# QUASI-EINSTEIN TOTALLY REAL SUBMANIFOLDS OF THE NEARLY KÄHLER 6-SPHERE 

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#### Abstract

We investigate Lagrangian submanifolds of the nearly Kähler 6-sphere. In particular we investigate Lagrangian quasi-Einstein submanifolds of the 6 -sphere. We relate this class of submanifolds to certain tubes around almost complex curves in the 6 -sphere.


1. Introduction. In this paper, we investigate 3 -dimensional totally real submanifolds $M^{3}$ of the nearly Kähler 6-sphere $S^{6}$. A submanifold $M^{3}$ of $S^{6}$ is called totally real if the almost complex structure $J$ on $S^{6}$ interchanges the tangent and the normal space. It has been proven by Ejiri ([E1]) that such submanifolds are always minimal and orientable. In the same paper, he also classified those totally real submanifolds with constant sectional curvature. Note that 3-dimensional Einstein manifolds have constant sectional curvature. Here, we will investigate the totally real submanifolds of $S^{6}$ for which the Ricci tensor has an eigenvalue with multiplicity at least 2 . In general, a manifold $M^{n}$ whose Ricci tensor has an eigenvalue of multiplicity at least $n-1$ is called quasi-Einstein.

The paper is organized as follows. In Section 2, we recall the basic formulas about the vector cross product on $\boldsymbol{R}^{7}$ and the almost complex structure on $S^{6}$. We also relate the standard Sasakian structure on $S^{5}$ with the almost complex structure on $S^{6}$. Then, in Section 3, we derive a necessary and sufficient condition for a totally real submanifold of $S^{6}$ to be quasi-Einstein. Using this condition, we deduce from [C], see also [CDVV1] and [DV], that totally real submanifolds $M$ with $\delta_{M}=2$ are quasi-Einstein. Here, $\delta_{M}$ is the Riemannian invariant defined by

$$
\delta_{M}(p)=\tau(p)-\inf K(p)
$$

where $\inf K$ is the function assigning to each $p \in M$ the infimum of $K(\pi), K(\pi)$ denoting the sectional curvature of a 2-plane $\pi$ of $T_{p} M$, where $\pi$ runs over all 2-planes in $T_{p} M$ and $\tau$ is the scalar curvature of $M$ defined by $\tau=\sum_{i<j} K\left(e_{i} \wedge e_{j}\right)$. Totally real submanifolds of $S^{6}$ with $\delta_{M}=2$ have been classified in [DV]. Essentially, these submanifolds are either local lifts of holomorphic curves in $\boldsymbol{C} P^{2}$ or tubes with radius $\pi / 2$ in the direction of $N N^{2}$, where $N^{2}$ is a non-totally geodesic almost complex curve and $N N^{2}$ denotes the vector bundle whose fibres are planes orthogonal to the first osculating space of $N^{2}$.

[^0]In Section 4, we then construct some other examples of 3-dimensional quasi-Einstein totally real submanifolds by considering also tubes with different radii. More specifically, we prove

THEOREM 1. Let $\phi: N^{2} \rightarrow S^{6}$ be an almost complex curve in $S^{6}(1)$ without totally geodesic points. Denote by $U N^{2}$ the unit tangent bundle of $N^{2}$. Define

$$
\psi_{\gamma}: U N^{2} \rightarrow S^{6}: v \mapsto \cos \gamma \phi+\sin \gamma v \times \frac{\alpha(v, v)}{\|\alpha(v, v)\|}
$$

where $\alpha$ denotes the second fundamental form of the surface $N^{2}$. Then $\psi$ is an immersion on an open dense subset of $U N^{2}$. Moreover $\psi$ is totally real if and only if either
(1) $\gamma=\pi / 2$, or
(2) $\cos ^{2} \gamma=5 / 9$ and $N^{2}$ is a superminimal surface.

Further, in both cases the immersion defines a quasi-Einstein metric on $U N^{2}$ and if (1) holds, then with respect to this metric $\delta_{U N^{2}}=2$ and if (2) holds, then $\delta_{U N^{2}}<2$.

The above theorem also generalizes results obtained by Ejiri [E2], who only considered tubes around superminimal almost complex curves and who omitted Case (1).

Next, in Section 5, we prove the following converse:
THEOREM 2. Let $F: M^{3} \rightarrow S^{6}$ be a totally real immersion of a 3-dimensional quasiEinstein manifold. Then either $\delta_{M^{3}}=2$ or there exists an open dense subset $W$ of $M$ such that each point $p$ of $W$ has a neighborhood $W_{1}$ such that either
(1) $F\left(W_{1}\right)=\psi_{\gamma}\left(U N^{2}\right)$, where $N^{2}$ is a superminimal linearly full almost complex curve in $S^{6}$, and $\psi_{\gamma}$ with $\cos ^{2} \gamma=5 / 9$ is as defined in Theorem 1 , or
(2) $\quad F\left(W_{1}\right)$ is an open subset of $\tilde{\psi}\left(S^{3}\right)$, where $\tilde{\psi}$ is as defined in Section 4.

Case (2) can be considered as a limit case of the previous one, by taking for $N^{2}$ a totally geodesic almost complex curve. Note also that Theorem 2 together with the classification theorems of [DV] provides a complete classification of the totally real quasi-Einstein submanifolds of $S^{6}$.
2. The vector cross product and the almost complex structure on $S^{6}$. We give a brief exposition of how the standard nearly Kähler structure on $S^{6}$ arises in a natural manner from the Cayley multiplication. For further details about the Cayley numbers and their automorphism group $G_{2}$, we refer the reader to [W] and [HL].

The multiplication on the Cayley numbers $\mathcal{O}$ may be used to define a vector cross product on the purely imaginary Cayley numbers $\boldsymbol{R}^{7}$ using the formula

$$
\begin{equation*}
u \times v=(1 / 2)(u v-v u) \tag{2.1}
\end{equation*}
$$

while the standard inner product on $\boldsymbol{R}^{7}$ is given by

$$
\begin{equation*}
(u, v)=-(1 / 2)(u v+v u) . \tag{2.2}
\end{equation*}
$$

It is now elementary to show that

$$
\begin{equation*}
u \times(v \times w)+(u \times v) \times w=2(u, w) v-(u, v) w-(w, v) u, \tag{2.3}
\end{equation*}
$$

and that the triple scalar product $(u \times v, w)$ is skew symmetric in $u, v, w$, see [HL] for proofs.
Conversely, denoting by $\operatorname{Re}(\mathcal{O})$ (respectively $\operatorname{Im}(\mathcal{O})$ ) the real (respectively imaginary) Cayley numbers, the Cayley multiplication on $\mathcal{O}$ is given in terms of the vector cross product and the inner product by

$$
\begin{align*}
& (r+u)(s+v)=r s-(u, v)+r v+s u+(u \times v),  \tag{2.4}\\
& r, s \in \operatorname{Re}(\mathcal{O}), u, v \in \operatorname{Im}(\mathcal{O}) .
\end{align*}
$$

In view of (2.1), (2.2) and (2.4), it is clear that the group $G_{2}$ of automorphisms of $\mathcal{O}$ is precisely the group of isometries of $\boldsymbol{R}^{7}$ preserving the vector cross product.

An ordered orthonormal basis $e_{1}, \ldots, e_{7}$ is said to be a $G_{2}$-frame if

$$
\begin{equation*}
e_{3}=e_{1} \times e_{2}, \quad e_{5}=e_{1} \times e_{4}, \quad e_{6}=e_{2} \times e_{4}, \quad e_{7}=e_{3} \times e_{4} . \tag{2.5}
\end{equation*}
$$

For example, the standard basis of $\boldsymbol{R}^{7}$ is a $G_{2}$-frame. Moreover, if $e_{1}, e_{2}, e_{4}$ are mutually orthogonal unit vectors with $e_{4}$ orthogonal to $e_{1} \times e_{2}$, then $e_{1}, e_{2}, e_{4}$ determine a unique $G_{2}$-frame $e_{1}, \ldots, e_{7}$ and $\left(\boldsymbol{R}^{7}, \times\right)$ is generated by $e_{1}, e_{2}, e_{4}$ subject to the relations

$$
\begin{equation*}
e_{i} \times\left(e_{j} \times e_{k}\right)+\left(e_{i} \times e_{j}\right) \times e_{k}=2 \delta_{i k} e_{j}-\delta_{i j} e_{k}-\delta_{j k} e_{i} \tag{2.6}
\end{equation*}
$$

Therefore, for any $G_{2}$-frame, we have the following multiplication table:

| $\times$ | $e_{1}$ | $e_{2}$ | $e_{3}$ | $e_{4}$ | $e_{5}$ | $e_{6}$ | $e_{7}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $e_{1}$ | 0 | $e_{3}$ | $-e_{2}$ | $e_{5}$ | $-e_{4}$ | $-e_{7}$ | $e_{6}$ |
| $e_{2}$ | $-e_{3}$ | 0 | $e_{1}$ | $e_{6}$ | $e_{7}$ | $-e_{4}$ | $-e_{5}$ |
| $e_{3}$ | $e_{2}$ | $-e_{1}$ | 0 | $e_{7}$ | $-e_{6}$ | $e_{5}$ | $-e_{4}$ |
| $e_{4}$ | $-e_{5}$ | $-e_{6}$ | $-e_{7}$ | 0 | $e_{1}$ | $e_{2}$ | $e_{3}$ |
| $e_{5}$ | $e_{4}$ | $-e_{7}$ | $e_{6}$ | $-e_{1}$ | 0 | $-e_{3}$ | $e_{2}$ |
| $e_{6}$ | $e_{7}$ | $e_{4}$ | $-e_{5}$ | $-e_{2}$ | $e_{3}$ | 0 | $-e_{1}$ |
| $e_{7}$ | $-e_{6}$ | $e_{5}$ | $e_{4}$ | $-e_{3}$ | $-e_{2}$ | $e_{1}$ | 0 |

