) \in U(m) \tag{41}$$

is surjective by surjectivity of the mapping  $\pi_m$  (see [2, Lemma 3.1]). Hence, the set  $S^m$  coincides with the unit sphere in  $\mathbb{C}^m$  and is independent on the choice of an element  $\mathfrak{a}_m$ . By *n*-homogeneity, we have  $\langle x^{\otimes n} | \psi_n \rangle_{\otimes_{\kappa}^n \mathbb{C}^m} = 0$  for all  $x \in \mathbb{C}^m$ .

Apply the following polarization formula for symmetric tensor products (see, e.g., [10, Section 1.5]):

$$z_1 \odot \cdots \odot z_n = \frac{1}{2^n n!} \sum_{1 \le t \le n} \sum_{\delta_t = \pm 1} \delta_1 \cdots \delta_n \ x^{\otimes n}, \qquad (42)$$

with  $x = \sum_{t=1}^{n} \delta_t z_t \in \mathbb{C}^m$ , which is valid for all  $z_1, \ldots, z_n \in \mathbb{C}^m$ . It follows that  $\langle z_1 \odot \cdots \odot z_n | \psi_n \rangle_{\otimes_h^n \mathbb{C}^m} = 0$  for all elements  $z_1, \ldots, z_n \in \mathbb{C}^m$ . Hence,  $\psi_n = 0$ , because the subset of all elements  $z_1 \odot \cdots \odot z_n$  is total in  $\odot_{\mathfrak{h}}^n \mathbb{C}^m$ . As a consequence, the subset

$$\left\{x^{\otimes n} = \left[\pi_m^{\otimes n}\left(u\right)\right]\left(\mathfrak{a}_m^{\otimes n}\right) : u \in \mathfrak{U}\right\}$$
(43)

is also total in  $\odot_{h}^{n} \mathbb{C}^{m}$ . It immediately yields the correspondence (40). 

Consider the symmetric Fock space F and its closed subspace  $F_m$ , where

$$\begin{aligned} \mathsf{F} &:= \mathbb{C} \oplus \mathsf{E} \oplus \left( \odot^2_{\mathfrak{h}} \mathsf{E} \right) \oplus \left( \odot^3_{\mathfrak{h}} \mathsf{E} \right) \oplus \cdots , \\ \mathsf{F}_m &:= \mathbb{C} \oplus \mathbb{C}^m \oplus \left( \odot^2_{\mathfrak{h}} \mathbb{C}^m \right) \oplus \left( \odot^3_{\mathfrak{h}} \mathbb{C}^m \right) \oplus \cdots . \end{aligned}$$

We will use the following notations:

$$(m) := (m1, ..., mm),$$

$$k_{(m)} := (k_{m1}, ..., k_{mm}) \in \mathbb{Z}_{+}^{m},$$

$$|k_{(m)}| := k_{m1} + \dots + k_{mm},$$

$$k_{(m)}! := k_{m1}! \cdot \ldots \cdot k_{mm}!.$$
(45)

As is well known (see, e.g., [11]), the system of symmetric tensor elements, indexed by the set  $k_{(m)}$ ,

$$\mathscr{E}\left(\odot_{\mathfrak{f}}^{n}\mathbb{C}^{m}\right) = \left\{ \mathfrak{e}_{(m)}^{\otimes k_{(m)}} = \mathfrak{e}_{m1}^{\otimes k_{m1}} \odot \cdots \odot \mathfrak{e}_{mm}^{\otimes k_{mm}} : \\ k_{(m)} \in \mathbb{Z}_{+}^{m}; \ \left|k_{(m)}\right| = n \right\}$$

$$(46)$$

forms an orthogonal basis in the subspace

$$\odot^{n}_{\mathfrak{h}}\mathbb{C}^{m}\subset\mathsf{F}_{m}.$$
(47)

\*\*\*\*\*\*

(111)

We will also use the following notations:

$$[m] := \{(11), (21, 22), \dots, (m1, \dots, mm)\},$$

$$\{k\} := \{k_{(1)}, \dots, k_{(m)}\} \in \bigvee_{r=1}^{m} \mathbb{Z}_{+}^{r},$$

$$|\{k\}| := |k_{(1)}| + \dots + |k_{(m)}|,$$

$$\{k\}! := k_{(1)}! \cdot \dots \cdot k_{(m)}!.$$
(48)

