# Correction to: On Regular Fréchet-Lie Groups I 

(Tokyo Journal of Mathematics, vol. 3, no. 2, 1980, pp. 353-390)

Hideki OMORI, Yoshiaki MAEDA and Akira YOSHIOKA

Okayama University, Keio University and Tokyo Metropolitan University

The paper with the above title contains the misprints and an omission. The omisson occurs in Lemma 3.5.

The correction in the statement of Lemma 3.5 is: Page 380 in Lemma 3.5: $\Phi_{2}$ in the statement and the proof should be understood as a mapping involving $X_{1}$-variable, i.e., $\Phi_{2}(x ; \xi)$ should be replaced by $\Phi_{2}\left(x ; X_{1}, \xi\right)$.

By the above reason, the proof of Proposition 4.1 is not correct, for $b_{2}(x ; \tilde{\xi})$ in (53) contains $X_{1}$-variable. This gap is repaired as follows:

Denote $\phi\left(x ; \tilde{\xi}, X_{1}\right)=\left\langle\tilde{\xi} \mid \widetilde{S}\left(x ; X_{1}, \bar{X}_{0}(x ; \xi(x ; \widetilde{\xi}))\right)\right\rangle$, and set

$$
\begin{aligned}
\psi\left(x ; \tilde{\xi}, Y, \zeta, X_{1}\right) & =\langle\zeta \mid Y\rangle+\phi\left(x ; \tilde{\xi}+\zeta, X_{1}\right)-\phi\left(x ; \tilde{\xi}, X_{1}\right) \\
& =\left\langle\zeta \left\lvert\, Y+\int_{0}^{1} \frac{\partial \phi}{\partial \tilde{\xi}}\left(x ; \tilde{\xi}+t \zeta, X_{1}\right) d t\right.\right\rangle .
\end{aligned}
$$

Note that if $\rho=$ id., then $\phi=\left\langle\tilde{\xi} \mid X_{1}\right\rangle$, hence $\psi=\left\langle\zeta \mid Y+X_{1}\right\rangle$. Therefore, one may assume that $\int_{0}^{1} \partial \phi / \partial \tilde{\xi}\left(x ; \tilde{\xi}+t \zeta, X_{1}\right) d t$ is sufficiently close to $X_{1}$ in the $C^{2}$-topology.

By Lemma 3.5, the given operator can be written by

$$
\begin{equation*}
\iint a\left(x ; \tilde{\xi}, X_{1}\right) e^{-i \phi\left(x ; \tilde{\xi}, X_{1}\right)} \nu_{1}(x, z) u(z) \mathrm{d} z \mathrm{~d} \tilde{\xi}, \quad z={ }_{x} X_{1} \tag{1}
\end{equation*}
$$

where $\nu_{1}$ is the cut off function defined in (46). (Cf. (34)~(40)). Since the breadth of $\nu_{1}$ is sufficiently small, one may assume $\phi\left(x ; \tilde{\xi}, X_{1}\right) \equiv$ $\left\langle\tilde{\xi} \mid X_{1}\right\rangle$ for $\left.\left.\left|X_{1}\right|\right\rangle\right\rangle 0$. For amplitude $a \in \widetilde{\Sigma}_{c}^{\beta}$ we consider the following equation:

$$
\begin{equation*}
a\left(x ; \tilde{\xi}, X_{1}\right)=\iint e^{-i \psi\left(x: \tilde{\xi}, Y, \zeta, X_{1}\right)} b(x ; \tilde{\xi}, Y) \mathrm{d} Y \mathrm{~d} \zeta \tag{2}
\end{equation*}
$$

Received September 25, 1981
where $(x ; \tilde{\xi})$ is understood as a parameter. To do that, we have to fix a function space $\widetilde{S}_{c}^{\beta}$ as the totality of $g\left(x ; \xi, X_{1}\right) \in C^{\infty}\left(T^{*} N \oplus T N\right)$ such that $g$ is rapidly decreasing in $X$ and $g$ has the following asymptotic expansion:

$$
g(x ; r \hat{\xi}, X) \sim g_{\beta}(x ; \hat{\xi}, X) \mu(r)^{\beta}+g_{\beta-1}(x ; \hat{\xi}, X) \mu(r)^{\beta-1}+\cdots,
$$

where $g_{\beta-j}$ 's are $C^{\infty}$ functions on $S^{*} N \oplus T N$, rapidly decreasing in $X$.
Note that by virture of the cut off function $\nu$ and the asymptotic expansion (11) (cf. p. 365), one may assume $a\left(x ; \tilde{\xi}, X_{1}\right)$ in (1) is an element of $\widetilde{S}_{c}^{\beta}$. If (2) can be solved in $\widetilde{S}_{c}^{\beta}$ for a given $a \in \widetilde{S}_{c}^{\beta}$, then (1) can be replaced as follows:

$$
\iint a e^{-i \phi} \nu_{1} u \mathrm{~d} z \mathrm{~d} \tilde{\xi}=\int b^{\prime}\left(x ; \xi_{1}\right) \tilde{\nu}\left(\varphi\left(x ; \xi_{1}\right)\right) \mathrm{d} \xi_{1}+(K \circ u)(x),
$$

where

$$
\begin{equation*}
b^{\prime}\left(x ; \xi_{1}\right)=\iint b\left(x ; \xi_{1}-\zeta, Y\right) e^{-i\langle\zeta \mid Y\rangle} \mathrm{d} Y \mathrm{~d} \zeta \tag{3}
\end{equation*}
$$