The standard nearly Kähler structure on $S^{6}$ is then obtained as follows.

$$
J u=x \times u, \quad u \in T_{x} S^{6}, x \in S^{6} .
$$

It is clear that $J$ is an orthogonal almost complex structure on $S^{6}$. In fact, $J$ is a nearly Kähler structure in the sense that the $(2,1)$-tensor field $G$ on $S^{6}(1)$ defined by

$$
G(X, Y)=\left(\tilde{\nabla}_{X} J\right)(Y)
$$

where $\tilde{\nabla}$ is the Levi-Civita connection on $S^{6}(1)$, is skew-symmetric. A straightforward computation also shows that

$$
G(X, Y)=X \times Y-\langle x \times X, Y\rangle x .
$$

For more information on the properties of the Cayley multiplication, $J$ and $G$, we refer to [Ca2], [BVW] and [DVV].

Finally, we recall the explicit construction of a Sasakian structure on $S^{5}(1)$ starting from $C^{3}$ and its relation with the nearly Kähler structure on $S^{6}$. For more details about Sasakian
structures we refer the reader to [B]. We consider $S^{5}$ as the hypersphere in $S^{6} \subset \boldsymbol{R}^{7}$ given by the equation $x_{4}=0$ and define

$$
j: S^{5}(1) \rightarrow C^{3}:\left(x_{1}, x_{2}, x_{3}, 0, x_{5}, x_{6}, x_{7}\right) \mapsto\left(x_{1}+i x_{5}, x_{2}+i x_{6}, x_{3}+i x_{7}\right)
$$

Then at a point $p$ the structure vector field is given by

$$
\xi(p)=\left(x_{5}, x_{6}, x_{7}, 0,-x_{1},-x_{2},-x_{3}\right)=e_{4} \times p
$$

and for a tangent vector $v=\left(v_{1}, v_{2}, v_{3}, 0, v_{5}, v_{6}, v_{7}\right)$, orthogonal to $\xi$, we have

$$
\phi(v)=\left(-v_{5},-v_{6},-v_{7}, 0, v_{1}, v_{2}, v_{3}\right)=v \times e_{4} .
$$

Also, $\phi \xi(p)=0=\left(e_{4} \times p\right) \times e_{4}-\left\langle\left(e_{4} \times p\right) \times e_{4}, p\right\rangle p$, from which we deduce for any tangent vector $w$ to $S^{5}$ that

$$
\begin{equation*}
\phi(w)=w \times e_{4}-\left\langle w \times e_{4}, p\right\rangle p \tag{2.7}
\end{equation*}
$$

3. A pointwise characterization. Let $M^{3}$ be a totally real submanifold of $S^{6}$. From [E1], we know that $M^{3}$ is minimal and that for tangent vector fields $X$ and $Y, G(X, Y)$ is a normal vector field on $M^{n}$. Moreover, $\langle h(X, Y), J Z\rangle$ is symmetric in $X, Y$ and $Z$. Denote by $S$ the Ricci tensor of $M^{3}$ defined by

$$
S(Y, Z)=\operatorname{trace}\{X \mapsto R(X, Y) Z\}
$$

and denote by Ric the associated 1-1 tensor field, i.e.,

$$
\langle\operatorname{Ric}(Y), Z\rangle=S(Y, Z)
$$

Let $p \in M$ and assume that $p$ is not a totally geodesic point of $M^{3}$. Then, we know from [V] that there exists an orthonormal basis $\left\{e_{1}, e_{2}, e_{3}\right\}$ at the point $p$ such that

$$
\begin{array}{ll}
h\left(e_{1}, e_{1}\right)=\lambda_{1} J e_{1}, & h\left(e_{2}, e_{2}\right)=\lambda_{2} J e_{1}+a J e_{2}+b J e_{3}, \\
h\left(e_{1}, e_{2}\right)=\lambda_{2} J e_{2}, & h\left(e_{2}, e_{3}\right)=b J e_{2}-a J e_{3}, \\
h\left(e_{1}, e_{3}\right)=\lambda_{3} J e_{3}, & h\left(e_{3}, e_{3}\right)=\lambda_{3} J e_{1}-a J e_{2}-b J e_{3},
\end{array}
$$

where $\lambda_{1}+\lambda_{2}+\lambda_{3}=0, \lambda_{1}>0, \lambda_{1}-2 \lambda_{2} \geq 0$ and $\lambda_{1}-2 \lambda_{3} \geq 0$. If $\lambda_{2}=\lambda_{3}$, we can choose $e_{2}$ and $e_{3}$ such that $b=0$. Then by a straightforward computation, we have the following lemma:

Lemma 3.1. Let $\left\{e_{1}, e_{2}, e_{3}\right\}$ be the basis constructed above. Then

$$
\left(S\left(e_{i}, e_{j}\right)\right)=\left(\begin{array}{ccc}
2-\lambda_{1}^{2}-\lambda_{2}^{2}-\lambda_{3}^{2} & -\left(\lambda_{2}-\lambda_{3}\right) a & -\left(\lambda_{2}-\lambda_{3}\right) b \\
-\left(\lambda_{2}-\lambda_{3}\right) a & 2-2 \lambda_{2}^{2}-2 a^{2}-2 b^{2} & 0 \\
-\left(\lambda_{2}-\lambda_{3}\right) b & 0 & 2-2 \lambda_{3}^{2}-2 a^{2}-2 b^{2}
\end{array}\right)
$$

Remark that if $\lambda_{2}=\lambda_{3}$, it immediately follows from Lemma 3.1 that $M$ is quasiEinstein.

Lemma 3.2. Let $M^{3}$ be a 3-dimensional totally real submanifold of $S^{6}$ with the second fundamental form $h$. Then the Ricci tensor $S$ has a double eigenvalue at a point pof $M^{3}$
if and only if $p$ is a totally geodesic point or there exists an orthonormal basis $\left\{e_{1}, e_{2}, e_{3}\right\}$ of $T_{p} M$ such that either

$$
\begin{array}{ll}
h\left(e_{1}, e_{1}\right)=\lambda J e_{1}, & h\left(e_{2}, e_{2}\right)=-\lambda J e_{1}, \\
h\left(e_{1}, e_{2}\right)=-\lambda J e_{2}, & h\left(e_{2}, e_{3}\right)=0, \\
h\left(e_{1}, e_{3}\right)=0, & h\left(e_{3}, e_{3}\right)=0 \\
h\left(e_{1}, e_{1}\right)=2 \lambda J e_{1}, & h\left(e_{2}, e_{2}\right)=-\lambda J e_{1}+a J e_{2},  \tag{2}\\
h\left(e_{1}, e_{2}\right)=-\lambda J e_{2}, & h\left(e_{2}, e_{3}\right)=-a J e_{3}, \\
h\left(e_{1}, e_{3}\right)=-\lambda J e_{3}, & h\left(e_{3}, e_{3}\right)=-\lambda J e_{1}-a J e_{2},
\end{array}
$$

where $\lambda$ is a non-zero number.
Proof. If $p$ is a totally geodesic point of $M^{3}$, there is nothing to prove. Hence, we may assume that $p$ is not totally geodesic and we can use the basis $\left\{e_{1}, e_{2}, e_{3}\right\}$ constructed above. So we see that

$$
\begin{aligned}
& \operatorname{Ric}\left(e_{2}\right)=-\left(\lambda_{2}-\lambda_{3}\right) a e_{1}+2\left(1-\lambda_{2}^{2}-a^{2}-b^{2}\right) e_{2} \\
& \left\langle\operatorname{Ric}\left(e_{1}\right), e_{3}\right\rangle=-\left(\lambda_{2}-\lambda_{3}\right) b
\end{aligned}
$$

Since $M^{3}$ is quasi-Einstein, we know that $e_{2}, \operatorname{Ric}\left(e_{2}\right)$ and $\operatorname{Ric}\left(\operatorname{Ric}\left(e_{2}\right)\right)$ have to be linearly dependent. Hence the above formulas imply that

$$
a b\left(\lambda_{2}-\lambda_{3}\right)^{2}=0
$$

If $\lambda_{2}=\lambda_{3}$, we see that $\left\{e_{1}, e_{2}, e_{3}\right\}$ satisfies Case (2) of Lemma 3.2 by rechoosing $e_{2}$ and $e_{3}$ if necessary. Therefore, we may assume that $\lambda_{2} \neq \lambda_{3}$. Then, if necessary by interchanging $e_{2}$ and $e_{3}$, we may assume that $b=0$.

