Structure of the scattering operator for time-periodic Schrödinger equations

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§1. Introduction.

In this paper we study the structure of the scattering operator for timeperiodic Schrödinger equations with period ω :

(1.1)
$$i\frac{\partial}{\partial t}\psi(t, x) = (-\Delta + V(t, x))\psi(t, x), \qquad \psi(t, \cdot) \in \mathcal{H} = L^2(\mathbf{R}^n),$$

(1.2)
$$V(t+\boldsymbol{\omega}, x) = V(t, x) \in \boldsymbol{R}, \quad (t \in \boldsymbol{R}, x \in \boldsymbol{R}^n).$$

Under suitable conditions on V(t, x) to be specified below, (1.1) generates a unitary evolution operator U(t, s), $-\infty < t$, $s < \infty$, and for each $s \in \mathbf{R}$, the wave operators defined by

(1.3)
$$W_{\pm}(s) = \operatorname{s-lim}_{t \to \pm \infty} U(s, t) e^{-i(t-s)H_0}, \qquad H_0 = -\Delta,$$

exist and are complete: $\operatorname{Ran} W_{\pm}(s) = \mathcal{H}^{\operatorname{ac}}(U(s + \omega, s))$ (see Yajima [14], Howland [6], Kitada-Yajima [8], Nakamura [11]). Then the scattering operator defined by

(1.4)
$$S(s) = W_+(s)^* W_-(s)$$

is unitary, and by virtue of the time-periodicity, it satisfies

$$S(s)e^{-i\omega H_0} = e^{-i\omega H_0}S(s).$$

It follows that, if we denote a spectral representation of H_0 by $(\tilde{F}(\lambda), \mathcal{X}(\lambda), d\lambda)$, S(s) is decomposed as

(1.6)
$$\begin{cases} S(s) = \sum_{\mu \in \mathbb{Z}} S_{\mu}, \\ \tilde{F}\left(\lambda - \frac{2\pi}{\omega}\mu\right) S_{\mu}\psi = \tilde{S}_{\mu}(\lambda)\tilde{F}(\lambda)\psi \quad (\text{a.e. }\lambda), \\ \tilde{S}_{\mu}(\lambda) \in B\left(\mathfrak{X}(\lambda), \ \mathfrak{X}\left(\lambda - \frac{2\pi}{\omega}\mu\right)\right), \end{cases}$$

for any $\psi \in \mathcal{H}$. We call $\{\widetilde{S}_{\mu}(\lambda)\}$ S-matrices (see §2 for details).

In this paper we are concerned with the structure of S(s) or $\{\tilde{S}_{\mu}(\lambda)\}\)$, and show that the decay of $\tilde{S}_{\mu}(\lambda)$ as μ tends to infinity is completely determined by