REGULARIZED SELF-INTERSECTION LOCAL TIMES OF PLANAR BROWNIAN MOTION¹

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Let

$$T_k^{\epsilon}(\lambda; t_1, \ldots, t_k) = \rho(X_{t_1}) q^{\epsilon}(X_{t_2} - X_{t_1}) \cdots q^{\epsilon}(X_{t_k} - X_{t_{k-1}}),$$

where X_t is a Brownian motion in \mathbb{R}^2 , $\lambda(dx) = \rho(x) dx$ and q^{ϵ} converges to Dirac's delta function as $\epsilon \downarrow 0$. The self-intersection local times of order k are described by a generalized random field

$$T_k\big(\lambda;\,t_1,\ldots,t_k\big) = \lim_{\varepsilon \downarrow 0} T_k^\varepsilon\big(\lambda;\,t_1,\ldots,t_k\big), \quad \text{for } 0 < t_1 < \, \cdots \, < t_k.$$

The field "blows up" as $t_i - t_j \to 0$ for some $i \neq j$. We show that with a proper choice of the coefficients $B_k^l(\varepsilon)$, a generalized random field

$$\mathscr{T}_{k}(\lambda; t_{1}, \ldots, t_{k}) = \lim_{\epsilon \downarrow 0} \left[T_{k}^{\epsilon}(\lambda; t_{1}, \ldots, t_{k}) + \sum_{l=1}^{k-1} \left[B_{k}^{l}(\epsilon) T_{l}^{\epsilon} \right] (\lambda; t_{1}, \ldots, t_{k}) \right]$$

is well defined for all $0 \le t_1 \le \cdots \le t_k$ and it coincides with $T_k(\lambda; t_1, \ldots, t_k)$ for $t_1 < \cdots < t_k$.

1. Main results.

1.1. We denote by (X_t, P_{μ}) the Brownian motion in \mathbb{R}^2 with the initial law μ (which can be any σ -finite measure on \mathbb{R}^2). If $0 < t_1 < \cdots < t_n$, then the joint probability density for X_{t_1}, \ldots, X_{t_n} is given by the formula

$$(1.1) p_{\mu}(t,x) = \int \mu(dx_0) p_{t_1}(x_1-x_0) p_{t_2-t_1}(x_2-x_1) \cdots p_{t_n-t_{n-1}}(x_n-x_{n-1}).$$

Here

$$(1.2) p_t(x) = t^{-1}p(x/\sqrt{t}), p(z) = (2\pi)^{-1}e^{-|z|^2/2}.$$

We start from a probability density q(z) on \mathbb{R}^2 such that

(1.3)
$$\int |\ln|x| \, |^k q(x) \, dx < \infty, \quad \text{for all } k > 0,$$

$$\int e^{\beta|x|} q(x) \, dx < \infty, \quad \text{for some } \beta > 0.$$

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Put

$$q^{\varepsilon}(x) = \varepsilon^{-2}q(x/\varepsilon),$$

fix a measure $\lambda(dx) = \rho(x) dx$ and consider a sequence of functions

(1.5)
$$T_{k}^{\epsilon}(\lambda; t_{1}, \dots, t_{k}) = \rho(X_{t_{1}})q^{\epsilon}(X_{t_{2}} - X_{t_{1}}) \cdots q^{\epsilon}(X_{t_{k}} - X_{t_{k-1}}),$$
$$(t_{1}, \dots, t_{k}) \in D_{k}, \text{ for } k = 1, 2, \dots.$$

Here

$$(1.6) D_k = \{0 \le t_1 \le \cdots \le t_k\}.$$

1.2. Let D be a region in \mathbb{R}^k . A generalized random field (g.r.f.) over D is a continuous linear mapping F from a space \mathcal{D} of functions on D (test functions) to a space L of random variables [i.e., measurable functions on $(\Omega, \mathcal{F}, P_\mu)$].

In this paper we use the following test functions. Put

$$\|\varphi\|_{\beta} = \sup_{t \in D} e^{\beta|t|} |\varphi(t)|, \text{ for } \beta > 0, \qquad |t| = |t_1| + \cdots + |t_k|,$$

and denote by $\mathcal{D}^{\beta}(D)$ the set of all functions φ on D which are infinitely differentiable (including the boundary) and satisfy the condition $|\mathbb{D}^l \varphi||_{\beta} < \infty$ for all $l = (l_1, \ldots, l_k)$. Here $\mathbb{D}^l = \mathbb{D}^{l_1}_1 \cdots \mathbb{D}^{l_k}_k$ with $\mathbb{D}_i = \partial/\partial t_i$. For every positive integer n we put

$$\|\varphi\|_{\beta, n} = \sup_{|l| < n} |\mathbb{D}^{l}\varphi\|_{\beta}.$$

The space of test functions $\mathcal{D}(D)$ is the union of $\mathcal{D}^{\beta}(D)$ over all $\beta > 0$. The convergence $\varphi_n \to \varphi$ in $\mathcal{D}(D)$ means that all φ_n and φ belong to the same $\mathcal{D}^{\beta}(D)$ and $\|\varphi_n - \varphi\|_{\beta, n} \to 0$ for all n.

The space L of random variables is defined as the intersection of $L^p(P_\mu)$ over all $p \geq 2$; the convergence $Y_n \to Y$ in L means that $P_\mu |Y_n - Y|^p \to 0$ for all $p \geq 2$.

The formula

$$T_k^{\epsilon}(\lambda; \varphi) = \int_{D_{\epsilon}} T_k^{\epsilon}(\lambda; t) \varphi(t) dt, \qquad \varphi \in \mathscr{D}(D_k),$$

defines a g.r.f. if

- (a) λ has a bounded density; and
- (b) either μ is finite or μ has a bounded density and then λ is finite.

We write $f_1 \simeq f_2$ if $f_1(\varepsilon) - f_2(\varepsilon) = O(|\varepsilon|^{\alpha})$ for some $\alpha > 0$ as $|\varepsilon| \to 0$ [in fact, everywhere we use the notation \simeq , $f_1(\varepsilon) - f_2(\varepsilon) = O(|\varepsilon|^{\alpha})$ for all $0 < \alpha < 1$ and, in many cases, for all $0 < \alpha < 2$].

1.3. Suppose that B is a continuous linear mapping from $\mathcal{D}(D_k)$ to $\mathcal{D}(D_l)$. To every g.r.f. F over D_l there corresponds a g.r.f. BF over D_k defined by the formula $(BF)(\varphi) = F(B\varphi)$.

For every $l \leq k$ we consider a mapping B_k^l from $\mathcal{D}(D_k)$ to $\mathcal{D}(D_l)$ given as

$$(1.7) (B_k^l \varphi)(t_1,\ldots,t_l) = \sum_{\sigma} \varphi(t_{\sigma_1},\ldots,t_{\sigma_k}),$$

where the sum is taken over all mappings σ from $\{1, 2, ..., k\}$ onto $\{1, 2, ..., l\}$ such that $\sigma_i \leq \sigma_i$ for i < j. For instance,

$$\begin{split} \left(B_{k}^{k}\varphi\right)\!(t_{1},\ldots,t_{k}) &= \varphi(t_{1},\ldots,t_{k}), \\ \left(B_{k}^{k-1}\varphi\right)\!(t_{1},\ldots,t_{k-1}) &= \varphi(t_{1},t_{1},t_{2},\ldots,t_{k-1}) + \varphi(t_{1},t_{2},t_{2},\ldots,t_{k-1}) \\ &\quad + \cdots + \varphi(t_{1},t_{2},\ldots,t_{k-1},t_{k-1}), \\ &\vdots \\ \left(B_{k}^{1}\varphi\right)\!(t_{1}) &= \varphi(t_{1},\ldots,t_{1}). \end{split}$$

Our main result is stated in the following theorem.