Then, the system of symmetric tensor elements with a fixed n, indexed by the sets [m] and  $\{k\}$ ,

$$\mathscr{E}_{n} = \bigcup_{m \in \mathbb{N}} \left\{ e_{[m]}^{\otimes \{k\}} = e_{(1)}^{\otimes k_{(1)}} \odot \cdots \odot e_{(m)}^{\otimes k_{(m)}} : \\ e_{(1)}^{\otimes k_{(1)}} \in \mathscr{E}\left( \odot_{\mathfrak{h}}^{|k_{(1)}|} \mathbb{C} \right), \dots, e_{(m)}^{\otimes k_{(m)}} \in \mathscr{E}\left( \odot_{\mathfrak{h}}^{|k_{(m)}|} \mathbb{C}^{m} \right) \\ \text{with fixed } |\{k\}| = n \right\},$$

$$(49)$$

forms an orthogonal basis in the subspace  $\bigcirc_{h}^{n} \mathsf{E} \subset \mathsf{F}$ . Thus, the system

$$\mathscr{E} = \left\{ \mathscr{E}_n : n \in \mathbb{Z}_+ \right\} \tag{50}$$

forms an orthogonal basis in the symmetric Fock space F.

By virtue of the one-to-one mapping (40), the system of symmetric tensor elements  $\mathscr{E}(\odot_{\mathfrak{h}}^{n}\mathbb{C}^{\overline{m}})$  uniquely defines the following corresponding system:

$$\mathscr{E}_{m,n}^* \in \mathscr{P}_{\mathfrak{h}}^n(\mathbb{C}^m), \qquad (51)$$

of the following  $\chi_m$ -integrable cylindrical functions:

$$\mathbf{e}_{(m)}^{*k_{(m)}}(u) := \left\langle \left[\pi_{m}^{\otimes n}(u)\right]\left(\mathbf{e}_{m1}^{\otimes n}\right) \mid \mathbf{e}_{(m)}^{\otimes k_{(m)}}\right\rangle_{\bigotimes_{\mathfrak{g}}^{n}\mathbb{C}^{m}}$$

$$= \prod_{r=1}^{m} \left\langle \left(\pi_{m} \circ u\right)\left(\mathbf{e}_{m1}\right) \mid \mathbf{e}_{mr}\right\rangle_{\mathbb{C}^{m}}^{k_{mr}},$$
(52)

of the variable  $u \in \mathfrak{U}$ , where we take  $\mathfrak{a}_m = \mathfrak{e}_{m1}$ . Consider the system of functions of the variable  $u \in \mathfrak{U}$ ,

$$\mathscr{C}_{n}^{*} = \bigcup_{m \in \mathbb{N}} \left\{ \mathbf{e}_{[m]}^{*\{k\}} = \mathbf{e}_{(1)}^{*k_{(1)}} \cdot \cdots \cdot \mathbf{e}_{(m)}^{*k_{(m)}} : \\ \mathbf{e}_{(1)}^{*k_{(1)}} \in \mathscr{C}_{1,|k_{(1)}|}^{*}, \dots, \mathbf{e}_{(m)}^{*k_{(m)}} \in \mathscr{C}_{m,|k_{(m)}|}^{*} \right\}$$
with fixed  $|\{k\}| = n$ , (53)

generated by the system of symmetric tensor elements  $\mathscr{C}_n$ . All these functions belong to the space  $\mathscr{L}^{\infty}_{\chi}$  by their definition. Denote

$$\mathscr{C}^* = \left\{ \mathscr{C}_n^* : n \in \mathbb{Z}_+ \right\}, \qquad \mathscr{C}_m^* = \left\{ \mathscr{C}_{m,n}^* : n \in \mathbb{Z}_+ \right\}.$$
(54)

## 6. The Hardy-Type Space

Let  $L_{\chi}^2$  be the space of square  $\chi$ -integrable complex functions , f on the space of virtual matrices  ${\mathfrak U}.$  Since  $\chi$  is a probability measure, the embedding  $\mathscr{L}^{\infty}_{\chi} \subset L^{2}_{\chi}$  holds and

$$\|f\|_{L^{2}_{\chi}} \leq \operatorname{ess\,sup}_{u \in \mathfrak{U}} |f(u)|, \quad f \in \mathscr{L}^{\infty}_{\chi}.$$
(55)

Denote by  $\mathscr{H}^2_{\chi_m}$  the  $L^2_{\chi}$ -closure of complex linear spans of the subsystem  $\mathscr{C}^*_m$ . As is well known (see, e.g., [12, Theorem 5.6.8]), the space  $\mathscr{H}^2_{\chi_m}$  is isomorphic to the classic Hardy space  $\mathscr{H}^{2}_{\chi_{m}}(\mathsf{B}^{m})$  of analytic complex functions on the open unit ball  $B^m = \{x_m \in \mathbb{C}^m : \|x_m\|_{\mathbb{C}^m} < 1\}$ . Therefore, the following more general definition seems natural (see, also [8]).