Suppose now that $a=0$. Hence $e_{1}, e_{2}$ and $e_{3}$ are eigenvectors of Ric. Since we assumed that $\lambda_{2} \neq \lambda_{3}$, we see that (if necessary after interchanging $e_{2}$ and $e_{3}$, which is allowed in this case since $a$ and $b$ both vanish) $M^{3}$ is quasi-Einstein if and only if

$$
2-\lambda_{1}^{2}-\lambda_{2}^{2}-\lambda_{3}^{2}=2-2 \lambda_{2}^{2}
$$

which reduces to

$$
-2 \lambda_{1}^{2}-2 \lambda_{1} \lambda_{2}=0
$$

Hence, since $\lambda_{1} \neq 0$, we see that $\lambda_{2}=-\lambda_{1}$ and $\lambda_{3}=0$. Thus $\left\{e_{1}, e_{2}, e_{3}\right\}$ is a basis as described in Case (1) of Lemma 3.2.

Finally, we consider the case that $\lambda_{2} \neq \lambda_{3}$ and $a \neq 0$. Since $a \neq 0$, we see that $M^{3}$ is quasi-Einstein if and only if $2-2 \lambda_{3}^{2}-2 a^{2}$ is a double eigenvalue of $S$. This is the case if and only if

$$
\operatorname{det}\left(\begin{array}{cc}
\lambda_{3}^{2}-\lambda_{1}^{2}-\lambda_{2}^{2}+2 a^{2} & \left(\lambda_{3}-\lambda_{2}\right) a \\
\left(\lambda_{3}-\lambda_{2}\right) a & 2\left(\lambda_{3}^{2}-\lambda_{2}^{2}\right)
\end{array}\right)=0
$$

Since $\lambda_{2} \neq \lambda_{3}$ and $\lambda_{3}=-\lambda_{1}-\lambda_{2}$, this is the case if and only if

$$
\operatorname{det}\left(\begin{array}{cc}
2 \lambda_{1} \lambda_{2}+2 a^{2} & -\left(\lambda_{1}+2 \lambda_{2}\right) a \\
a & -2 \lambda_{1}
\end{array}\right)=0
$$

i.e., if and only if

$$
\begin{equation*}
4 \lambda_{1}^{2} \lambda_{2}=-3 \lambda_{1} a^{2}+2 \lambda_{2} a^{2} . \tag{3.1}
\end{equation*}
$$

Now, we consider the following change of basis

$$
\begin{aligned}
& u_{1}=\frac{1}{\sqrt{a^{2}+4 \lambda_{1}^{2}}}\left(a e_{1}-2 \lambda_{1} e_{2}\right), \\
& u_{2}=\frac{1}{\sqrt{a^{2}+4 \lambda_{1}^{2}}}\left(2 \lambda_{1} e_{1}+a e_{2}\right), \\
& u_{3}=e_{3}
\end{aligned}
$$

Then, using (3.1), we have

$$
\begin{aligned}
h\left(a e_{1}-2 \lambda_{1} e_{2}, a e_{1}-2 \lambda_{1} e_{2}\right) & =\left(a^{2} \lambda_{1}+4 \lambda_{1}^{2} \lambda_{2}\right) J e_{1}+\left(-4 a \lambda_{1} \lambda_{2}+4 a \lambda_{1}^{2}\right) J e_{2} \\
& =-2\left(\lambda_{1}-\lambda_{2}\right) a\left(a J e_{1}-2 \lambda_{1} J e_{2}\right), \\
h\left(a e_{1}-2 \lambda_{1} e_{2}, e_{3}\right) & =a\left(\lambda_{1}-\lambda_{2}\right) J e_{3}, \\
h\left(2 \lambda_{1} e_{1}+a e_{2}, e_{3}\right) & =\left(2 \lambda_{1} \lambda_{3}-a^{2}\right) J e_{3}, \\
h\left(a e_{1}-2 \lambda_{1} e_{2}, 2 \lambda_{1} e_{1}+a e_{2}\right) & =\left(2 a \lambda_{1}^{2}-2 a \lambda_{1} \lambda_{2}\right) J e_{1}+\left(a^{2} \lambda_{2}-2 a^{2} \lambda_{1}-4 \lambda_{1}^{2} \lambda_{2}\right) J e_{2} \\
& =a\left(\lambda_{1}-\lambda_{2}\right)\left(2 \lambda_{1} J e_{1}+a J e_{2}\right) .
\end{aligned}
$$

Using now the minimality of $M$, together with the fact that $\langle h(X, Y), J Z\rangle$ is totally symmetric it follows that the basis $\left\{u_{1}, u_{2}, u_{3}\right\}$ satisfies Case (2) of Lemma 3.2.

Remark 3.3. An elementary computation shows that if Case (1) of Lemma 3.2 is satisfied, the Ricci tensor has eigenvalues $2\left(1-\lambda^{2}\right), 2\left(1-\lambda^{2}\right)$ and 2 , while if Case (2) is satisfied its eigenvalues are $2-6 \lambda^{2}, 2-2 \lambda^{2}-2 a^{2}$ and $2-2 \lambda^{2}-2 a^{2}$.

Remark 3.4. Submanifolds satisfying Case (1) of Lemma 3.2 are exactly those totally real submanifolds of $S^{6}$ which satisfy Chen's equality (see [CDVV1], [CDVV2] and [DV]). A complete classification of these submanifolds was obtained in [DV].

## 4. Examples of totally real submanifolds.

Example 4.1. We recall from [DVV] the following example: Consider the unit sphere

$$
S^{3}=\left\{\left(y_{1}, y_{2}, y_{3}, y_{4}\right) \in \boldsymbol{R}^{4} \mid y_{1}^{2}+y_{2}^{2}+y_{3}^{2}+y_{4}^{2}=1\right\}
$$

in $\boldsymbol{R}^{4}$. Let $X_{1}, X_{2}$ and $X_{3}$ be the vector fields defined by

$$
\begin{aligned}
& X_{1}\left(y_{1}, y_{2}, y_{3}, y_{4}\right)=\left(y_{2},-y_{1}, y_{4},-y_{3}\right), \\
& X_{2}\left(y_{1}, y_{2}, y_{3}, y_{4}\right)=\left(y_{3},-y_{4},-y_{1}, y_{2}\right), \\
& X_{3}\left(y_{1}, y_{2}, y_{3}, y_{4}\right)=\left(y_{4}, y_{3},-y_{2},-y_{1}\right) .
\end{aligned}
$$

Then $X_{1}, X_{2}$ and $X_{3}$ form a basis of tangent vector fields to $S^{3}$. Moreover, we have $\left[X_{1}, X_{2}\right]=$ $2 X_{3},\left[X_{2}, X_{3}\right]=2 X_{1}$ and $\left[X_{3}, X_{1}\right]=2 X_{2}$. Inspired by $[\mathrm{M}]$, we define a metric $\langle., .\rangle_{1}$ on $S^{3}$ such that $X_{1}, X_{2}$ and $X_{3}$ are orthogonal and such that $\left\langle X_{2}, X_{2}\right\rangle_{1}=\left\langle X_{3}, X_{3}\right\rangle_{1}=8 / 3$ and
$\left\langle X_{1}, X_{1}\right\rangle_{1}=4 / 9$. Then $E_{1}=(3 / 2) X_{1}, E_{2}=(\sqrt{3} / 2 \sqrt{2}) X_{2}$ and $E_{3}=-(\sqrt{3} / 2 \sqrt{2}) X_{3}$ form an orthonormal basis on $S^{3}$. We denote the Levi-Civita connection of $\langle\text {., . }\rangle_{1}$ by $\nabla$. We recall from [DVV] that there exists an isometric totally real immersion $\psi$ from $\left(S^{3},\langle., .\rangle_{1}\right)$ given by

$$
\psi: S^{3}(1) \rightarrow S^{6}(1):\left(y_{1}, y_{2}, y_{3}, y_{4}\right) \mapsto\left(x_{1}, x_{2}, x_{3}, x_{4}, x_{5}, x_{6}, x_{7}\right)
$$

where

$$
\begin{aligned}
& x_{1}=(1 / 9)\left(5 y_{1}^{2}+5 y_{2}^{2}-5 y_{3}^{2}-5 y_{4}^{2}+4 y_{1}\right) \\
& x_{2}=-(2 / 3) y_{2} \\
& x_{3}=(2 \sqrt{5} / 9)\left(y_{1}^{2}+y_{2}^{2}-y_{3}^{2}-y_{4}^{2}-y_{1}\right) \\
& x_{4}=(\sqrt{3} / 9 \sqrt{2})\left(-10 y_{3} y_{1}-2 y_{3}-10 y_{2} y_{4}\right) \\
& x_{5}=(\sqrt{3} \sqrt{5} / 9 \sqrt{2})\left(2 y_{1} y_{4}-2 y_{4}-2 y_{2} y_{3}\right) \\
& x_{6}=(\sqrt{3} \sqrt{5} / 9 \sqrt{2})\left(2 y_{1} y_{3}-2 y_{3}+2 y_{2} y_{4}\right) \\
& x_{7}=(\sqrt{3} / 9 \sqrt{2})\left(10 y_{1} y_{4}+2 y_{4}-10 y_{2} y_{3}\right)
\end{aligned}
$$