THEOREM 1.1. Let measures μ , λ satisfy conditions (a), (b) in Section 1.2 and let q satisfy condition (1.3). Put

(1.8)
$$\mathscr{T}_{k}^{\epsilon}(\lambda) = \sum_{l=1}^{k} h_{\epsilon}^{k-l} (B_{k}^{l} T_{l}^{\epsilon})(\lambda),$$

where B_k^l are given by (1.7),

(1.9)
$$h_{\varepsilon} = \frac{1}{\pi} \left\{ \ln \varepsilon + \int \left[C + \ln \frac{|y|}{\sqrt{2}} \right] q(y) \, dy \right\}$$

and C = 0.5772157... is Euler's constant.

There exist generalized random fields $\mathcal{T}_k(\lambda)$ (independent of q) such that

$$(1.10) \mathcal{F}_k(\lambda,\varphi) = \lim_{\epsilon \downarrow 0} \mathcal{F}_k^{\epsilon}(\lambda;\varphi), \quad \text{in L for all $\varphi \in \mathscr{D}(D_k)$.}$$

Moreover, for every $m \geq 2$ and each $\varphi \in \mathcal{D}(D_k)$,

(1.11)
$$P_{\mu}[|\mathcal{T}_{k}(\lambda,\varphi) - \mathcal{T}_{k}^{\epsilon}(\lambda;\varphi)|^{m}] \simeq 0.$$

REMARK. The limit of the field (1.8) exists with h_{ε} replaced by $h_{\varepsilon} + \kappa$ with an arbitrary constant κ . For example, we can take $h_{\varepsilon} = (1/\pi) \ln \varepsilon$. Our choice of κ is made to get the limit independent of q.

1.4. Put $T_k(\varepsilon, \lambda, u) = T_k^{\varepsilon}(\lambda; \psi_{ku})$, $\mathscr{T}_k(\varepsilon, \lambda, u) = \mathscr{T}_k^{\varepsilon}(\lambda; \psi_{ku})$, where $\psi_{ku}(t) = 1_{t_k < u}$, $t \in D_k$, u > 0. Since $B_k^l \psi_{ku} = \begin{bmatrix} k-1 \\ l-1 \end{bmatrix} \psi_{lu}$, we have

(1.12)
$$\mathscr{T}_{k}(\varepsilon,\lambda,u) = \sum_{l=1}^{k} \begin{bmatrix} k-1 \\ l-1 \end{bmatrix} h_{\varepsilon}^{k-l} T_{l}(\varepsilon,\lambda,u).$$

It is proved in Dynkin (1987) that there exist random variables $\mathcal{F}_k(\lambda, u)$ such that

$$\int_0^\infty du \, e^{-ru} P_\mu |\mathcal{T}_k(\varepsilon,\lambda,u) - \mathcal{T}_k(\lambda,u)|^p \simeq 0,$$

for all r > 0 and all $p \ge 2$. This implies

$$P_{\mu} \bigg| \int_0^{\infty} du \, e^{-ru} \mathcal{T}_k(\varepsilon, \lambda, u) - \int_0^{\infty} du \, e^{-ru} \mathcal{T}_k(\lambda, u) \bigg|^p \simeq 0.$$

Note that

$$\int_0^\infty du \, e^{-ru} \mathcal{T}_k(\varepsilon, \lambda, u) = \mathcal{T}_k^{\varepsilon}(\lambda, \varphi_r),$$

where $\varphi_r(t) = r^{-1}e^{-rt_k}$ belongs to $\mathcal{D}(D_k)$. It follows from (1.11) that

(1.13)
$$\mathscr{T}_k(\lambda;\varphi_r) = \int_0^\infty du \, e^{-ru} \mathscr{T}_k(\lambda,u).$$

1.5. The expression for $\mathcal{F}_k^{\epsilon}(\lambda)$ can be simplified by the change of variables

$$v_1 = t_1, \quad v_i = t_i - t_{i-1}, \text{ for } i = 2, ..., k.$$

To every function f there corresponds a function \tilde{f} such that $\tilde{f}(v_1,\ldots,v_k)=f(t_1,\ldots,t_k)$. In particular,

$$(1.14) \quad \tilde{T}_{k}^{\epsilon}(\lambda, \mathbf{v}) = \rho(X_{v_{1}})q^{\epsilon}(X_{v_{1}+v_{2}} - X_{v_{1}}) \cdots q^{\epsilon}(X_{v_{1}+\cdots+v_{k}} - X_{v_{1}+\cdots+v_{k-1}}).$$

Let $\tilde{\varphi}$ be a function of $\mathbf{v}=(v_1,\ldots,v_k)$ and let $\Lambda=\{i_1,\ldots,i_l\}\subset\{1,\ldots,k\}$. We put $v_{\Lambda}=(v_{i_1},\ldots,v_{i_l})$ and we denote by $\tilde{\varphi}(v)_{\Lambda}$ the function of v_{Λ} obtained from $\tilde{\varphi}(v)$ by setting $v_j=0$ for all $j\notin \Lambda$. In this notation we can rewrite (1.7) and (1.8) as

(1.15)
$$\tilde{\mathcal{T}}_{k}^{\epsilon}(\lambda, \tilde{\varphi}) = \sum_{\Lambda} h_{\epsilon}^{k-l} \int dv_{\Lambda} \, \tilde{\varphi}(v)_{\Lambda} \tilde{T}_{l}^{\epsilon}(\lambda, v_{\Lambda}),$$

where Λ runs over all subsets of the set $\{1, \ldots, k\}$ which contain 1 and $l = |\Lambda|$ is the cardinality of Λ .

1.6. We shall see in Section 3 that if φ vanishes near the boundary of D_k , then there exists

(1.16)
$$T_k(\lambda, \varphi) = \lim_{\epsilon \downarrow 0} T_k^{\epsilon}(\lambda, \varphi) \quad \text{in } L$$

and

$$P_{\mu}[T_{k}(\lambda_{1},\varphi_{1})\cdots T_{k}(\lambda_{n},\varphi_{n})]$$

$$(1.17) = \int \mu(dz_0) \lambda_1(dz_1) \cdots \lambda_n(dz_n) \times \varphi_1(t^1) \cdots \varphi_n(t^n) dt^1 \cdots dt^n$$
$$\times m_{k_1 \cdots k_n}(z_0, \dots, z_n; t^1, \dots, t^n).$$

The moment functions $m_{k_1 \cdots k_n}$ can be described as follows. If $t^a = (t_1^a, \dots, t_{k_a}^a)$ and if

$$(1.18) 0 < t_{b_1}^{a_1} < \cdots < t_{b_N}^{a_N}$$

(here $N = k_1 + \cdots + k_n$), then

$$(1.19) m_{k_1 \dots k_n}(z_0, \dots, z_n; t^1, \dots, t^n) = \prod_{i=1}^N p_{u_i}(z_{a_i} - z_{a_{i-1}}),$$

with $a_0 = 0$, and

$$(1.20) u_1 = t_{b_1}^{a_1}, u_2 = t_{b_2}^{a_2} - t_{b_1}^{a_1}, \dots, u_N = t_{b_N}^{a_N} - t_{b_{N-1}}^{a_{N-1}}.$$

If $a_i = a_{i-1}$, then the corresponding factor

(1.21)
$$p_{u_i}(0) = (2\pi u_i)^{-1}$$

is not integrable near the origin and this is the source of the trouble. However, the formula

(1.22)
$$\langle \xi(u_i), \varphi(u_i) \rangle = \int_0^\infty \frac{\varphi(u_i) - e^{-u_i} \varphi(0)}{2\pi u_i} du_i$$

defines a generalized function which can be interpreted as a regularization of (1.21).