If $b \in \widetilde{S}_{c}^{\beta}$, then we obtain $b^{\prime} \in \Sigma_{c}^{\beta}$. Thus, we have only to solve (2).
Now, remark that (2) is a Fourier-integral operator of order 0 on $\boldsymbol{R}^{n}$ for each fixed ( $x ; \tilde{\xi}$ ). Apply the adjoint operator to both sides of (2). The left hand side is

$$
\begin{equation*}
b^{\prime \prime}(x ; \tilde{\xi}, Z)=\iint e^{i \psi\left(x ; \tilde{\xi}, z, \eta, X_{1}\right)} a\left(x ; \tilde{\xi}, X_{1}\right) \mathrm{d} X_{1} \mathrm{~d} \eta \tag{4}
\end{equation*}
$$

Since $X_{2}=\int_{0}^{1}(\partial \phi / \partial \tilde{\xi})\left(x ; \tilde{\xi}+t \eta, X_{1}\right) d t$ is sufficiently close to $X_{1}, b^{\prime \prime}$ can be written as

$$
\begin{equation*}
b^{\prime \prime}(x ; \tilde{\xi}, Z)=\iint e^{\iota\left\langle\eta \mid Z+x_{2}\right\rangle} a\left(x ; \tilde{\xi}, X_{1}\left(x ; \tilde{\xi}, \eta, X_{2}\right)\right) \frac{d X_{1}}{d X_{2}}\left(x ; \tilde{\xi}, \eta, X_{2}\right) \mathrm{d} X_{2} \mathrm{~d} \eta \tag{5}
\end{equation*}
$$

and belongs to $\widetilde{S}_{c}^{\beta}$. Thus, we have only to solve

$$
\begin{equation*}
b^{\prime \prime}\left(x ; \tilde{\xi}, Z_{1}\right)=\iiint \int e^{i \psi\left(x: \tilde{\xi}, z, \eta, x_{1}\right)-i \psi\left(x: \tilde{\xi}, Y, \zeta, x_{1}\right)} b(x ; \tilde{\xi}, Y) \mathrm{d} Y \mathrm{~d} \zeta \mathrm{~d} X_{1} \mathrm{~d} \eta \tag{6}
\end{equation*}
$$

Note that (6) is a pseudo-differential equation.
For each fixed ( $x ; \tilde{\xi}, Z, \eta, Y$ ), compute out the critical point and value of the above phase function. Then, using that $\int_{0}^{1}(\partial \phi / \partial \tilde{\xi})(x ; \tilde{\xi}+\zeta+$ $\left.t(\eta-\zeta), X_{1}\right) d t$ is sufficiently close to $X_{1}$, one can obtain that
(7) $\quad b^{\prime \prime}(x ; \tilde{\xi}, Z)=\iint\left[\iint \frac{d\left(\zeta, X_{1}\right)}{d\left(\zeta^{\prime}, X^{\prime}\right)} e^{-i\left\langle\zeta^{\prime} \mid X^{\prime}\right\rangle} \mathrm{d} X^{\prime} \mathrm{d}^{\prime} \zeta^{\prime}\right] e^{-i\langle\eta \mid Y-z\rangle} b(x ; \tilde{\xi}, Y) \mathrm{d} Y \mathrm{~d} \eta$.

Note that

$$
c(x ; \tilde{\xi}, Z, \eta, Y)=\iint \frac{d\left(\zeta, X_{1}\right)}{d\left(\zeta^{\prime}, X^{\prime}\right)} e^{-i\left\langle\zeta^{\prime} \mid X^{\prime}\right\rangle} \mathrm{d} X^{\prime} \mathrm{d} \zeta^{\prime}
$$

is sufficiently close to 1 . Thus, we see that (7) is an invertible pseudodifferential operator. Therefore, $b(x ; \tilde{\xi}, Y)$ is obtained in the following form

$$
\begin{equation*}
b(x ; \tilde{\xi}, Y)=\iint f(x ; \tilde{\xi}, Z, \eta, Y) e^{-i\langle\eta \mid Z-Y\rangle} b^{\prime \prime}(x ; \tilde{\xi}, Z) \mathrm{d} Z \mathrm{~d} \eta \tag{8}
\end{equation*}
$$

Hence by (5), we obtain $b \in \widetilde{S}_{c}^{\beta}$ and $b^{\prime} \in \Sigma_{c}^{\beta}$. This computation is not so easy one, but the tiresome computation using (4)-(8) and (3) leads us directly to the conclusion.

Other miscellaneous errata are as follows:
p. $353 l \uparrow 3: \quad h(s, t) \equiv 1$ should be read $h(0, t) \equiv 1$.
p. $358 l \downarrow 9$ : $b(x, \hat{\xi})$ should be read $b(x ; r \hat{\xi})$.
p. 385 in Proposition 5.3: The domain of integrations should be $T_{x}^{*}$. p. $387 l \downarrow 9$ : The second line in the computation of $I_{\lambda}$ should be read

$$
=\lim _{s \rightarrow \infty} \int_{0}^{1} \int_{S_{x}^{*} N}(i \hat{\xi})^{\beta}(1-\kappa)\left(\frac{1}{i} \frac{d}{d t}\right)^{\lambda}\left(t^{|\beta|+n-1} b(x ; t \hat{\xi})\right) e^{-i t s\left(\hat{\xi}\left|Z-(1 / s) X_{0}^{*}(x ; \hat{\xi}\rangle\right\rangle\right.} \mathrm{d} \hat{\xi} \mathrm{~d} t
$$

## Present Address:

Department of Mathematics College of General Education Tsushimanaka, Okayama 700

Department of Mathematics Keio University Hiyoshi, Kohoku-ku, Yokohama 223 and Department of Mathematics Faculty of Sciences Tokyo Metropolitan University Fukazawa, Setagaya-ku, Tokyo 158