Its connection is given by

$$
\begin{array}{lll}
\nabla_{E_{1}} E_{1}=0, & \nabla_{E_{2}} E_{2}=0, & \nabla_{E_{3}} E_{3}=0, \\
\nabla_{E_{1}} E_{2}=-(11 / 4) E_{3}, & \nabla_{E_{1}} E_{3}=(11 / 4) E_{2}, & \nabla_{E_{2}} E_{3}=-(1 / 4) E_{1}, \\
\nabla_{E_{2}} E_{1}=(1 / 4) E_{3}, & \nabla_{E_{3}} E_{1}=-(1 / 4) E_{2}, & \nabla_{E_{3}} E_{2}=(1 / 4) E_{1},
\end{array}
$$

and its second fundamental form satisfies

$$
\begin{array}{ll}
h\left(E_{1}, E_{1}\right)=(\sqrt{5} / 2) J E_{1}, & h\left(E_{3}, E_{1}\right)=-(\sqrt{5} / 4) J E_{3} \\
h\left(E_{1}, E_{2}\right)=-(\sqrt{5} / 4) J E_{2}, & h\left(E_{3}, E_{2}\right)=0, \\
h\left(E_{2}, E_{2}\right)=-(\sqrt{5} / 4) J E_{1}, & h\left(E_{3}, E_{3}\right)=-(\sqrt{5} / 4) J E_{1}
\end{array}
$$

Hence $\tilde{\psi}$ is quasi-Einstein.
Example 4.2. Here, we will consider tubes in the direction of the orthogonal complement of the first osculating space on an almost complex curve. In [E2], N. Ejiri already showed that a tube with radius $\cos ^{2} \gamma=5 / 9$ on a superminimal almost complex curve defines a totally real submanifold of $S^{6}$, and in [DV] it was shown that a tube with radius $\pi / 2$ on any almost complex curve defines a totally real submanifold $M$ with $\delta_{M}=2$.

An immersion $\bar{\phi}: N \rightarrow S^{6}(1)$ is called almost complex if $J$ preserves the tangent space, i.e., $J_{p} \bar{\phi}_{\star}\left(T_{p} N\right)=\bar{\phi}_{\star}\left(T_{p} N\right)$. It is well-known that such immersions are always minimal, and as indicated in [BVW] there are essentially 4 types of almost complex immersions in $S^{6}(1)$, namely, those which are
(I) linearly full in $S^{6}(1)$ and superminimal,
(II) linearly full in $S^{6}$ (1) but not superminimal,
(III) linearly full in some totally geodesic $S^{5}(1)$ in $S^{6}(1)$ (and thus by [Ca1] necessarily not superminimal),
(IV) totally geodesic.

Now, let $\bar{\phi}: N \rightarrow S^{6}(1)$ be an almost complex curve. We denote its position vector in $\boldsymbol{R}^{7}$ also by $\bar{\phi}$. For the proof of elementary properties of such surfaces, we refer to [S]. Here, we simply recall that for tangent vector fields $X$ and $Y$ to $N$ and for a normal vector field $\eta$, we have

$$
\begin{align*}
& \alpha(X, J Y)=J \alpha(X, Y)  \tag{4.1}\\
& A_{J \eta}=J A_{\eta}=-A_{\eta} J  \tag{4.2}\\
& \nabla_{\frac{1}{X}}^{\perp} J \eta=G(X, \eta)+J \nabla_{X}^{\perp} \eta  \tag{4.3}\\
& (\nabla \alpha)(X, Y, J Z)=J(\nabla \alpha)(X, Y, Z)+G\left(\bar{\phi}_{\star} X, \alpha(Y, Z)\right), \tag{4.4}
\end{align*}
$$

where $\alpha$ denotes the second fundamental form of the immersion and the pull-back of $J$ to $N$ is also denoted by $J$.

Next, if necessary, by restricting ourselves to an open dense subset of $N$, we may assume that $N$ does not contain any totally geodesic points. Let $p \in N$ and $V$ be an arbitrary unit tangent vector field defined on a neighborhood $W$ of $p$. We define a local non vanishing function $\mu=\|\alpha(V, V)\|$ and an orthogonal tangent vector field $U$ such that $\bar{\phi}_{\star} U=J \bar{\phi}_{\star} V=$ $\bar{\phi} \times \bar{\phi}_{\star} V$. Then, using the properties of the vector cross product, it is easy to see that $F_{1}=\bar{\phi}$, $F_{2}=\bar{\phi}_{\star} V, F_{3}=J \bar{\phi}_{\star} V, F_{4}=\alpha(V, V) / \mu, F_{5}=\alpha(V, J V) / \mu=J \alpha(V, V) / \mu=F_{1} \times F_{4}$, $F_{6}=F_{2} \times \alpha(V, V) / \mu$ and $F_{7}=F_{3} \times \alpha(V, V) / \mu$ form a $G_{2}$-frame and hence satisfy the multiplication table as defined in Section 2.

Since $F_{4}, \ldots, F_{7}$ form a basis for the normal space along $N$, it is clear that we can write any normal vector field as a linear combination of these basis vector fields. Thus there exist functions $a_{1}, \ldots, a_{4}$ such that

$$
\begin{equation*}
(\nabla \alpha)(V, V, V)=\mu\left(a_{1} F_{4}+a_{2} F_{5}+a_{3} F_{6}+a_{4} F_{7}\right) . \tag{4.5}
\end{equation*}
$$

Then using (4.4) and the multiplication table, we get that

$$
\begin{equation*}
(\nabla \alpha)(V, V, U)=\mu\left(-a_{2} F_{4}+a_{1} F_{5}+\left(1+a_{4}\right) F_{6}-a_{3} F_{7}\right) \tag{4.6}
\end{equation*}
$$

From (4.5) and (4.6), it is immediately clear that
(1) $N$ is an almost complex curve of Type (I) if and only if $a_{3}=0$ and $a_{4}=-1 / 2$.
(2) $N$ is an almost complex curve of Type (III) if and only if $a_{4}+a_{3}^{2}+a_{4}^{2}=0$.

Introducing local functions $\mu_{1}$ and $\mu_{2}$ on $N$ by

$$
\nabla_{V} V=\mu_{1} U, \quad \nabla_{U} U=\mu_{2} V, \quad \nabla_{V} U=-\mu_{1} V, \quad \nabla_{U} V=-\mu_{2} U,
$$

it follows from (4.5) and (4.6) that $a_{1}=V(\mu) / \mu$ and $a_{2}=-U(\mu) / \mu$.
Now, in order to construct explicitly the totally real immersion from the unit tangent bundle, we recall a technical lemma from [DV].

Lemma 4.1. Denote by $D$ the standard connection on $\boldsymbol{R}^{7}$. Then, we have

$$
\begin{aligned}
& D_{V}\left(\mu F_{4}\right)=\mu\left(-\mu F_{2}+a_{1} F_{4}+\left(a_{2}+2 \mu_{1}\right) F_{5}+a_{3} F_{6}+a_{4} F_{7}\right), \\
& D_{U}\left(\mu F_{4}\right)=\mu\left(\mu F_{3}-a_{2} F_{4}+\left(a_{1}-2 \mu_{2}\right) F_{5}+\left(1+a_{4}\right) F_{6}-a_{3} F_{7}\right), \\
& D_{V}\left(\mu F_{5}\right)=\mu\left(-\mu F_{3}-\left(a_{2}+2 \mu_{1}\right) F_{4}+a_{1} F_{5}+\left(1+a_{4}\right) F_{6}-a_{3} F_{7}\right), \\
& D_{U}\left(\mu F_{5}\right)=\mu\left(-\mu F_{2}-\left(a_{1}-2 \mu_{2}\right) F_{4}-a_{2} F_{5}-a_{3} F_{6}-a_{4} F_{7}\right), \\
& D_{V}\left(\mu F_{6}\right)=\mu\left(-a_{3} F_{4}-\left(a_{4}+1\right) F_{5}+a_{1} F_{6}+\left(a_{2}+3 \mu_{1}\right) F_{7}\right), \\
& D_{U}\left(\mu F_{6}\right)=\mu\left(-\left(a_{4}+1\right) F_{4}+a_{3} F_{5}-a_{2} F_{6}+\left(a_{1}-3 \mu_{2}\right) F_{7}\right), \\
& D_{V}\left(\mu F_{7}\right)=\mu\left(-a_{4} F_{4}+a_{3} F_{5}-\left(a_{2}+3 \mu_{1}\right) F_{6}+a_{1} F_{7}\right), \\
& D_{U}\left(\mu F_{7}\right)=\mu\left(a_{3} F_{4}+a_{4} F_{5}+\left(3 \mu_{2}-a_{1}\right) F_{6}-a_{2} F_{7}\right) .
\end{aligned}
$$

Proof of Theorem 1. We define a map

$$
\bar{\psi}: U N \rightarrow S^{6}(1): v_{p} \mapsto \cos \gamma \bar{\phi}(p)+\sin \gamma \bar{\phi}_{\star}(v) \times \frac{\alpha(v, v)}{\|\alpha(v, v)\|} .
$$

Using the above vector fields, we can write $v_{p}=\cos (t / 3) V+\sin (t / 3) U$; and an easy computation shows that the map $\bar{\psi}$ can be locally parameterized by

$$
\begin{equation*}
\bar{\psi}(q, t)=\cos \gamma F_{1}(q)+\sin \gamma\left(\cos t F_{6}(q)+\sin t F_{7}(q)\right), \tag{4.7}
\end{equation*}
$$

where $q \in W$ and $t \in \boldsymbol{R}$. Since the case with $\gamma=\pi / 2$ was already treated in [DV], we restrict ourselves here to the case that $\cos \gamma \neq 0$. We immediately see that

$$
\begin{equation*}
\bar{\psi}_{\star}\left(\frac{\partial}{\partial t}\right)=\sin \gamma\left(-\sin t F_{6}+\cos t F_{7}\right) . \tag{4.8}
\end{equation*}
$$