It turns out that (1.17) remains true with $T_{k_i}(\lambda, \varphi_i)$ replaced by $\mathcal{F}_{k_i}(\lambda, \varphi_i)$ if we replace in (1.19) every function (1.21) with its regularization (1.22).

1.7. We formulate this statement in a more precise way by using the concept of the direct (tensor) product of generalized functions.

A generalized function F of positive variables u_1, \ldots, u_k is a continuous linear functional on $\mathcal{D}(\mathbb{R}_+^k)$ where $\mathbb{R}_+ = [0, \infty)$. Writing $\langle F(u), \varphi(u) \rangle$ means the same as $F(\varphi)$.

Formula (1.22) defines a generalized function of one positive variable u_i . Another example is the delta function $\langle \delta(u_i), \varphi(u_i) \rangle = \varphi(0)$. The formula

(1.23)
$$\langle F(u), \varphi(u) \rangle = \int f(u) \varphi(u) du, \quad \varphi \in \mathcal{D}(\mathbb{R}^k_+),$$

defines a generalized function of $u = (u_1, \dots, u_k)$ if f is a Borel function with the property

(1.24)
$$\int |f(u)|e^{-\beta|u|} du < \infty, \text{ for every } \beta > 0.$$

Suppose that F_1 is a generalized function of positive variables $\mathbf{u}=(u_1,\ldots,u_k)$ and F_2 is a generalized function of positive variables $\mathbf{v}=(v_1,\ldots,v_l)$. Then there exists a unique generalized function $F(\mathbf{u},\mathbf{v})$ such that

$$\langle F(\mathbf{u}, \mathbf{v}), \varphi(\mathbf{u}, \mathbf{v}) \rangle = \langle F_1(\mathbf{u}), \langle F_2(\mathbf{v}), \varphi(\mathbf{u}, \mathbf{v}) \rangle \rangle,$$

for all $\varphi \in \mathcal{D}(\mathbb{R}^{k+l}_+)$. It is called *the direct product of* F_1 *and* F_2 and it is denoted by $F_1(\mathbf{u}) \times F_2(\mathbf{v})$. [See, e.g., Yosida (1980), Chapter 1, Section 14.] Analogously, we define the direct product $F(\mathbf{u}) \times T(\mathbf{v})$ of a generalized function $F(\mathbf{u})$ and a g.r.f. $T(\mathbf{v})$ and the direct product of two g.r.f.'s.

Suppose that the set $\{1, ..., N\}$ is partitioned into disjoint subsets $I_1, ..., I_r$ and let $u_L = \{u_i, i \in I_l\}$. If F_l is generalized function of u_L , then

$$\prod_{l=1}^r F_l(u_{I_l}) = F_1(u_{I_1}) \times \cdots \times F_r(u_{I_r})$$

is a generalized function of $\mathbf{u} = (u_1, \dots, u_N)$.

Formula (1.15) can be interpreted as

(1.25)
$$\tilde{\mathcal{T}}_{k}^{\epsilon}(\lambda, \mathbf{v}) = \sum_{\Lambda} h_{\epsilon}^{k-l} \delta(v_{B}) \times \tilde{T}_{l}^{\epsilon}(\lambda, v_{\Lambda}),$$

where B is the complement of Λ in $\{1, \ldots, k\}$ and $\delta(v_{\Lambda}) = \prod_{i \in \Lambda} \delta(v_i)$. Heuristically, we can rewrite the sum on the right-hand side as a product

(1.26)
$$\rho(X_{v_1}) \prod_{i=2}^k \left[q^{\epsilon} (X_{v_1 + \dots + v_i} - X_{v_1 + \dots + v_{i-1}}) + h_{\epsilon} \delta(v_i) \right].$$

[In general, the product of generalized functions with the same argument is not defined. In our case it can be *defined* as the sum in (1.25).]

1.8. To prove the results stated in Sections 1.3 and 1.6, first, we investigate the moment functions

(1.27)
$$\mathscr{M}_{k}^{\varepsilon}(\mu,\lambda,q;t) = P_{\mu} \prod_{i=1}^{n} T_{k_{i}}^{\varepsilon}(\lambda_{i},q_{i};t^{i})$$

of the random field (1.5). (We deal here simultaneously with several density functions q and, to avoid confusion, we write q as an extra argument.) Then we study the moments

(1.28)
$$\mathcal{N}_{k}^{\epsilon}(\mu, \lambda, q; \varphi) = P_{\mu} \prod_{i=1}^{n} \mathcal{F}_{k_{i}}^{\epsilon}(\lambda_{i}, q_{i}; \varphi_{i}).$$

Consider the set S of pairs (a, b) which is the union of the disjoint ordered sets

(1.29)
$$S_a = \{(a,1), (a,2), \dots, (a,k_a)\}.$$

Denote by Γ the set of all orderings

(1.30)
$$\gamma = \{(a_1, b_1), \dots, (a_N, b_N)\},$$

of S which are compatible with the order within each subset S_a . To every $\gamma \in \Gamma$ there corresponds a set D_{γ} in \mathbb{R}^N_+ described by (1.18). The union of these sets coincides with $D = D_{k_1} \times \cdots \times D_{k_n}$. Formula (1.20) defines a 1–1 linear mapping C_{γ} from D_{γ} onto \mathbb{R}^N_+ . Put

(1.31)
$$I_{\gamma} = \{1\} \cup \{i: a_i \neq a_{i-1}\},$$

$$(1.32) p_{\gamma}(\mu, \lambda; u_{I_{\gamma}}) = \int \mu(dz_0) \lambda_1(dz_1) \cdots \lambda_n(dz_n) \prod_{i \in I_{\gamma}} p_{u_i}(z_{a_i} - z_{a_{i-1}}),$$

with $a_0 = 0$.