Definition 5. The Hardy-type space  $\mathscr{H}^2_{\chi}$  on the space of virtual unitary matrices  $\mathfrak{U}$  is defined to be the  $L^2_{\gamma}$ -closure of the complex linear span of the system  $\,\mathscr{C}^*.$ 

**Theorem 6.** The system  $\mathscr{C}^*$  of all functions  $\mathbf{e}_{[m]}^{*\{k\}}$  $\mathbf{e}_{(1)}^{*k_{(1)}} \cdots \mathbf{e}_{(m)}^{*k_{(m)}}$  with  $m \in \mathbb{N}$ , such that  $\mathbf{e}_{(r)}^{*k_{(r)}} \in \mathscr{E}_{r,|k_{(r)}|}^{*}$  as  $r = 1, \ldots, m$ , forms an orthogonal basis in the Hardy-type spaces  $\mathcal{H}_{\chi}^{2}$  with norms

$$\left\| \mathbf{e}_{[m]}^{*\{k\}} \right\|_{L^{2}_{\chi}} = \left( \prod_{r=1}^{m} \frac{(r-1)!(k)_{r}!}{(r-1+\left|(k)_{r}\right|)!} \right)^{1/2}.$$
 (56)

*Proof.* If  $|\{k\}| \neq |\{q\}|$ , then from (28), it follows that

$$\int_{\mathfrak{U}} \mathbf{e}_{[m]}^{*\{k\}} \cdot \overline{\mathbf{e}}_{[n]}^{*\{q\}} d\chi$$

$$= \int_{\mathfrak{U}} \mathbf{e}_{[m]}^{*\{k\}} (\exp(\mathfrak{i}\vartheta) u) \cdot \overline{\mathbf{e}}_{[n]}^{*\{q\}} (\exp(\mathfrak{i}\vartheta) u) d\chi (u)$$

$$= \frac{1}{2\pi} \int_{\mathfrak{U}} \mathbf{e}_{[m]}^{*\{k\}} \overline{\mathbf{e}}_{[n]}^{*\{q\}} d\chi \int_{-\pi}^{\pi} \exp\left(\mathfrak{i}\left(|\{k\}| - |\{q\}|\right)\vartheta\right) d\vartheta$$

$$= 0.$$
(57)

So,  $\mathbf{e}_{[m]}^{*\{k\}} \perp \mathbf{e}_{[n]}^{*\{q\}}$  in the space  $L_{\chi}^2$  if  $|\{k\}| \neq |\{q\}|$  for all indices [m], [n].

Let  $|\{k\}| = |\{q\}|$  and m > n for definiteness. If the elements  $\mathbf{e}_{[m]}^{*\{k\}}$  and  $\mathbf{e}_{[n]}^{*\{q\}}$  are different, then there exists a subindex  $ms \in \{11, 21, 22, ..., m1, ..., mm\}$  in the blockindex [m] = [(11), (21, 22), ..., (m1, ..., mm)] such that  $ms \notin$  $\{11, 21, 22, \dots, n1, \dots, nn\}$ , where  $[n] = [(11), (21, 22), \dots, nn]$  $(n1, \ldots, nn)$ ]. The formula (28) implies that for the group of shifts  $Q_{g_{ms}(\vartheta)}$  generated by elements  $g_{ms}(\vartheta) \in U_s^2(m)$  with the subindex ms,

$$\int_{\mathfrak{U}} \mathbf{e}_{[m]}^{*\{k\}} \cdot \mathbf{\bar{e}}_{[n]}^{*\{q\}} d\chi$$

$$= \int_{\mathfrak{U}} Q_{g_{ms}(\vartheta)} \mathbf{e}_{[m]}^{*\{k\}} \cdot Q_{g_{ms}(\vartheta)} \mathbf{\bar{e}}_{[n]}^{*\{q\}} d\chi \qquad (58)$$

$$= \frac{1}{2\pi} \int_{\mathfrak{U}} \mathbf{e}_{[m]}^{*\{k\}} \cdot \mathbf{\bar{e}}_{[n]}^{*\{q\}} d\chi \int_{-\pi}^{\pi} \exp\left(\mathbf{i}k_{ms}\vartheta\right) d\vartheta = 0.$$

Hence,  $\mathbf{e}_{[m]}^{*\{k\}} \perp \mathbf{e}_{[n]}^{*\{q\}}$  in  $L^2_{\chi}$ .