Using Lemma 4.1, we then obtain that

$$
\begin{align*}
\bar{\psi}_{\star}= & \cos \gamma D_{V} F_{1}+\sin \gamma\left(\cos t D_{V} F_{6}+\sin t D_{V} F_{7}\right) \\
= & \cos \gamma F_{2}+\sin \gamma\left(-\cos t(V(\mu) / \mu) F_{6}\right) \\
& +\cos t\left(-a_{3} F_{4}-\left(a_{4}+1\right) F_{5}+a_{1} F_{6}+\left(a_{2}+3 \mu_{1}\right) F_{7}\right)-\sin t(V(\mu) / \mu) F_{7}  \tag{4.9}\\
& +\sin t\left(-a_{4} F_{4}+a_{3} F_{5}-\left(a_{2}+3 \mu_{1}\right) F_{6}+a_{1} F_{7}\right) \\
= & \cos \gamma F_{2}+\sin \gamma\left(\left(-a_{3} \cos t-a_{4} \sin t\right) F_{4}\right. \\
& \left.+\left(a_{3} \sin t-\left(a_{4}+1\right) \cos t\right) F_{5}\right)+\left(3 \mu_{1}-(U(\mu) / \mu)\right) \bar{\psi}_{\star}\left(\frac{\partial}{\partial t}\right) .
\end{align*}
$$

Using similar computations, we also get that

$$
\begin{align*}
\bar{\psi}_{\star}(U)= & \cos \gamma F_{3}+\sin \gamma\left(\left(a_{3} \sin t-\left(1+a_{4}\right) \cos t\right) F_{4}\right. \\
& \left.+\left(a_{3} \cos t+a_{4} \sin t\right) F_{5}\right)+\left(-3 \mu_{2}+(V(\mu) / \mu)\right) \bar{\psi}_{\star}\left(\frac{\partial}{\partial t}\right) . \tag{4.10}
\end{align*}
$$

From (4.8), (4.9) and (4.10), we see that $\bar{\psi}$ is an immersion at every point $(q, t)$.

Now, we put

$$
\begin{aligned}
& X=V-\left(3 \mu_{1}-(U(\mu) / \mu)\right) \frac{\partial}{\partial t} \\
& Y=U-\left(-3 \mu_{2}+(V(\mu) / \mu)\right) \frac{\partial}{\partial t}
\end{aligned}
$$

A straightforward computation, using the multiplication table of Section 2, then shows that

$$
\begin{aligned}
\bar{\psi} \times \bar{\psi}_{\star}\left(\frac{\partial}{\partial t}\right)= & -\sin ^{2} \gamma F_{1}+\cos \gamma \sin \gamma\left(\cos t F_{6}+\sin t F_{7}\right), \\
\bar{\psi} \times \bar{\psi}_{\star}(X)= & \sin ^{2} \gamma\left(a_{3} \cos 2 t+a_{4} \sin 2 t+(1 / 2) \sin 2 t\right) F_{2} \\
& +\left(\cos ^{2} \gamma+\sin ^{2} \gamma\left(a_{3} \sin 2 t-a_{4} \cos 2 t-\cos ^{2} t\right)\right) F_{3} \\
& +\cos \gamma \sin \gamma\left(-a_{3} \sin t+\left(a_{4}+2\right) \cos t\right) F_{4} \\
& +\cos \gamma \sin \gamma\left(-a_{3} \cos t-\left(a_{4}-1\right) \sin t\right) F_{5} .
\end{aligned}
$$

Consequently, $\bar{\psi}$ is a totally real immersion if and only if

$$
\left\langle\bar{\psi} \times \bar{\psi}_{\star}(X), \bar{\psi}_{\star}(Y)\right\rangle=0,
$$

i.e., if and only if

$$
\begin{aligned}
\cos \gamma\left(\cos ^{2} \gamma+\sin ^{2} \gamma\right. & \left(3 a_{3} \sin 2 t-a_{4} \cos 2 t-\cos ^{2} t-a_{3}^{2}\right. \\
& \left.\left.-\left(a_{4}+2\right)\left(1+a_{4}\right) \cos ^{2} t-a_{4}\left(a_{4}-1\right) \sin ^{2} t\right)\right)=0 .
\end{aligned}
$$

Hence, since we assumed $\cos \gamma \neq 0$, we find that

$$
3 a_{3} \sin ^{2} \gamma \sin 2 t-3\left(a_{4}+1 / 2\right) \sin ^{2} \gamma \cos 2 t+\cos ^{2} \gamma-\sin ^{2} \gamma\left(a_{4}+a_{3}^{2}+a_{4}^{2}+3 / 2\right)=0
$$

Since the above formula has to be satisfied for every value of $t$, we deduce that $a_{3}=0$, $a_{4}=-1 / 2$ and $\cos ^{2} \gamma=5 / 9$. Hence $N^{2}$ is a superminimal almost complex curve in $S^{6}$ and the radius of the tube satisfies $\cos ^{2} \gamma=5 / 9$. Using the above values for $a_{3}$ and $a_{4}$, we then obtain that

$$
\begin{aligned}
& \bar{\psi}_{\star}(X)=\cos \gamma F_{2}+(1 / 2) \sin \gamma\left(\sin t F_{4}-\cos t F_{5}\right), \\
& \bar{\psi}_{\star}(Y)=\cos \gamma F_{3}-(1 / 2) \sin \gamma\left(\cos t F_{4}+\sin t F_{5}\right), \\
& J \bar{\psi}_{\star}\left(\frac{\partial}{\partial t}\right)=-\sin ^{2} \gamma F_{1}+\cos \gamma \sin \gamma\left(\cos t F_{6}+\sin t F_{7}\right), \\
& J \bar{\psi}_{\star}(X)=(1 / 3) F_{3}+(3 / 2) \cos \gamma \sin \gamma\left(\cos t F_{4}+\sin t F_{5}\right), \\
& J \bar{\psi}_{\star}(Y)=-(1 / 3) F_{2}+(3 / 2) \cos \gamma \sin \gamma\left(\sin t F_{4}-\cos t F_{5}\right) .
\end{aligned}
$$

Therefore, by a straightforward computation, we obtain that

$$
\begin{aligned}
& D_{\frac{\partial}{\partial t}} \bar{\psi}_{\star}(X)=(1 / 2) \sin \gamma\left(\cos t F_{4}+\sin t F_{5}\right)=-(1 / 2)\left((1 / 3) \bar{\psi}_{\star}(Y)-\cos \gamma J \bar{\psi}_{\star}(X)\right), \\
& D_{\frac{\partial}{\partial t}} \bar{\psi}_{\star}(Y)=-(1 / 2) \sin \gamma\left(-\sin t F_{4}+\cos t F_{5}\right)=(1 / 2)\left((1 / 3) \bar{\psi}_{\star}(X)+\cos \gamma J \bar{\psi}_{\star}(Y)\right), \\
& D_{\frac{\partial}{\partial t}} \bar{\psi}_{\star}\left(\frac{\partial}{\partial t}\right)=-\sin \gamma\left(\cos t F_{6}+\sin t F_{7}\right)=-\left((4 / 9) \bar{\psi}+\cos \gamma J \bar{\psi}_{\star}\left(\frac{\partial}{\partial t}\right)\right) .
\end{aligned}
$$

So, if we put $E_{1}=(3 / 2)(\partial / \partial t), E_{2}=(\sqrt{3} / \sqrt{2}) X$ and $E_{3}=(\sqrt{3} / \sqrt{2}) Y$ we see that $E_{1}, E_{2}$ and $E_{3}$ form an orthonormal basis of the tangent space to $U N$ and

$$
\begin{aligned}
& h\left(E_{1}, E_{1}\right)=-(3 / 2) \cos \gamma J \bar{\psi}_{\star}\left(E_{1}\right), \\
& h\left(E_{1}, E_{2}\right)=(3 / 4) \cos \gamma J \bar{\psi}_{\star}\left(E_{2}\right), \\
& h\left(E_{1}, E_{3}\right)=(3 / 4) \cos \gamma J \bar{\psi}_{\star}\left(E_{3}\right) .
\end{aligned}
$$

Since $U N$ is totally real (and thus minimal) and $\langle h(X, Y), J Z\rangle$ is totally symmetric in $X$, $Y$ and $Z$, the above formulas and Lemma 3.1 imply that $\bar{\psi}$ is quasi-Einstein. Since the first normal space is 3-dimensional, with respect to the induced metric we have $\delta_{U N}<2$ (see [C]). Hence $U N$ satisfies Case (2) of Lemma 3.2.
5. Proof of Theorem 2. Throughout this section we will assume that $F: M^{3} \rightarrow S^{6}$ is a totally real immersion which is quasi-Einstein. Unless otherwise indicated, we will identify $M^{3}$ with its image in $S^{6}$.

First, we remark that if $\delta_{M}=2$, there is nothing to prove. Next, we assume that $M^{3}$ is Einstein. Since a 3-dimensional Einstein manifold has constant sectional curvatures, it follows from [E1] that a neighborhood of $p$ is $G_{2}$-congruent with an open part of the image of the totally real immersion of $S^{3}(1 / 16)$ in $S^{6}(1)$ as described in [E1] (see also [DVV]). From [E2], we also know that we can consider this image as the tube with radius $\gamma$, with $\cos ^{2} \gamma=5 / 9$ on the almost complex curve with constant Gaussian curvature $1 / 6$. This completes the proof in this case.