Theorem 1.2. Suppose that we are given, for every $i=1,\ldots,n$, a pair of measures μ , λ_i subject to conditions (a), (b), Section 1.2, and a density q_i which satisfies (1.3). Let $\mathcal{T}_k^{\varepsilon}$ be random fields defined by (1.8). Then for every $\varphi_i \in \mathcal{D}(D_k)$, $i=1,\ldots,n$,

(1.33)
$$\mathcal{N}_{k}^{\epsilon}(\mu,\lambda,q;\varphi) \simeq \sum_{\gamma \in \Gamma} m_{\gamma}(\mu,\lambda)(\tilde{\varphi}_{\gamma}),$$

where

(1.34)
$$\tilde{\varphi}_{\nu}(\mathbf{u}) = \varphi(C_{\nu}^{-1}\mathbf{u}),$$

with $\varphi(t^1,\ldots,t^n)=\varphi_1(t^1)\cdots\varphi_n(t^n)$ and

(1.35)
$$m_{\gamma}(\mu, \lambda; \mathbf{u}) = p_{\gamma}(\mu, \lambda; u_{I_{\gamma}}) \times \prod_{j \notin I_{\gamma}} \xi(u_{j}).$$

Theorem 1.2 will be proved in Section 3 with tools developed in Section 2. We also prove that

(1.36)
$$\lim_{\beta \downarrow 0} \lim_{\epsilon \downarrow 0} P_{\mu} \mathcal{F}_{k}^{\epsilon}(\lambda, \varphi) \mathcal{F}_{k}^{\beta}(\lambda, \varphi) = \lim_{\epsilon \downarrow 0} P_{\mu} \mathcal{F}_{k}^{\epsilon}(\lambda, \varphi) \mathcal{F}_{k}^{\epsilon}(\lambda, \varphi).$$

Theorem 1.1 follows easily from Theorem 1.2 and (1.36).

1.9. Let u = C(t), where C is a unimodular matrix. For every generalized function F(u), the formula

$$\langle \hat{F}(t), \varphi(t) \rangle = \langle F(u), \varphi(C^{-1}(u)) \rangle$$

determines a generalized function $\hat{F}(t)$. We put $\hat{F} = C(F)$ and we say that it is obtained from F by the change of variables u = C(t).

Using this notation, we get from (1.33) and (1.34) the formula

$$(1.37) P_{\mu} \prod_{i=1}^{n} \mathcal{T}_{k_{i}}(\lambda_{i}, t^{i}) = \sum_{\gamma \in \Gamma} C_{\gamma} \left\{ p_{\gamma}(\mu, \lambda, u_{I_{\gamma}}) \times \prod_{j \notin I_{\gamma}} \xi(u_{j}) \right\},$$

which is the promised precise version of the recipe described in Section 1.6.

1.10. Interest in the self-intersections of the Brownian motion has increased considerably in connection with Symanzik's ideas in quantum field theory.

The functionals $\mathcal{T}_2(\lambda, u)$ mentioned in Section 1.4 have been introduced in a pioneering work by Varadhan (1969) published as an appendix to a Symanzik article. They also have been studied by Dynkin (1985, 1987), Le Gall (1985), Rosen (1986a) and Yor (1985a, b, 1986). The functionals $\mathcal{T}_k(\lambda, u)$ for k > 2 first appeared in Dynkin (1984a, b) as a tool for a probabilistic representation of $P(\varphi)_2$ fields. In Dynkin (1986b) various families of functionals which converge to $\mathcal{T}_k(\lambda, u)$ have been investigated and the moment functions of $\mathcal{T}_k(\lambda, u)$ have been evaluated. Theorems 1.1 and 1.2 were, first, announced in Dynkin (1986a).

A different renormalization of the self-intersection local times was proposed in Rosen (1986b) where the existence of an L^2 -limit

$$(1.38) \ I^{k}(B) = \lim_{\epsilon \downarrow 0} \int_{B} \prod_{i=1}^{k} \left[\left[p_{\epsilon} (X_{t_{k}} - X_{t_{k-1}}) \right] - \frac{1}{2\pi (t_{k} - t_{k-1} + \epsilon)} \right] dt_{1}, \dots, dt_{k}$$

was proved for every bounded Borel set $B \subset D_k$. [For k = 3 this is also done in Yor (1985c) by a different method.] Heuristically,

$$\tilde{I}_k(v) = \sum_{\Lambda} (-1)^{k-l} \xi(v_B) \times \tilde{\mathcal{T}}_l(\lambda, v_{\Lambda}),$$

where $I^k(t) = \tilde{I}^k(v)$, λ is Lebesgue measure, Λ and B have the same meaning as in (1.25) and $\xi(V_B) = \prod_{i \in B} \xi(v_i)$.

We refer to Dynkin (1988) for more bibliographical information.

2. Preliminaries.

2.1. Suppose that a generalized function F^{ε} is given for every $\varepsilon \in (0, \varepsilon_0)$. We say that F^{ε} is bounded if, for every test function φ , $F^{\varepsilon}(\varphi)$ is a bounded real-valued function of ε . Let $F^{\varepsilon} = F_1^{\varepsilon} \times F_2^{\varepsilon}$. Standard arguments [see, e.g., Gel'fand and Shilov (1968), Chapter 1, Section 4.4] show that, if F_1^{ε} and F_2^{ε} are bounded, then so is F^{ε} .

Let a^{ϵ} be an arbitrary real-valued function. We write $F^{\epsilon} = O(a^{\epsilon})$ if $F^{\epsilon}/a^{\epsilon}$ is bounded. We write $F_1^{\epsilon} \simeq F_2^{\epsilon}$ if $F_1^{\epsilon} - F_2^{\epsilon} = O(\epsilon^{\alpha})$ for some $\alpha > 0$. (If F^{ϵ} is a real-valued function of ϵ , this is consistent with the notation introduced in Section 1.2.)

2.2. We need some estimates for Green's function.

(2.1)
$$g_{\beta}(x) = \int_0^{\infty} e^{-\beta t} p_t(x) dt.$$

(We drop the subscript β if it is equal to 1.)

LEMMA 2.1. For every $\beta > 0$ and every integer k > 0.

Suppose that a random variable Y has a probability density q which satisfies condition (1.3). Then there exist constants β_k such that

(2.3)
$$E\left[g(\varepsilon Y)^{k}\right] \leq \beta_{k} |\ln \varepsilon|^{k},$$

for all sufficiently small ε . We also have

$$(2.4) Eg(\varepsilon Y) \simeq -h_{\varepsilon}.$$

PROOF. It is well known [see, e.g., Itô and McKean (1965), page 233] that

(2.5)
$$g_{\beta}(x) = \frac{1}{\pi} K_0(\sqrt{2\beta}|x|),$$

where K_0 is a modified Bessel function which can be described [see Watson (1952), 3.71.14 and 3.7.2] by the formula

(2.6)
$$K_0(r) = -I_0(r) \ln \frac{r}{2} + B(r).$$

Here

(2.7)
$$I_0(r) = \sum_{m=0}^{\infty} a_m r^{2m} / (2m)!, \qquad a_m = \begin{bmatrix} 2m \\ m \end{bmatrix} 2^{2-m},$$

(2.8)
$$B(r) = -C + \sum_{i=0}^{\infty} a_m (1 + 1/2 + \cdots + 1/m - C) r^{2m} / (2m!).$$

Formula (2.2) follows from (2.5), (2.6) and (1.3).

Formula (2.6) also implies that

(2.9)
$$\frac{1}{\pi}K_0(2\varepsilon r) = \tilde{h}_{\varepsilon}\varphi_{\varepsilon}(r) + \psi_{\varepsilon}(r),$$

with

$$ilde{h}_{arepsilon} = -rac{1}{\pi} {
m ln} arepsilon, \qquad arphi_{arepsilon}(r) = I_0(2arepsilon r), \qquad \psi_{arepsilon}(r) = rac{1}{\pi} igl[B(2arepsilon r) - I_0(2arepsilon r) {
m ln} \, r igr]$$

and h_s given by (1.9).