Let now  $|\{k\}| = |\{q\}|$  and m = n. If  $e_{[m]}^{*\{k\}} \neq e_{[n]}^{*\{q\}}$ , then  $\{k\} \neq \{q\}$ . Hence, there exists a sub-index rs in the blockindex [m] = [n] such that  $k_{rs} \neq q_{rs}$ . Similarly as previous mentioned, applying the formula (28) to the group of shifts  $Q_{q_{rs}(\vartheta)}$  generated by elements  $g_{rs}(\vartheta) \in U_s^2(r)$  with the subindex rs, we get

$$\int_{\mathfrak{U}} \mathbf{e}_{[m]}^{*\{k\}} \cdot \overline{\mathbf{e}}_{[n]}^{*\{q\}} d\chi$$

$$= \frac{1}{2\pi} \int_{\mathfrak{U}} \mathbf{e}_{[m]}^{*\{k\}} \overline{\mathbf{e}}_{[n]}^{*\{q\}} d\chi \int_{-\pi}^{\pi} \exp\left(\mathbf{i} \left(k_{rs} - q_{rs}\right) \vartheta\right) d\vartheta$$

$$= 0.$$
(59)

Hence, in this case also  $e_{[m]}^{*\{k\}} \perp e_{[n]}^{*\{q\}}$  under the measure  $\chi$ . Let  $g_r = (\mathbb{1}_r, w_r) \in U^2(r)$  and  $u \in \mathfrak{U}$ . Using (11) and (52),

we have

$$\begin{split} \int_{U^{2}(r)} Q_{g_{r}} \left| \mathbf{e}_{(r)}^{*(k)_{r}} \right|^{2}(u) d\left(\chi_{r} \otimes \chi_{r}\right) \left(g_{r}\right) \\ &= \int_{U(r)} \prod_{l=1}^{r} \left| \left\langle \left[w_{r}^{-1} \pi_{r}\left(u\right)\right] \left(\mathbf{e}_{r1}\right) \mid \mathbf{e}_{rl} \right\rangle_{\mathbb{C}^{r}}^{k_{rl}} \right|^{2} d\chi_{r}\left(w_{r}\right). \end{split}$$

$$\tag{60}$$

However, the previous integral with the Haar measure  $\chi_r$  is independent of  $\pi_r(u) \in U(r)$ . It follows that

$$\begin{aligned} \int_{U^{2}(r)} Q_{g_{r}} \left| \mathbf{e}_{(r)}^{*(k)_{r}} \right|^{2}(u) d\left(\chi_{r} \otimes \chi_{r}\right) \left(g_{r}\right) \\ &= \int_{U(r)} \prod_{l=1}^{r} \left| \left\langle w_{r}^{-1}\left(\mathbf{e}_{r1}\right) \mid \mathbf{e}_{rl} \right\rangle_{\mathbb{C}^{r}}^{k_{rl}} \right|^{2} d\chi_{r}\left(w_{r}\right) \qquad (61) \\ &= \frac{(r-1)!(k)_{r}!}{(r-1+|(k)_{r}|)!} = \left\| \mathbf{e}_{(r)}^{*(k)_{r}} \right\|_{L^{2}_{\chi_{r}}}^{2} \end{aligned}$$

by the well-known formula [12, Section 1.4.9]. Using the formula (27) *m*-times for  $r = 1, \ldots, m$ , we get

$$\begin{aligned} \int_{\mathfrak{U}} \left| \mathbf{e}_{[m]}^{*\{k\}} \right|^{2} d\chi \\ &= \int_{\mathfrak{U}} d\chi \left( u \right) \prod_{r=1}^{m} \int_{U^{2}(r)} Q_{g_{r}} \left| \mathbf{e}_{(r)}^{*(k)_{r}} \right|^{2} \left( u \right) d\left( \chi_{r} \otimes \chi_{r} \right) \left( g_{r} \right) \\ &= \prod_{r=1}^{m} \left\| \mathbf{e}_{(r)}^{*k_{(r)}} \right\|_{L^{2}_{\chi_{r}}}^{2}, \end{aligned}$$
(62)

because  $\int_{M} d\chi = 1$ . It follows that

$$\left\|\mathbf{e}_{[m]}^{*\{k\}}\right\|_{L^{2}_{\chi}}^{2} = \prod_{r=1}^{m} \left\|\mathbf{e}_{(r)}^{*k_{(r)}}\right\|_{L^{2}_{\chi_{r}}}^{2} = \prod_{r=1}^{m} \frac{(r-1)!(k)_{r}!}{(r-1+|(k)_{r}|)!}, \quad (63)$$

for all 
$$e_{[m]}^{*\{k\}} = e_{(1)}^{*k_{(m)}} \cdot \cdots \cdot e_{(m)}^{*k_{(m)}}$$
.