Next assume that $p \in M$ such that Ric has an eigenvalue with multiplicity 2 and $\delta_{M}(p) \neq$ 2. Since in a neighborhood of $p, M$ is quasi-Einstein, but not Einstein, there exist local orthonormal vector fields $\left\{E_{1}, E_{2}, E_{3}\right\}$ such that $E_{1}$ spans the 1-dimensional eigenspace and $\left\{E_{2}, E_{3}\right\}$ span the 2 -dimensional eigenspace. Hence, applying Lemma 3.2, we see that there exist local functions $\lambda, a$ and $b$ such that

$$
\begin{array}{ll}
h\left(E_{1}, E_{1}\right)=2 \lambda J E_{1}, & h\left(E_{2}, E_{2}\right)=-\lambda J E_{1}+a J E_{2}, \\
h\left(E_{1}, E_{2}\right)=-\lambda J E_{2}, & h\left(E_{2}, E_{3}\right)=-a J E_{3}, \\
h\left(E_{1}, E_{3}\right)=-\lambda J E_{3}, & h\left(E_{3}, E_{3}\right)=-\lambda J E_{1}-a J E_{2} .
\end{array}
$$

If necessary by changing the sign of $E_{3}$, we may assume that $G\left(E_{1}, E_{2}\right)=J E_{3}, G\left(E_{2}, E_{3}\right)=$ $J E_{1}$ and $G\left(E_{3}, E_{1}\right)=J E_{2}$. We now introduce local functions $a_{1}, a_{2}, a_{3}, b_{1}, b_{2}, b_{3}, c_{1}, c_{2}$, $c_{3}$ by

$$
\begin{array}{lll}
\nabla_{E_{1}} E_{1}=a_{1} E_{2}+a_{2} E_{3}, & \nabla_{E_{1}} E_{2}=-a_{1} E_{1}+a_{3} E_{3}, & \nabla_{E_{1}} E_{3}=-a_{2} E_{1}-a_{3} E_{2}, \\
\nabla_{E_{2}} E_{1}=b_{1} E_{2}+b_{2} E_{3}, & \nabla_{E_{2}} E_{2}=-b_{1} E_{1}+b_{3} E_{3}, & \nabla_{E_{2}} E_{3}=-b_{2} E_{1}-b_{3} E_{2}, \\
\nabla_{E_{3}} E_{1}=c_{1} E_{2}+c_{2} E_{3}, & \nabla_{E_{3}} E_{2}=-c_{1} E_{1}+c_{3} E_{3}, & \nabla_{E_{3}} E_{3}=-c_{2} E_{1}-c_{3} E_{2} .
\end{array}
$$

LEMMA 5.1. The function $\lambda$ satisfies $\lambda=\sqrt{5} / 4$. Moreover, $a_{1}=a_{2}=c_{2}=b_{1}=0$ and $b_{2}=-c_{1}=1 / 4$.

Proof. A straightforward computation shows that

$$
\begin{aligned}
(\nabla h)\left(E_{2}, E_{1}, E_{1}\right)= & 2 E_{2}(\lambda) J E_{1}-2 \lambda J E_{3}+4 \lambda\left(b_{1} J E_{2}+b_{2} J E_{3}\right), \\
(\nabla h)\left(E_{1}, E_{2}, E_{1}\right)= & -E_{1}(\lambda) J E_{2}-\lambda J E_{3}-\lambda\left(-a_{1} J E_{1}\right) \\
& +2 \lambda a_{1} J E_{1}-a_{1}\left(-\lambda J E_{1}+a J E_{2}\right)-a_{2}\left(-a J E_{3}\right) .
\end{aligned}
$$

Hence, it follows from the Codazzi equation $(\nabla h)\left(E_{2}, E_{1}, E_{1}\right)=(\nabla h)\left(E_{1}, E_{2}, E_{1}\right)$ that

$$
\begin{align*}
& E_{2}(\lambda)=2 \lambda a_{1},  \tag{5.1}\\
& E_{1}(\lambda)=-4 \lambda b_{1}-a a_{1},  \tag{5.2}\\
& 4 b_{2}=1+(a / \lambda) a_{2} . \tag{5.3}
\end{align*}
$$

Similarly, we obtain from the Codazzi equation $(\nabla h)\left(E_{3}, E_{1}, E_{1}\right)=(\nabla h)\left(E_{1}, E_{3}, E_{1}\right)$ that

$$
\begin{align*}
& E_{3}(\lambda)=2 \lambda a_{2},  \tag{5.4}\\
& E_{1}(\lambda)=-4 \lambda c_{2}+a a_{1},  \tag{5.5}\\
& 4 c_{1}=-1+(a / \lambda) a_{2} . \tag{5.6}
\end{align*}
$$

Comparing (5.5) and (5.2), we get that

$$
\begin{equation*}
c_{2}-b_{1}=(a / 2 \lambda) a_{1} \tag{5.7}
\end{equation*}
$$

A straightforward computation, using (5.1) and (5.4), then shows that

$$
\begin{align*}
& (\nabla h)\left(E_{1}, E_{2}, E_{3}\right)=a a_{2} J E_{1}+\left(3 a a_{3}+a-a_{2} \lambda\right) J E_{2}-\left(a_{1} \lambda+E_{1}(a)\right) J E_{3},  \tag{5.8}\\
& (\nabla h)\left(E_{2}, E_{1}, E_{3}\right)=\left(4 \lambda b_{2}-\lambda\right) J E_{1}+b_{2} a J E_{2}+\left(b_{1} a-2 \lambda a_{1}\right) J E_{3},  \tag{5.9}\\
& (\nabla h)\left(E_{3}, E_{1}, E_{2}\right)=\left(4 \lambda c_{1}+\lambda\right) J E_{1}-\left(a c_{1}+2 \lambda a_{2}\right) J E_{2}+a c_{2} J E_{3} . \tag{5.10}
\end{align*}
$$

Therefore, using the Codazzi equations and (5.3), (5.6) and (5.7), we get that

$$
\left(\left(a^{2} / 2 \lambda\right)+2 \lambda\right) a_{2}=0, \quad\left(\left(a^{2} / 2 \lambda\right)+2 \lambda\right) a_{1}=0 .
$$

Hence $a_{1}=a_{2}=0$ and we deduce from the previous equations that

$$
c_{2}=b_{1}, \quad c_{1}=-1 / 4, \quad b_{2}=1 / 4 .
$$

This implies that the function $\lambda$ is a solution of the following system of differential equations:

$$
E_{1}(\lambda)=-4 \lambda b_{1}, \quad E_{2}(\lambda)=0, \quad E_{3}(\lambda)=0
$$

Since $\left[E_{2}, E_{3}\right]=-(1 / 2) E_{1}-b_{3} E_{2}-c_{3} E_{3}$, it immediately follows from the integrability conditions that $b_{1}=0$ and hence $\lambda$ is a constant.

To compute the actual value of $\lambda$ we use the Gauss equation. We have

$$
R\left(E_{1}, E_{2}\right) E_{1}=-E_{2}+\lambda^{2} E_{2}+2 \lambda^{2} E_{2}=\left(3 \lambda^{2}-1\right) E_{2} .
$$

On the other hand, we have

$$
\begin{aligned}
R\left(E_{1}, E_{2}\right) E_{1} & =\nabla_{E_{1}} \nabla_{E_{2}} E_{1}-\nabla_{E_{2}} \nabla_{E_{1}} E_{1}-\nabla_{\left[E_{1}, E_{2}\right]} E_{1} \\
& =\nabla_{E_{1}}\left((1 / 4) E_{3}\right)-\left(a_{3}-1 / 4\right) \nabla_{E_{3}} E_{1} \\
& =-(1 / 4) a_{3} E_{2}+(1 / 4)\left(a_{3}-1 / 4\right) E_{2}=-(1 / 16) E_{2} .
\end{aligned}
$$

Hence

$$
\lambda^{2}=5 / 16
$$

Since $\lambda$ is positive, this completes the proof of the lemma.
Now, in order to complete the proof of the theorem, we have to make a distinction between $M$ and its image under $F$ in $S^{6}$. First, we consider the case that $a$ is identically zero in a neighborhood of the point $p$. Then, we have the following lemma:

Lemma 5.2. There exists an orthonormal basis $\left\{E_{1}, E_{2}, E_{3}\right\}$ with $G\left(E_{1}, E_{2}\right)=J E_{3}$, $G\left(E_{2}, E_{3}\right)=J E_{1}$ and $G\left(E_{3}, E_{1}\right)=J E_{2}$, defined on a neighborhood of the point $p$ such that

$$
\begin{array}{ll}
h\left(E_{1}, E_{1}\right)=(\sqrt{5} / 2) J E_{1}, & h\left(E_{2}, E_{2}\right)=-(\sqrt{5} / 4) J E_{1}, \\
h\left(E_{1}, E_{2}\right)=-(\sqrt{5} / 4) J E_{2}, & h\left(E_{2}, E_{3}\right)=0, \\
h\left(E_{1}, E_{3}\right)=-(\sqrt{5} / 4) J E_{3}, & h\left(E_{3}, E_{3}\right)=-(\sqrt{5} / 4) J E_{1} .
\end{array}
$$