Since $a_m \to 0$ and $a_m(1 + 1/2 + \cdots + 1/m - C) \to 0$ as $m \to \infty$, there exist constants $\gamma_1, \gamma_2, \gamma_3$ such that

$$(2.10) \quad \varphi_{\varepsilon}(r) \leq \gamma_1 e^{2\varepsilon r}, \quad |\psi_{\varepsilon}(r)| \leq (\gamma_2 + \gamma_3 |\ln r|) e^{2\varepsilon r}, \quad \text{for all } r > 0.$$

Put $N = |Y| / \sqrt{2}$. By (2.5), (2.6) and (2.9),

$$(2.11) g(\varepsilon Y) = \tilde{h}_{\varepsilon} \varphi_{\varepsilon}(N) + \psi_{\varepsilon}(N) \leq (\gamma_{1} \tilde{h}_{\varepsilon} + \gamma_{2} + \gamma_{3} |\ln N|) e^{2\varepsilon N}.$$

The estimate (2.3) follows from (2.11) and (1.3). By (2.11)

$$(2.12) \quad Eg(\varepsilon Y) = a(\varepsilon)\tilde{h}_{\varepsilon} + b(\varepsilon), \qquad a(\varepsilon) = E\varphi_{\varepsilon}(N), \qquad b(\varepsilon) = E\psi_{\varepsilon}(N).$$

The functions $a(\varepsilon)$ and $b(\varepsilon)$ are even and analytic in a neighborhood of 0. Since a(0)=1, $b(0)=-h_{\varepsilon}-\tilde{h}_{\varepsilon}$ [cf. (1.9)], we have $a(\varepsilon)=1+O(\varepsilon^2)$, $b(\varepsilon)=-h_{\varepsilon}-\tilde{h}_{\varepsilon}+O(\varepsilon^2)$ which implies (2.4). \square

REMARK. Using Hölder's inequality, we conclude from (2.2) that, if λ has a bounded density, then

(2.13)
$$\sup_{x_1,\ldots,x_n} \int \lambda(dz) \prod_{i=1}^n g_{\beta}(x_i-z) < \infty,$$

for every $\beta > 0$.

2.3.

LEMMA 2.2. For every Y described in Lemma 2.1 and for every $\varepsilon > 0$,

(2.14)
$$\langle K^{\epsilon}(u), \varphi(u) \rangle = E \int_{0}^{\infty} du \, \varphi(u) p_{u}(\epsilon Y)$$

is a continuous linear functional on $\mathcal{D}(\mathbb{R}_+)$ and

(2.15)
$$K^{\varepsilon}(u) \simeq -h_{\varepsilon}\delta(u) + \xi(u),$$

where h_{ε} is defined by (1.9) and

(2.16)
$$\langle \xi(u), \varphi(u) \rangle = \int_0^\infty [\varphi(u) - e^{-u} \varphi(0)] / 2\pi u \, du.$$

PROOF. We note that

$$(2.17) R^{\varepsilon}(\varphi) = K^{\varepsilon}(\varphi) + h_{\varepsilon}\varphi(0) - \xi(\varphi) = r^{\varepsilon}\varphi(0) - R_{1}^{\varepsilon}(\varphi),$$

where

$$(2.18) r^{\varepsilon} = Eg(\varepsilon Y) + h_{\varepsilon}$$

and

(2.19)
$$R_1^{\epsilon}(\varphi) = E \int_0^{\infty} \left[1 - e^{-\epsilon^2 Y^2/2u} \right] f(u) du,$$

with

(2.20)
$$f(u) = [\varphi(u) - e^{-u}\varphi(0)]/2\pi u.$$

Since $1 - e^{-a} \le \sqrt{a}$ for every $a \ge 0$, we have $R_1^{\epsilon}(\varphi) \le c_{\varphi} \epsilon$, where

$$c_{\varphi} = \int_{0}^{\infty} |f(u)| (2u)^{-1/2} du E|Y| < \infty.$$

Obviously, this implies (2.15). \square

2.4.

LEMMA 2.3. Let Y be the random variable introduced in Lemma 2.1 and let

(2.21)
$$g_{\beta}^{u}(x) = \int_{0}^{u} e^{-\beta t} p_{t}(x) dt.$$

Then

(2.22)
$$E \int \left[g_{\beta}^{u}(z) - g_{\beta}^{u}(z - \epsilon Y) \right]^{2} dz \simeq 0$$

uniformly in u.

PROOF. The left-hand side in (2.22) is equal to $2E[Q(0) - Q(\varepsilon Y)]$, where

$$Q(y) = \int g_{\beta}^{u}(z)g_{\beta}^{u}(z-y) dz.$$

We have

$$0 \le Q(0) - Q(y) \le (2\pi)^{-1} \int_0^\infty (1 - e^{-y^2/2t}) e^{-\beta t} dt$$
$$\le (2\pi)^{-1} |y| \int_0^\infty (2t)^{-1/2} e^{-\beta t} dt,$$

which implies (2.22). \square

2.5.

LEMMA 2.4. Let h_1, \ldots, h_n be positive Borel functions and let

(2.23)
$$H_r^{\beta}(u) = \int_0^u e^{-\beta t} h_r(t) dt.$$

Suppose that $H_r^{\beta}(\infty) < \infty$ for all $\beta > 0$ and all r. Then for every $\varphi \in \mathcal{D}(D_n)$, there exists a $\beta > 0$ such that

(2.24)
$$\int_{D_n} h_1(t_1) h_2(t_2 - t_1) \cdots h_n(t_n - t_{n-1}) \varphi(t) dt$$
$$= \int_{\mathbb{R}^n} H_1^{\beta}(u_1) \cdots H_n^{\beta}(u_n) \psi(u) du,$$

where

$$(2.25) \quad \psi(u) = \mathbb{D}_1, \dots, \mathbb{D}_n \left[e^{\beta(u_1 + \dots + u_n)} \varphi(u_1, u_1 + u_2, \dots, u_1 + u_2 + \dots + u_n) \right].$$

PROOF. We start with n = 1. Let $\varphi \in \mathbb{D}^{\beta}(\mathbb{R}^+)$ and let $0 < \beta < \beta'$. Integration by parts yields

(2.26)
$$\int_0^c \varphi(t)h_1(t) dt = e^{\beta c}\varphi(c)H_1^{\beta}(c) - \int_0^c (e^{\beta t}\varphi(t))'H_1^{\beta}(t) dt.$$

Since $e^{\beta c} \varphi(c) \to 0$ as $c \to \infty$, we get (2.24). Now using (2.26) we prove that (2.24) holds for n+1 if it holds for n. \square

3. Proofs of Theorems 1.1 and 1.2.

3.1. To evaluate the moment functions $\mathcal{M}_{k}^{\epsilon}(\mu,\lambda,q;t)$ we consider N independent random variables Y_{b}^{a} , $(a,b)\in S$, where Y_{1}^{a} has the distribution λ_{a} and $Y_{2}^{a},\ldots,Y_{k_{a}}^{a}$ are distributed with the density q_{a} . Let

$$(3.1) V_1^a(\varepsilon) = Y_1^a, V_b^a(\varepsilon) = V_{b-1}^a(\varepsilon) + \varepsilon Y_b^a, \text{for } b > 1.$$

