

On the equations of one-dimensional motion of compressible viscous fluids

By

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§1. Introduction

We consider the equations of the one-dimensional motion of a compressible, viscous and heat-conductive fluid in Lagrangian coordinates:

$$(1.1) \quad \begin{cases} \rho_t + \rho^2 u_x = 0, \\ u_t + p_x = (\mu \rho u_x)_x, \\ \theta_t + \frac{\theta p_\theta}{c_V} u_x = \frac{1}{c_V} \{(\kappa \rho \theta_x)_x + \mu \rho u_x^2\}, \end{cases}$$

where t is time and x denotes the Lagrangian mass coordinate. Here the unknown functions ρ , u and θ represent the density, velocity and absolute temperature of the fluid; the pressure p and the heat capacity at constant volume c_V are related to the thermodynamic quantities $\rho > 0$ and $\theta > 0$ by the equations of state, and p_θ denotes $\partial p / \partial \theta$; μ and κ are the coefficients of viscosity and heat conduction respectively.

We assume the following conditions on the system (1.1).

A_1 : $p = p(\rho, \theta)$ and $c_V = c_V(\rho, \theta)$ are smooth functions of $(\rho, \theta) \in \mathcal{D}_{\rho, \theta} \equiv \{\rho > 0, \theta > 0\}$ and satisfy the general equations of state on $\mathcal{D}_{\rho, \theta}$, that is,

$$p_\rho \equiv \frac{\partial p}{\partial \rho} > 0, \quad c_V > 0, \quad \frac{\partial c_V}{\partial \rho} = -\theta p_{\theta\theta} / \rho^2,$$

where $p_{\theta\theta} = \partial^2 p / \partial \theta^2$.

A_2 : $\mu = \mu(\rho)$ and $\kappa = \kappa(\rho)$ are smooth functions of $\rho > 0$ (independent of $\theta > 0$) and satisfy $\mu > 0$ and $\kappa > 0$ for $\rho > 0$.

The assumptions $p_\rho > 0$ and $c_V > 0$ imply that the system (1.1) with $\mu = \kappa = 0$ is hyperbolic, while the assumptions $\mu > 0$ and $\kappa > 0$ imply that the equations of u and θ in (1.1) are parabolic.