Moreover, they satisfy

$$
\begin{array}{lll}
\nabla_{E_{1}} E_{1}=0, & \nabla_{E_{1}} E_{2}=-(11 / 4) E_{3}, & \nabla_{E_{1}} E_{3}=(11 / 4) E_{2}, \\
\nabla_{E_{2}} E_{1}=(1 / 4) E_{3}, & \nabla_{E_{2}} E_{2}=0, & \nabla_{E_{2}} E_{3}=-(1 / 4) E_{1}, \\
\nabla_{E_{3}} E_{1}=-(1 / 4) E_{2}, & \nabla_{E_{3}} E_{2}=(1 / 4) E_{1}, & \nabla_{E_{3}} E_{3}=0 .
\end{array}
$$

Proof. We take the local orthonormal basis $\left\{E_{1}, E_{2} . E_{3}\right\}$ constructed in the previous lemma. Clearly, this basis already satisfies the first condition. Since $a_{1}=a_{2}=c_{2}=b_{1}=0$ and $b_{2}=-c_{1}=1 / 4$, the Gauss equations $\left\langle R\left(E_{1}, E_{2}\right) E_{2}, E_{3}\right\rangle=0,\left\langle R\left(E_{1}, E_{3}\right) E_{3}, E_{2}\right\rangle=0$ and $\left\langle R\left(E_{2}, E_{3}\right) E_{3}, E_{2}\right\rangle=21 / 16$ reduce to

$$
\begin{align*}
& E_{1}\left(b_{3}\right)-E_{2}\left(a_{3}\right)-c_{3}\left(a_{3}-1 / 4\right)=0  \tag{5.11}\\
& -E_{1}\left(c_{3}\right)+E_{3}\left(a_{3}\right)-b_{3}\left(a_{3}-1 / 4\right)=0  \tag{5.12}\\
& E_{3}\left(b_{3}\right)-E_{2}\left(c_{3}\right)-(1 / 2) a_{3}-b_{3}^{2}-c_{3}^{2}=11 / 8 \tag{5.13}
\end{align*}
$$

Now, we use the following transformation of the local frame $\left\{E_{1}, E_{2}, E_{3}\right\}$ :

$$
\begin{aligned}
& U_{1}=E_{1} \\
& U_{2}=\cos \theta E_{2}+\sin \theta E_{3} \\
& U_{3}=-\sin \theta E_{2}+\cos \theta E_{3} .
\end{aligned}
$$

where $\theta$ is an arbitrary locally defined function on $M$. It is immediately clear that $\left\{U_{1}, U_{2}, U_{3}\right\}$ satisfies the conditions of the lemma if and only if the function $\theta$ satisfies the following system of differential equations:

$$
\begin{aligned}
& d \theta\left(E_{1}\right)+a_{3}+11 / 4=0, \\
& d \theta\left(E_{2}\right)+b_{3}=0, \\
& d \theta\left(E_{3}\right)+c_{3}=0,
\end{aligned}
$$

i.e., $d \theta=-\left(a_{3}+11 / 4\right) \theta_{1}-b_{3} \theta_{2}-c_{3} \theta_{3}$, where $\left\{\theta_{1}, \theta_{2}, \theta_{3}\right\}$ is that dual basis of $\left\{E_{1}, E_{2}, E_{3}\right\}$. Now, this system locally has a solution if and only if the 1-form $\omega=\left(a_{3}+11 / 4\right) \theta_{1}+b_{3} \theta_{2}+$
$c_{3} \theta_{3}$ is closed. One can easily verify that $d \omega=0$ is equivalent with (5.11), (5.12) and (5.13).

The proof now follows from the Cartan-Ambrose-Hicks Theorem and the uniqueness theorem for totally real immersions.

Finally, we deal with the case that $a(p) \neq 0$. Then, we have the following lemma:
Lemma 5.3. Let $\left\{E_{1}, E_{2}, E_{3}\right\}$ be the local orthonormal frame constructed before. Then we have

$$
E_{1}(a)=0 \text { and } a_{3}=-1 / 4 .
$$

Proof. We look again at the proof of Lemma 5.2. From (5.8), (5.9), (5.10) and the Codazzi equation, we get $E_{1}(a)=0$ and $3 a\left(1 / 4+a_{3}\right)=0$. Since $a(p) \neq 0$, this completes the proof.

From now on we will make a distinction between $M$ and its image under $F$ in $S^{6}$. We will also write explicitly $J_{p} v$ as $p \times v$, since we will be using the almost complex structure at different points of $S^{6}$. Let us recall that we have a local basis $\left\{E_{1}, E_{2}, E_{3}\right\}$ on $U$ such that

$$
\begin{array}{lll}
\nabla_{E_{1}} E_{1}=0, & \nabla_{E_{1}} E_{2}=-(1 / 4) E_{3}, & \nabla_{E_{1}} E_{3}=(1 / 4) E_{2}, \\
\nabla_{E_{2}} E_{1}=(1 / 4) E_{3}, & \nabla_{E_{2}} E_{2}=b_{3} E_{3}, & \nabla_{E_{2}} E_{3}=-(1 / 4) E_{1}-b_{3} E_{2}, \\
\nabla_{E_{3}} E_{1}=-(1 / 4) E_{2}, & \nabla_{E_{3}} E_{2}=(1 / 4) E_{1}+c_{3} E_{3}, & \nabla_{E_{3}} E_{3}=-c_{3} E_{2} ; \\
h\left(E_{1}, E_{1}\right)=(\sqrt{5} / 2) F \times F_{\star} E_{1}, & h\left(E_{2}, E_{2}\right)=-(\sqrt{5} / 4) F \times F_{\star} E_{1}+a F \times F_{\star} E_{2}, \\
h\left(E_{1}, E_{2}\right)=-(\sqrt{5} / 4) F \times F_{\star} E_{2}, & h\left(E_{2}, E_{3}\right)=-a F \times F_{\star} E_{3}, \\
h\left(E_{1}, E_{3}\right)=-(\sqrt{5} / 4) F \times F_{\star} E_{3}, & h\left(E_{3}, E_{3}\right)=-(\sqrt{5} / 4) F \times F_{\star} E_{1}-a F \times F_{\star} E_{2} ; \\
& F_{\star} E_{1} \times F_{\star} E_{2}=F \times F_{\star} E_{3}, \\
F_{\star} E_{2} \times F_{\star} E_{3}=F \times F_{\star} E_{1}, \\
F_{\star} E_{3} \times F_{\star} E_{1}=F \times F_{\star} E_{2} .
\end{array}
$$

We now define a mapping $G: U \rightarrow S^{6}$, where $U$ is a neighborhood of $p$, by

$$
G(q)=(\sqrt{5} / 3) F(q)+(2 / 3) F \times F_{\star}\left(E_{1}(q)\right) .
$$

Then, using the above formulas, we find that

$$
\begin{aligned}
D_{E_{1}} G & =(\sqrt{5} / 3) F_{\star}\left(E_{1}\right)+(2 / 3) F \times h\left(E_{1}, E_{1}\right)=0, \\
D_{E_{2}} G & =(\sqrt{5} / 3) F_{\star}\left(E_{2}\right)+(2 / 3)\left(F_{\star}\left(E_{2}\right) \times F_{\star}\left(E_{1}\right)+F \times F_{\star}\left(\nabla_{E_{2}} E_{1}\right)+F \times h\left(E_{2}, E_{1}\right)\right) \\
& =(\sqrt{5} / 2) F_{\star}\left(E_{2}\right)-(1 / 2) F \times F_{\star}\left(E_{3}\right), \\
D_{E_{3}} G & =(\sqrt{5} / 3) F_{\star}\left(E_{3}\right)+(2 / 3)\left(F_{\star}\left(E_{3}\right) \times F_{\star}\left(E_{1}\right)+F \times F_{\star}\left(\nabla_{E_{3}} E_{1}\right)+F \times h\left(E_{3}, E_{1}\right)\right) \\
& =(\sqrt{5} / 2) F_{\star}\left(E_{3}\right)+(1 / 2) F \times F_{\star}\left(E_{2}\right),
\end{aligned}
$$

from which it follows that $G$ is not an immersion.
Using [ Sp, Vol. 1, p. 204], we can identify a neighborhood of $p$ with a neighborhood $I \times W_{1}$ of the origin in $\boldsymbol{R}^{3}$ (with coordinates $(t, u, v)$ ) such that $p=(0,0,0)$ and $E_{1}=\partial / \partial t$.

Then there exist functions $\alpha_{1}$ and $\alpha_{2}$ on $W_{1}$ such that $E_{2}+\alpha_{1} E_{1}$ and $E_{3}+\alpha_{2} E_{1}$ form a basis for the tangent space to $W_{1} \subset U$ at the point $q=(0, u, v)$.