Assuming that $\lambda_i(dz_i) = \rho_i(z_i) dz_i$, we have the following expression for the

joint density of $V_b^a(\varepsilon)$,

$$(3.2) q^{\varepsilon}(\lambda; v) = \prod_{a=1}^{n} \rho_a(v_1^a) q_a^{\varepsilon}(v_2^a - v_1^a) \cdots q_a^{\varepsilon}(v_{k_a}^a - v_{k_a-1}^a).$$

By substituting $X_{t_b^a}$ for v_b^a we get

(3.3)
$$\prod_{i=1}^{n} T_{k_i}^{\varepsilon}(\lambda_i, q_i; t^i) = q^{\varepsilon}(\lambda; X_t).$$

If $P_{u}{X_{t} \in dx} = p_{u}(t, x) dx$, then

$$(3.4) \ \mathscr{M}_{k}^{\varepsilon}(\mu,\lambda,q;t) = P_{\mu}q^{\varepsilon}(\lambda,X_{t}) = \int q^{\varepsilon}(\lambda,x)p_{\mu}(t,x) dx = Ep_{\mu}(t,V(\varepsilon)).$$

Note that

(3.5)
$$p_{\mu}(t,x) = \int \mu(dx_0) \prod_{r=1}^{N} p_{u_r} (x_{b_r}^{a_r} - x_{b_{r-1}}^{a_{r-1}}),$$

where $\gamma = \{(a_1, b_1), \dots, (a_N, b_N)\}$ is the ordering of S defined by (1.18), u_1, \dots, u_N are given by (1.20) and $x_{b_0}^{a_0} = x_0$.

It follows from (3.4) that

(3.6)
$$\int_{D} \mathscr{M}_{k}^{\varepsilon}(\mu,\lambda,q;t) \varphi(t) dt = \sum_{\gamma \in \Gamma} F_{\gamma}^{\varepsilon}(\varphi),$$

where

(3.7)
$$F_{\gamma}^{\varepsilon}(\varphi) = \int_{D_{\gamma}} Ep_{\mu}(t, V(\varepsilon)) \varphi(t) dt.$$

3.2. Now we investigate the limit behavior of $F_{\nu}^{\epsilon}(\varphi)$ as $\epsilon \to 0$. Note that

(3.8)
$$p_{\mu}(t, V(\varepsilon)) = \int \mu(dx_0) \prod_{r=1}^{N} p_{u_r}(\eta_r),$$

where

(3.9)
$$\eta_r = V_{b_r}^{a_r}(\varepsilon) - V_{b_{r-1}}^{a_{r-1}}(\varepsilon),$$

with $V_{b_0}^{a_0} = 0$. Put

(3.10)
$$\tilde{\eta}_r = Y_1^{a_r} - Y_1^{a_{r-1}}$$

and compare $F_{\gamma}^{\epsilon}(\varphi)$ with

(3.11)
$$\tilde{F}_{\gamma}^{\varepsilon}(\varphi) = \int_{D_{\gamma}} B^{\varepsilon}(t) \varphi(t) dt,$$

where

(3.12)
$$B^{\varepsilon}(t) = E \int \mu(dx_0) \prod_{i \in I_{\gamma}} p_{u_i}(\tilde{\eta}_i) \prod_{j \notin I_{\gamma}} p_{u_j}(\eta_j)$$

and I_{ν} is defined by (1.31). By Lemma 2.4

$$(3.13) F_{\gamma}^{\epsilon}(\varphi) = E \int \mu(dx_0) \int_{\mathbb{R}^N_+} du \, \psi(u) \prod_{r=1}^N g_{\beta}^{u_r}(\eta_r),$$

(3.14)
$$\tilde{F}_{\gamma}^{\varepsilon}(\varphi) = E \int \mu(dx_0) \int_{\mathbb{R}^N_+} du \, \psi(u) \prod_{i \in I_{\gamma}} g_{\beta}^{u_i}(\tilde{\eta}_i) \prod_{j \notin I_{\gamma}} g_{\beta}^{u_j}(\eta_j),$$

with g^u_{β} and ψ given by (2.21) and (2.25).

For the sake of brevity we denote by M the product of the measure P with respect to which the mathematical expectations are taken in (3.13) and (3.14) and the measure $\mu(dx_0)|\psi(u)|\,du$. We put $A_i=g^{u_i}_{\beta}(\eta_i)$, $\tilde{A_i}=g^{u_i}_{\beta}(\tilde{\eta}_i)$ and $\Delta_i=|A_i-\tilde{A_i}|$. Note that

$$|F_{\gamma}^{\epsilon}(\varphi) - F_{\gamma}^{\epsilon}(\tilde{\varphi})| \leq \int_{L} \int \prod_{j \notin I_{\gamma}} A_{j} \prod_{l \in L} \Delta_{l} \prod_{i \in I_{\gamma} \setminus L} \tilde{A}_{i} dM,$$

where the sum is taken over all nonempty $L \subset I_{\gamma}$. By Hölder's inequality each integral does not exceed a product of powers of the integrals

$$\int \! A_{j}^{k_{j}} \, dM, \qquad \int \! ilde{A}_{i}^{k_{i}} \, dM, \qquad \int \! \Delta_{l}^{k_{l}} \, dM,$$

and we can choose $k_l = 2$ for one l. Using estimates (2.3) and (2.22) we prove that

(3.15)
$$\tilde{F}_{\gamma}^{\varepsilon}(\varphi) \simeq F_{\gamma}^{\varepsilon}(\varphi).$$

By (3.1), $\eta_j = \varepsilon Y_{b_j}^{a_j}$ for $j \notin I_{\gamma}$. Since the random variables $\tilde{\eta}_i$, $i \in I_{\gamma}$, and the family $\{\eta_j, j \notin I_{\gamma}\}$ are mutually independent, we have from (3.6), (3.9), (3.10), (3.12) and (3.15) that

$$(3.16) \int_{D}^{\mathcal{A}_{k}^{\varepsilon}}(\mu, \lambda, q; t)\varphi(t) dt$$

$$\simeq \sum_{\gamma \in \Gamma} \int_{D_{\gamma}} dt \, \varphi(t) E \left[\int \mu(dx_{0}) \prod_{i \in I_{\gamma}} p_{u_{i}}(Y_{1}^{a_{r}} - Y_{1}^{a_{r-1}}) \right] \prod_{j \notin I_{\gamma}} E p_{u_{j}}(\varepsilon Y_{2}^{a_{j}}),$$

where according to (1.20)

$$(3.17) u_i = t_{b_i}^{a_i} - t_{b_{i-1}}^{a_{i-1}}.$$

3.3. The limit behavior of

(3.18)
$$\mathcal{M}_{k}^{\epsilon\alpha}(\mu,\lambda;t) = P_{\mu}T_{k_{1}}^{\epsilon}(\lambda_{1},t^{1})T_{k_{2}}^{\beta}(\lambda_{2},t^{2}),$$

as $\varepsilon \downarrow 0$ and β is fixed can be investigated in a similar way. We repeat the arguments in Sections 3.1 and 3.2 with random variables $V_b^2(\varepsilon)$ replaced by $V_b^2(\beta)$ and the set I_{γ} replaced by J_{γ} defined by the condition: $i \notin J_{\gamma}$ if and only if i=1 or $a_{i-1}=a_i=1$. In this way we get that

(3.19)
$$\mathcal{M}_{k}^{\epsilon\alpha}(\mu,\lambda;t) = \sum_{\gamma \in \Gamma} \phi_{\gamma}^{\epsilon\beta}(\varphi) + R(\beta,\epsilon).$$

Here $\varepsilon^{-\alpha}R(\beta,\varepsilon)\to 0$ (with some $\alpha>0$) as $\varepsilon\downarrow 0$ and β is fixed, and

$$(3.20) \qquad \phi_{\gamma}^{\epsilon\beta}(\varphi) = E \left[\int \mu(dx_0) \prod_{i \in J_{\gamma}} p_{u_i} \left(\tilde{V}_{b_i}^{a_i} - \tilde{V}_{b_{i-1}}^{a_{i-1}} \right) \right] \prod_{j \notin J_{\gamma}} E p_{u_j} \left(\varepsilon Y_2^1 \right),$$

where $\tilde{V}_h^1 = Y_1^1$ and $\tilde{V}_h^2 = V_h^2(\beta)$.