As is known (see, e.g., [11]), the system  $\mathcal{C}_m$  of symmetric tensors  $e_{(m)}^{\otimes(k)_m}$  with a fixed *m* forms an orthogonal basis in the symmetric Fock space  $F_m$  with norms  $\|\mathbf{e}_{(m)}^{\otimes(k)_m}\|_{F_m} =$  $\sqrt{(k)_m!/|(k)_m|!}$ . Similarly, the system  $\mathscr{E}$  of symmetric tensors  $\mathbf{e}_{[m]}^{\otimes \{k\}} = \mathbf{e}_{(1)}^{\otimes (k)_1} \odot \cdots \odot \mathbf{e}_{(m)}^{\otimes k_{(m)}}$  with all  $m \in \mathbb{N}$ , such that  $\mathbf{e}_{(r)}^{\otimes (k)_r} \in \mathscr{E}_{r,|(k)_r|}$  as  $r = 1, \ldots, m$ , forms an orthogonal basis in the symmetric Fock space F with norms  $\|\mathbf{e}_{[m]}^{\otimes\{k\}}\|_{\mathsf{F}} = \sqrt{\{k\}!/|\{k\}|!}$ .

Combining Lemma 4, Theorem 6, and [12, Theorem 5.6.8], we obtain the following.

Theorem 7. Antilinear extensions of the one-to-one mappings between the orthonormal bases

$$\frac{\mathbf{e}_{(m)}^{\otimes(k)_{m}}}{\left\|\mathbf{e}_{(m)}^{\otimes(k)_{m}}\right\|_{\mathsf{F}_{m}}} \rightleftharpoons \frac{\mathbf{e}_{(m)}^{*(k)_{m}}}{\left\|\mathbf{e}_{(m)}^{*(k)_{m}}\right\|_{L^{2}_{\chi_{m}}}},$$

$$\frac{\mathbf{e}_{[m]}^{\otimes\{k\}}}{\left\|\mathbf{e}_{[m]}^{\otimes\{k\}}\right\|_{\mathsf{F}}} \rightleftharpoons \frac{\mathbf{e}_{[m]}^{*\{k\}}}{\left\|\mathbf{e}_{[m]}^{*\{k\}}\right\|_{L^{2}_{\chi}}},$$
(64)

uniquely define the corresponding anti-linear isometric isomorphisms

$$\mathsf{F}_m \simeq \mathscr{H}^2_{\chi_m} \left( \mathsf{B}^m \right), \qquad \mathsf{F} \simeq \mathscr{H}^2_{\chi}.$$
 (65)

Reasoning by analogy with [8, Proposition 6.1 and Theorem 7.1], it is easy to show that the Hardy space  $\mathscr{H}_{\gamma}^{2}$  possesses the reproducing kernel of a Cauchy type

$$\mathfrak{C}(v,u) = \sum_{n \in \mathbb{Z}_{+}} \sum_{|\{k\}|=n} \frac{\mathbf{e}_{[m]}^{*\{k\}}(v) \,\overline{\mathbf{e}}_{[m]}^{*\{k\}}(u)}{\left\|\mathbf{e}_{[m]}^{*\{k\}}\right\|_{L^{2}_{\chi}}^{2}} \\ = \prod_{m=1}^{\infty} (1 - \left\langle \left(\pi_{m} \circ v\right) \left(\mathbf{e}_{m1}\right) \right| \left(\pi_{m} \circ u\right) \left(\mathbf{e}_{m1}\right) \right\rangle_{\mathsf{E}} \right)^{-m},$$
(66)

with  $u, v \in \mathfrak{U}$ , where the sum  $\sum_{|\{k\}|=n}$  is over all indices  $\{k\} \in \{\bigotimes_{r=1}^{m} \mathbb{Z}_{+}^{r} : m \in \mathbb{N}\}$  such that  $|\{k\}| = n$ . As a consequence, the integral representation of any function  $f \in \mathcal{H}_{\gamma}^{2}$ ,

$$f(\lambda v) = \int_{\mathfrak{U}} f(u) \mathfrak{C}(\lambda v, u) d\chi(u)$$
(67)

gives a unique analytic extension in the complex variable  $\lambda \in B^1$  for all elements  $v \in \mathfrak{U}$  such that

$$\sum_{m\in\mathbb{N}} m \left\| \left( \pi_m \circ \nu \right) \left( \mathbf{e}_{m1} \right) \right\|_{\mathbb{C}^m}^2 < \infty.$$
(68)

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