Now since $\nabla_{E_{1}} E_{1}=0$ and $h\left(E_{1}, E_{1}\right)=(\sqrt{5} / 2) F \times F_{\star} E_{1}$, it follows that the integral curve of $E_{1}$ through the point $F(q)$ is a circle with radius $2 / 3$, tangent vector $F_{\star} E_{1}(q)$ and normal vector $(\sqrt{5} / 3) F \times F_{\star} E_{1}(q)-(2 / 3) F(q)$. From this it is clear that $F(U)$ can be reconstructed from $W_{1}$ by

$$
\begin{align*}
F(t, u, v)= & (\sqrt{5} / 3)\left((\sqrt{5} / 3) F(0, u, v)+(2 / 3) J F_{\star} E_{1}(0, u, v)\right) \\
& +(2 / 3)\left(\sin (3 t / 2) E_{1}(0, u, v)-\cos (3 t / 2)((2 / 3) F(0, u, v)\right.  \tag{5.14}\\
& \left.\left.-(\sqrt{5} / 3) F \times F_{\star} E_{1}(0, u, v)\right)\right)
\end{align*}
$$

Now, we look at the restriction of the map $G$ to $W_{1}$. Since

$$
\begin{aligned}
& D_{E_{2}+\alpha_{1} E_{1}} G=(\sqrt{5} / 2) F_{\star}\left(E_{2}\right)-(1 / 2) F \times F_{\star}\left(E_{3}\right) \\
& D_{E_{3}+\alpha_{2} E_{1}} G=(\sqrt{5} / 2) F_{\star}\left(E_{3}\right)+(1 / 2) F \times F_{\star}\left(E_{2}\right)
\end{aligned}
$$

we see that $G$ is an immersion from $W_{1}$ into $S^{6}$. Moreover, since

$$
\begin{aligned}
\left((\sqrt{5} / 3) F+(2 / 3) F \times F_{\star} E_{1}\right) \times\left((\sqrt{5} / 2) F_{\star} E_{2}-\right. & \left.(1 / 2) F \times F_{\star} E_{3}\right) \\
& =(\sqrt{5} / 2) F_{\star} E_{3}+(1 / 2) F \times F_{\star} E_{2}
\end{aligned}
$$

we see that $G$ is an almost complex immersion (and hence minimal). A straightforward computation now shows that

$$
\begin{aligned}
& D_{E_{2}+\alpha_{1} E_{1}}\left((\sqrt{5} / 2) F_{\star} E_{2}-(1 / 2) F \times F_{\star} E_{3}\right) \\
& \quad=f\left((\sqrt{5} / 2) F_{\star} E_{3}+(1 / 2) F \times F_{\star} E_{2}\right)-(3 / 2) G+a\left((\sqrt{5} / 2) F \times F_{\star} E_{2}-(1 / 2) F_{\star} E_{3}\right)
\end{aligned}
$$

where $f$ is some function whose precise value is not essential. So, if we put $X=E_{2}+\alpha_{1} E_{1}$ and $Y=E_{3}+\alpha_{2} E_{1}$, we see that $X$ and $Y$ are orthogonal with respect to the induced metric and have the same constant length $\sqrt{3 / 2}$. We also see that

$$
h(X, X)=a\left((\sqrt{5} / 2) F \times F_{\star} E_{2}-(1 / 2) F_{\star} E_{3}\right)
$$

Since $G$ is an almost complex immersion, it follows that

$$
h(X, Y)=h(X, G \times X)=G \times h(X, X)
$$

Hence

$$
h(X, Y)=a\left(-(1 / 2) F_{\star} E_{2}-\left((\sqrt{5} / 2) F \times F_{\star} E_{3}\right)\right)
$$

So, we see that the image of the tangent space and the first normal space to the almost complex immersion are spanned by $F_{\star} E_{2}(q), F_{\star} E_{3}(q), F \times F_{\star}\left(E_{2}\right)(q)$ and $F \times F_{\star} E_{3}(q)$. Therefore, we get that its orthogonal complement in $S^{6}$ is spanned by $F_{\star}\left(E_{1}\right)(q)$ and $(2 / 3) F-$ $(\sqrt{5} / 3) F \times F_{\star}\left(E_{1}\right)$. Hence, the tube on the almost complex immersion $G$ with radius $\gamma$, with $\cos \gamma=\sqrt{5} / 3$ in the direction of the orthogonal complement of the first osculating space is given by (5.14) and corresponds therefore to the original totally real immersion $F$. This completes the proof of Theorem 2.

REMARK 5.4. The above construction can also be applied to the totally real immersion of $S^{3}$ into $S^{6}$ constructed in Example 4.1. However, in that case, the resulting almost complex curve is totally geodesic, and hence it is impossible to define the first normal bundle. Taking coordinates

$$
\begin{aligned}
& y_{1}=\cos (3 t / 2) z_{1}, \\
& y_{2}=-\sin (3 t / 2) z_{1}, \\
& y_{3}=\cos (3 t / 2) z_{3}+\sin (3 t / 2) z_{4}, \\
& y_{4}=-\sin (3 t / 2) z_{3}+\cos (3 t / 2) z_{4},
\end{aligned}
$$

we notice that $\partial / \partial t$ corresponds with the vector field $E_{1}$. Since

$$
\tilde{\psi}_{\star}\left(E_{1}\right)=\left(\begin{array}{c}
(2 / 3) y_{2} \\
y_{1} \\
-(\sqrt{5} / 3) y_{2} \\
-(\sqrt{3} / 3 \sqrt{2}) y_{4} \\
(\sqrt{15} / 3 \sqrt{2}) y_{3} \\
-(\sqrt{15} / 3 \sqrt{2}) y_{4} \\
(\sqrt{3} / 3 \sqrt{2}) y_{3}
\end{array}\right)
$$

a straightforward computation shows that the resulting totally geodesic almost complex curve has components $\left(u_{1}, \ldots, u_{7}\right)$ given by

$$
\begin{aligned}
& u_{1}=(\sqrt{5} / 3)\left(y_{1}^{2}+y_{2}^{2}-y_{3}^{2}-y_{4}^{2}\right), \\
& u_{2}=0, \\
& u_{3}=(2 / 3)\left(y_{1}^{2}+y_{2}^{2}-y_{3}^{2}-y_{4}^{2}\right), \\
& u_{4}=-(2 \sqrt{15} / 3 \sqrt{2})\left(y_{1} y_{3}+y_{2} y_{4}\right), \\
& u_{5}=(2 \sqrt{3} / 3 \sqrt{2})\left(y_{1} y_{4}-y_{2} y_{3}\right), \\
& u_{6}=(2 \sqrt{3} / 3 \sqrt{2})\left(y_{1} y_{3}+y_{2} y_{4}\right), \\
& u_{7}=-(2 \sqrt{15} / 3 \sqrt{2})\left(y_{1} y_{4}-y_{2} y_{3}\right) .
\end{aligned}
$$

Therefore, $\tilde{\psi}\left(S^{3}\right)$ can still be considered as some tube, given by (5.14), on a totally geodesic almost complex curve.

REMARK 5.5. It is clear that the examples satisfying Case (1) and (2) of Theorem 2 do not contain any totally geodesic points (the length of the second fundamental form is strictly greater then a positive constant) and that the eigenvalues of Ric are bounded by a constant strictly smaller than 2 . Therefore, it follows from Lemma 3.2 that these examples can not be put together differentiably with examples $M$ satisfying $\delta_{M}=2$.

REMARK 5.6. We recall that a Riemannian manifold $M$, of any dimension $n$, is locally symmetric when its Riemann-Christoffel curvature tensor $R$ is parallel, i.e., $\nabla R=0$, where $\nabla$ is the Levi Civita connection of its metric $\langle.,$.$\rangle , and that M$ is said to be semi-symmetric
when more generally, $R \cdot R=0$, meaning that $(R(X, Y) \cdot R)\left(X_{1}, X_{2}, X_{3}, X_{4}\right)=0$ for all tangent $X, Y, X_{1}, X_{2}, X_{3}, X_{4}$, where $R(X, Y)=\nabla_{X} \nabla_{Y}-\nabla_{Y} \nabla_{X}-\nabla_{[X, Y]}$ is the curvature operator of $M$. By pseudo-symmetric manifolds, we mean here the further generalisation of locally symmetric manifolds, namely those manifolds $M$ for which $R \cdot R=f Q(\langle.,), R$.$) ,$ where $f: M \rightarrow \boldsymbol{R}$ is a differentiable function and $Q(\langle.,\rangle, R$.$) is defined by$

$$
Q(\langle., .\rangle, R)\left(X, Y, X_{1}, X_{2}, X_{3}, X_{4}\right)=(X \wedge Y) \cdot R\left(X_{1}, X_{2}, X_{3}, X_{4}\right),
$$

where $(X \wedge Y) Z=\langle X, Z\rangle Y-\langle Y, Z\rangle X$, for all vector fields $Z$. From the extrinsic as well as from the intrinsic point of view, this notion turns out to be natural generalisation of localand semi-symmetry; for a survey on this see [Ver]. It is known that a 3-dimensional manifold is pseudo-symmetric if and only if it is quasi-Einstein. Therefore, Theorem 2 also provides a classification of all pseudo-symmetric 3-dimensional totally real submanifolds of $S^{6}$. The examples $M$ with $\delta_{M}=2$ satisfy $R \cdot R=Q(\langle.,), R$.$) , while the examples of Case (1) and$ (2) of Theorem 2 satisfy $R \cdot R=(1 / 16) Q(\langle.,\rangle, R$.$) . So the pseudo-symmetry conditions$ for these submanifolds are realised with constant functions $f$, being respectively $f=1$ and $f=1 / 16$ (the values of which we observe to be precisely the only possibilities for $K$ for totally real immersions of $M^{3}$ into $S^{6}(1)$ with constant sectional curvatures $K$ [E1]).

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