3.4. We use (1.25) to express $\mathcal{N}_k^{\epsilon}(\mu, \lambda, q; \varphi)$ through the moment functions $\mathcal{M}_l^{\epsilon}(\mu, \lambda, q; t)$. To use (1.25) as it is stated we need to introduce a new set of variables $v_b^a = t_b^a - t_{b-1}^a$ besides two sets $\{u_i\}$ and $\{t_b^a\}$ we deal with. To avoid cumbersome notation we prefer to return in (1.25) to original variables $t_i = v_i + \cdots + v_i$ and to use the following rules in manipulating with the δ 's:

$$(3.21) f(v)\delta(v) = f(0)\delta(v),$$

$$(3.22) F(s)\delta(t-s) = F(t)\delta(t-s)$$

For instance,

$$\begin{split} \delta(v_2)q\big(X_{v_1+v_3}-X_{v_1}\big) &= \delta(v_2)q\big(X_{v_1+v_2+v_3}-X_{v_1+v_2}\big) \\ &= \delta(t_2-t_1)q\big(X_{t_2}-X_{t_2}\big). \end{split}$$

Following these rules, we rewrite (1.25) in the form

(3.23)
$$\mathcal{F}_{k}^{\epsilon}(\lambda, t) = \sum_{\Lambda} \prod_{i=2}^{k} \left[h_{\epsilon} \delta(t_{i} - t_{i-1}) \right] T_{l}^{\epsilon}(\lambda, t_{\Lambda}).$$

This implies

$$(3.24) \qquad \mathscr{N}_{k}^{\epsilon}(\mu,\lambda,q;\varphi) = \sum_{\Lambda_{-}} \cdots \sum_{\Lambda_{-}} E \prod_{a=1}^{n} \left\{ T_{l_{a}}^{\epsilon}(\lambda_{a},t_{\Lambda_{a}}^{a}) \prod_{b=2}^{k_{a}} \left[h_{\epsilon} \delta(t_{b}^{a}-t_{b-1}^{a}) \right] \right\},$$

where Λ_a runs over all subsets of the set $\{1, \ldots, k_a\}$ which contain 1.

3.5. To simplify the presentation we say that the elements of $S_a = \{(a,1),\ldots,(a,k_a)\}$ have color a. By identifying $b \in \Lambda_a$ with the pair (a,b) we imbed Λ_a into S_a . The terms on the right-hand side of (3.24) are in a 1-1 correspondence with subsets $\Lambda = \Lambda_1 \cup \cdots \cup \Lambda_n$ of S which contain the first elements of each color. Denote the collection of all such sets by $\mathscr L$ and consider, for every $\Lambda \in \mathscr L$, the set Γ_Λ of all orderings of Λ which agree with the given order within each set S_a . Define the characteristic set for a pair (Λ, γ) , $\Lambda \in \mathscr L$, $\gamma \in \Gamma_\Lambda$, as the set of all $s \in \Lambda$ whose color is different from the color of the left neighbor (in Λ relative to γ). We say that an ordering $\tilde{\gamma} \in \Gamma = \Gamma_S$ is the standard continuation of $\gamma \in \Gamma_\Lambda$ if both orderings agree on Λ and if the characteristic sets for $(S,\tilde{\gamma})$ and (Λ,γ) coincide.

We claim that every $\gamma \in \Gamma_{\Lambda}$ has a unique standard continuation. Indeed, suppose that the characteristic set of (Λ, γ) is I and that $\alpha_1 < \cdots < \alpha_r$ are all red elements of I. The set of all red elements in Λ is the union of disjoint intervals $[\alpha_1, \tilde{\alpha}], \ldots, [\alpha_r, \tilde{\alpha}_r]$. The right neighbor β_i of $\tilde{\alpha}_i$ in Λ , obviously, belongs

to I. Since the first red elements of S belongs to Λ , each red element s of S satisfies the relation $\alpha_i \leq s \leq \alpha_{i+1}$ for some i or the relation $\alpha_r \leq s$. The only way to avoid expanding the characteristic set is to put s between α_i and β_{i+1} in the first case or after α_r in the second case. The exact place for s in now uniquely determined by the ordering of red elements.

3.6. It follows from (3.24) and (3.16) that

(3.25)
$$\mathcal{N}_{k}^{\varepsilon}(\mu, \lambda, q; \varphi) \simeq \sum_{\Lambda \in L} \sum_{\gamma \in \Gamma_{\Lambda}} \int_{D_{\gamma}} \Psi_{\gamma\Lambda}^{\varepsilon}(t) \varphi(t) dt,$$

where

(3.26)
$$\Psi_{\gamma\Lambda}^{\varepsilon}(t) = A(t)C^{\varepsilon}(t) \prod_{i \in \Lambda} \left[h_{\varepsilon} \delta(t_{i-1}^{a}, t_{i}^{a}) \right],$$

with

(3.27)
$$A(t) = \int \mu(dx) p_{t(a_1, b_1)}(Y_1^a - x) \prod_{a \neq \tilde{a}} p_{t(a, i) - t(\tilde{a}, j)}(Y_1^{\tilde{a}} - Y_1^a),$$

$$C^{\epsilon}(t) = \prod_{a = \tilde{a}} E p_{t(a, i) - t(\tilde{a}, j)}(\epsilon Y_2^a)$$

[for typographical reasons we write t(a,i) instead of t_i^a]. In (3.26) and (3.27) (\tilde{a},j) is the left neighbor of (a,i) in Λ relative to γ . By the rule (3.22) we can replace $t(\tilde{a},j)$ by $t(\tilde{a},\tilde{j})$, where (\tilde{a},\tilde{j}) is the left neighbor of (a,i) in S relative to the standard continuation $\tilde{\gamma}$ of γ . [Note that the left neighbors of (a,i) in Λ and S have the same color.]

The set of variables $\{t(a,i)-t(\tilde{a},\tilde{j}),\ a\neq\tilde{a}\}$ coincides with the set $\{u_i,\ I\in I_\gamma\}$ and the set $\{(t(a,i)-t(\tilde{a},\tilde{j}),\ a=\tilde{a}\}$ coincides with $\{u_i,i\notin I_\gamma\}$. Thus the sum of terms in (3.25) corresponding to the pairs (Λ,γ) with a given characteristic set is equal to

(3.28)
$$p_{\gamma}(\mu, \lambda; u_{I_{\gamma}}) \prod_{i \notin I_{\gamma}} E[p_{u_{i}}(\epsilon Y_{2}^{a}) - h_{\epsilon} \delta(u_{i})],$$

where the first factor is defined by (1.32). By the criterion (1.24), this factor can be interpreted as a generalized function in $u_{I\gamma}$. Indeed it follows from Hölder's inequality and (2.13) that under the conditions of Theorem 1.2,

$$\int \mu(dz_0)\lambda_1(dz_1)\cdots\lambda_n(dz_n)\prod_{i\in I_{\gamma}}g_{\beta}(z_{a_i}-z_{a_{i-1}})<\infty,$$

for every $\beta > 0$.

Formulas (3.25) and (3.28) and Lemma 1.2 imply Theorem 1.2.

3.7. Using the expression for $P_{\mu}T_{k_1}^{\epsilon}(\lambda_1, \varphi_1)T_{k_2}^{\epsilon}(\lambda_2, \varphi_2)$ given in Section 3.3, we prove that

$$P_{\mu}\mathcal{F}_{k_1}^{\epsilon}(\lambda_1,\varphi_1)T_{k_2}^{\epsilon}(\lambda_2,\varphi_2) \simeq \sum_{\gamma \in \Gamma} \int f_{\gamma}(\mu,\lambda;u_{J_{\gamma}}) \prod_{j \notin J_{\gamma}} \xi(u_j) \tilde{\varphi}_{\gamma}(u) du + R',$$

where $R' = o(\varepsilon^{\alpha})$ for some $\alpha > 0$ and

(3.29)
$$f_{\gamma}(\mu, \lambda; u_{J_{\gamma}}) = E \left[\int u(dx_0) \prod_{i \in J_{\gamma}} p_{u_i} (\tilde{V}_{b_i}^{a_i} - \tilde{V}_{b_{i-1}}^{a_{i-1}}) \right].$$

Using this formula we prove (1.36).

3.8.

PROOF OF THEOREM 1.1. Note that

(3.30)
$$\mathscr{T}_{1}(\lambda,\varphi) = T_{1}(\lambda,\varphi) = \int_{0}^{\infty} \rho(X_{t})\varphi(t) dt$$

do not depend on ε and q. Consider the set \mathscr{A} of all products $A = \mathscr{T}_1(\lambda_1, \varphi_1) \cdots \mathscr{T}_n(\lambda_n, \varphi_n)$, where $n = 1, 2, \ldots, \varphi_i \in \mathscr{D}(D_1)$ and λ_i are finite measures with continuous densities. It is easy to check that the span of \mathscr{A} is everywhere dense in $L^2(P_n)$.

Let $F_{\varepsilon} = \mathcal{F}_{k}^{\varepsilon}(\lambda, q; \varphi)$ with λ, μ, q subject to the conditions of Theorem 1.1. By Theorem 1.2, for every $A \in \mathcal{A}$, $P_{\mu}AF_{\varepsilon}$ converges to a finite limit as $\varepsilon \downarrow 0$; besides $P_{\mu}F_{\varepsilon}^{\varepsilon}$ is bounded. Hence F_{ε} converges weakly to an element F of $L^{2}(P_{\mu})$. By (1.36)

$$\lim_{\epsilon\downarrow 0} P_{\boldsymbol{\mu}} F_{\epsilon}^2 = \lim_{\beta\downarrow 0} \lim_{\epsilon\downarrow 0} P_{\boldsymbol{\mu}} F_{\epsilon} F_{\beta} = \lim_{\beta\downarrow 0} P_{\boldsymbol{\mu}} F F_{\beta} = P_{\boldsymbol{\mu}} F^2.$$

Hence $P_{\mu}(F_{\epsilon}-F)^2 \to 0$. By Theorem 1.2, $P_{\mu}(F-F_{\epsilon})^p$ is bounded for every $p \ge 1$, and by the Schwarz inequality

$$P_{\mu}|F-F_{\varepsilon}|^{m} \leq \operatorname{const}\left[P_{\mu}|F-F_{\varepsilon}|^{2}\right]^{1/2}$$

for every $m \geq 2$, which proves (1.11). For every $A \in \mathcal{A}$, $P_{\mu}FA$ does not depend on q. Therefore the same is true for F. \square

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REFERENCES

DYNKIN, E. B. (1984a). Local times and quantum fields. In Seminar on Stochastic Processes, 1983 (E. Çinlar, K. L. Chung and R. K. Getoor, eds.) 69–84. Birkhäuser, Boston.

DYNKIN, E. B. (1984b). Polynomials of the occupation field and related random fields. J. Funct. Anal. 58 20-52.

DYNKIN, E. B. (1985). Random fields associated with multiple points of the Brownian motion. J. Funct. Anal. 62 397-434.

DYNKIN, E. B. (1986a). Generalized random fields related to self-intersections of the Brownian motion. Proc. Nat. Acad. Sci. U.S.A. 83 3575-3576.

DYNKIN, E. B. (1986b). Functionals associated with self-intersections of the planar Brownian motion. Séminaire de Probabilités XX, 1984 / 85. Lecture Notes in Math. 1204 553-573. Springer, Berlin.

- DYNKIN, E. B. (1987). Self-intersection local times, occupation fields and stochastic integrals. Adv. in Math. 65 254-271.
- DYNKIN, E. B. (1988). Self-intersection gauge for random walks and for Brownian motion. *Ann. Probab.* 16 1-57.
- GEL'FAND, I. M. and SHILOV, G. E. (1968). Generalized Functions. Spaces of Fundamental and Generalized Functions 2. Academic, New York.
- Itô, K. and McKean, H. P., Jr. (1965). Diffusion Processes and Their Sample Paths. Springer, Berlin.
- LE GALL, J. F. (1985). Sur le temps local d'intersection du mouvement brownien plan et la méthode de renormalisation de Varadhan. Séminaire de Probabilités XIX, 1983 / 84. Lecture Notes in Math. 1123 314-331. Springer, Berlin.
- Rosen, J. (1986a). Tanaka's formula and renormalization for intersections of planar Brownian motion. *Ann. Probab.* 14 1245–1251.
- Rosen, J. (1986b). A renormalized local time for multiple intersections of planar Brownian motion. Séminaire de Probabilités XX, 1984 / 85. Lecture Notes in Math. 1204 515-531. Springer, Berlin.
- Varadhan, S. R. S. (1969). Appendix to "Euclidean quantum field theory," by K. Symanzik. In *Local Quantum Theory* (R. Jost, ed.). Academic, New York.
- WATSON, G. N. (1952). A Treatise on the Theory of Bessel Functions, 2nd ed. Cambridge Univ. Press, Cambridge.
- Yor, M. (1985a). Compléments aux formules de Tanaka-Rosen. Séminaire de Probabilités XIX, 1983 / 84. Lecture Notes in Math. 1123 332-349. Springer, Berlin.
- Yor, M. (1985b). Renormalisation et convergence en loi pour les temps locaux d'intersection du mouvement brownien dans \mathbb{R}^2 . Unpublished.
- YOR, M. (1985c). Renormalization results for some triple integrals of two-dimensional Brownian motion. Unpublished.
- Yor, M. (1986). Sur la représentation comme intégrales stochastiques des temps d'occupation du mouvement brownien dans ℝ^d. Séminaire de Probabilités XX, 1984 / 85. Lecture Notes in Math. 1204 543-552. Springer, Berlin.
- Yosida, K. (1980). Functional Analysis, 6th ed. Springer, Berlin.

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